



Anomalies in rare B decays

TU Dresden,
October 19, 2017

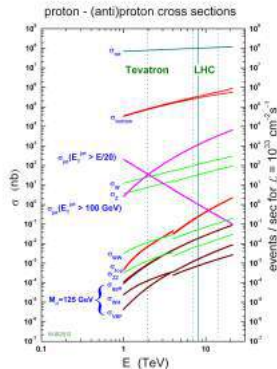
Michel De Cian, University of Heidelberg

B physics

	mass	charge	spin					
	$\sim 2.4 \text{ MeV}/c^2$	$2/3$	$1/2$	u	$\sim 1.275 \text{ GeV}/c^2$	$2/3$	$1/2$	c
				up				charm
	$\sim 172.46 \text{ GeV}/c^2$	$2/3$	$1/2$	t	$\sim 125.09 \text{ GeV}/c^2$	0	0	H
				top				Higgs
	$\sim 4.8 \text{ MeV}/c^2$	$-1/3$	$1/2$	d	$\sim 4.18 \text{ GeV}/c^2$	0	1	g
				down				gluon
	$\sim 511 \text{ MeV}/c^2$	-1	$1/2$	e	$\sim 91.18 \text{ GeV}/c^2$	0	1	Z
				electron				Z boson
	$\sim 105.67 \text{ MeV}/c^2$	-1	$1/2$	μ	$\sim 1.7768 \text{ GeV}/c^2$	1	1	W
				muon				W boson
	$\sim 1.7 \text{ MeV}/c^2$	0	$1/2$	ν_μ	$\sim 15.5 \text{ MeV}/c^2$	0	1	W
				muon neutrino				W boson
	$\sim 1.7768 \text{ GeV}/c^2$	0	1	ν_τ	$\sim 80.38 \text{ GeV}/c^2$	1	1	W
				tau neutrino				W boson
	$\sim 2.2 \text{ eV}/c^2$	0	$1/2$	ν_e	$\sim 1.7768 \text{ GeV}/c^2$	1	1	W
				electron neutrino				W boson

$$B^0 = |d\bar{b}\rangle, B^+ = |u\bar{b}\rangle, B_s^0 = |s\bar{b}\rangle, \Lambda_b^0 = |udb\rangle$$

- B physics is the study of bound states containing one b quark and their decays / dynamics.
- They decay in a multitude of final states, allowing the study of a wide range of physics.
- They are copiously produced at the LHC.



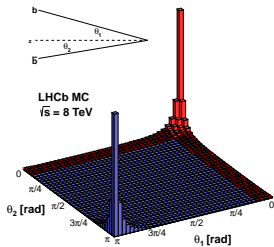
LHC physics



- Take a bunch of protons
- Smash them together and create a mess.
- Spend some millions to build a device to understand it.

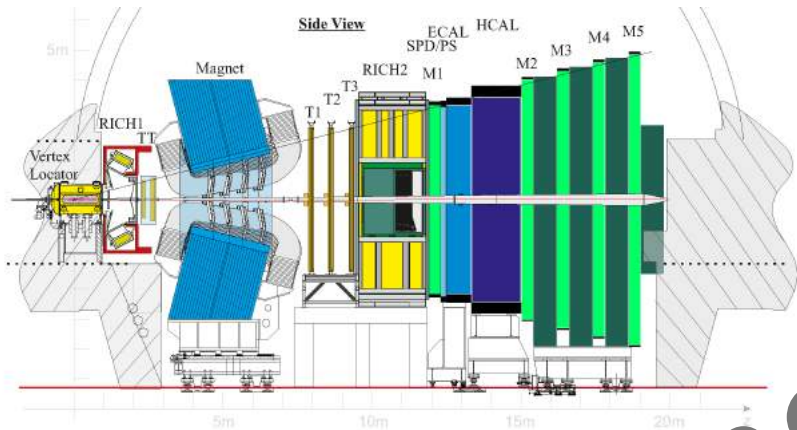


B physics at the LHC

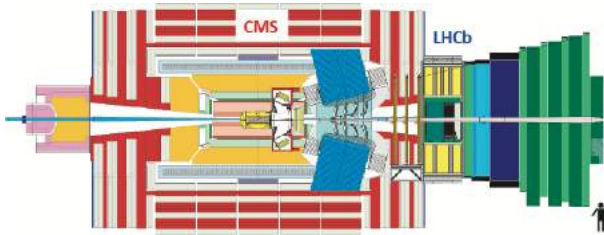


- b hadrons are moderately heavy ($\sim 5 \text{ GeV}$, $> \Lambda_{QCD}$) and are mainly produced in the forward or backward direction at the LHC
→ build a forward detector.
- B hadrons have "soft" decay products and travel $\sim 1 \text{ cm}$ before decaying
→ build a detector with low- p_T capability and good momentum / vertex resolution.
- B hadrons have a large variety of decay channels with different particle species in the final states
→ need a particle ID.

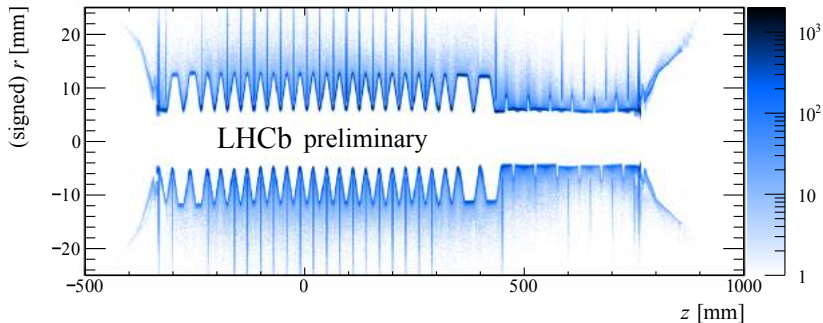
LHCb



LHCb: Comparison with CMS

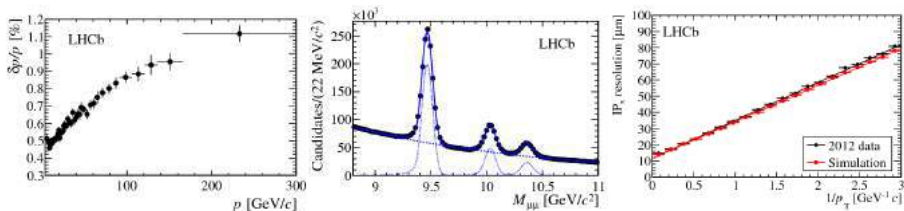


LHCb: Vertex Locator

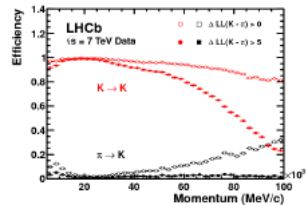


- Velo sensors 8mm from the beam position.
- Allows for very good Impact Parameter and Primary / Secondary Vertex resolution.

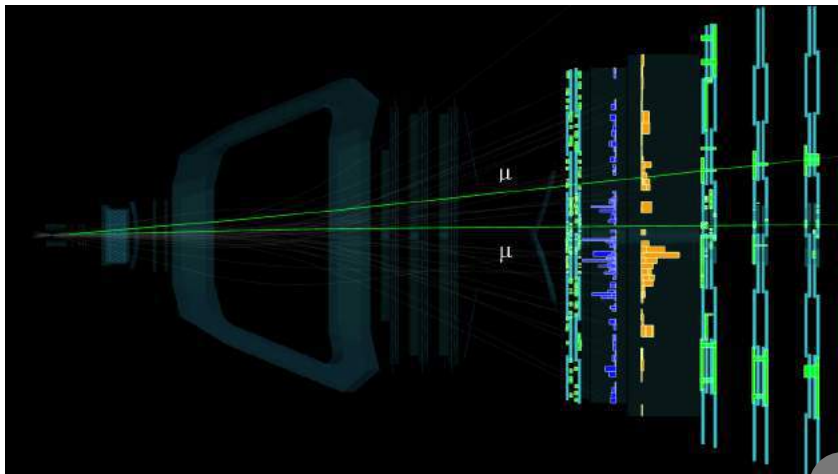
LHCb: Performance numbers



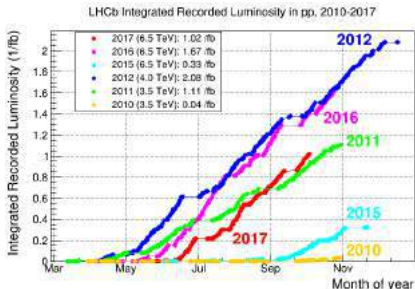
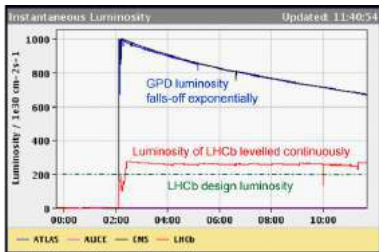
- Excellent momentum / mass resolution:
 - $\frac{\delta p}{p} = 0.5\%$ (10 GeV/c) - 1.0% (200 GeV/c)
 - $\sigma_m(B_s^0 \rightarrow \mu^+ \mu^-) \approx 20 \text{ MeV}/c^2$
- Impact parameter resolution:
 - $15 + 29/p_T$ [GeV/c] μm
- High particle identification efficiency.
 - $\varepsilon_\mu \approx 97\%$ with 1-3% $\pi \rightarrow \mu$ misidentification
 - $\varepsilon_K \approx 95\%$ with $\approx 5\%$ $\pi \rightarrow K$ misidentification



LHCb: Event display

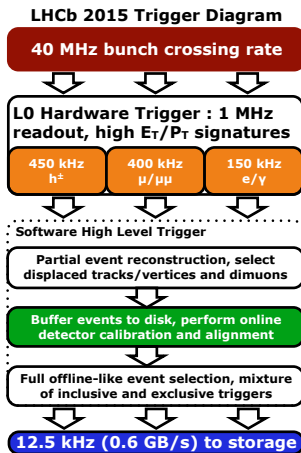


Luminosity levelling



- LHCb does not run at the maximum instantaneous luminosity, as the occupancy in the forward region would be too high.
 - $\mu \approx 1.1$ for 25ns running.
- Luminosity for LHCb is leveled such that it is constant within a fill.
- Achieved by displacing the beams.

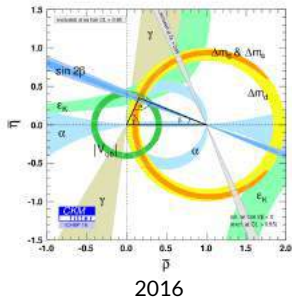
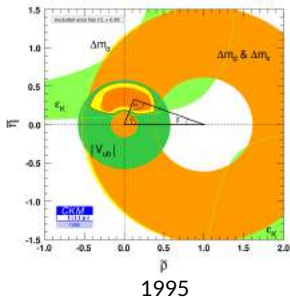
Trigger (Run II)



- Have the same reconstruction (charged and neutral particles) in the software trigger and offline.
- Perform a alignment & calibration after first stage of software trigger, *i.e.* automatically.

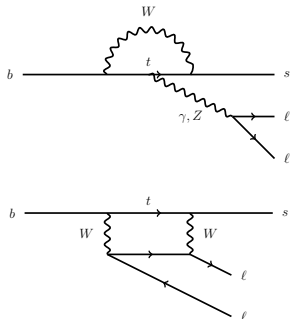
Rare B Decays

B decays as a laboratory



- B hadrons are a perfect laboratory to perform measurements of many fundamental physics quantities.
- Decays governed by (electro)weak interaction, but hadronic state itself by strong interaction.
- B physics is by definition flavour physics and strongly linked to the CKM matrix.

Rare B decays

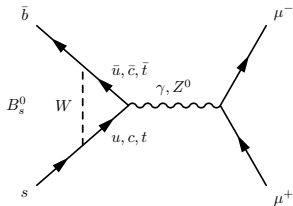


		neutral current		
charged current ↑ Quarks	2.4 MeV $\frac{2}{3}$ $\frac{1}{2}$ u up	1.27 GeV $\frac{2}{3}$ $\frac{1}{2}$ c charm	171.2 GeV $\frac{2}{3}$ $\frac{1}{2}$ t top	
	4.8 MeV $-\frac{1}{3}$ $\frac{1}{2}$ d down	104 MeV $-\frac{1}{3}$ $\frac{1}{2}$ s strange	4.2 GeV $-\frac{1}{3}$ $\frac{1}{2}$ b bottom	
	<2.2 eV 0 $\frac{1}{2}$ ν_e electron neutrino	<0.17 MeV 0 $\frac{1}{2}$ ν_μ muon neutrino	<15.5 MeV 0 $\frac{1}{2}$ ν_τ tau neutrino	
Leptons	0.511 MeV -1 $\frac{1}{2}$ e electron	105.7 MeV -1 $\frac{1}{2}$ μ muon	1.777 GeV -1 $\frac{1}{2}$ τ tau	

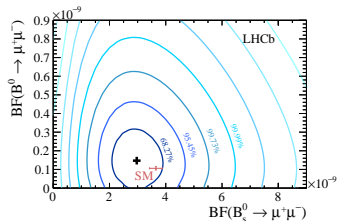
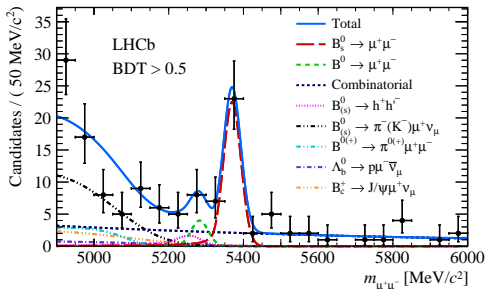
- Will take about flavour-changing neutral current decays of b quarks today.
- And only consider electroweak interactions (no gluonic penguins).
- Decays are strongly suppressed, but heavy new particles (beyond the SM) can appear in the loop and alter the final state distributions.
- "Rare B decays": $\mathcal{B} \sim 10^{-6}$

B

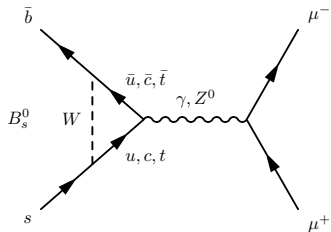
$$B_s^0 \rightarrow \mu^+ \mu^-$$



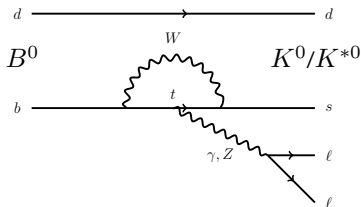
- Strongly helicity suppressed.
- Only leptons in the final state, very clean theoretically.
- $\mathcal{B}(B_s^0 \rightarrow \mu^+ \mu^-) = (3.0 \pm 0.6_{-0.2}^{+0.3}) \cdot 10^{-9}$
- $\mathcal{B}(B^0 \rightarrow \mu^+ \mu^-) < 3.4 \cdot 10^{-10}$
- All compatible with standard model predictions.



$b \rightarrow s \ell^+ \ell^-$



$$m^2(\gamma, Z) = q^2 = m^2(B_s^0)$$



$$m^2(\gamma, Z) = q^2 > 4m_\ell^2$$

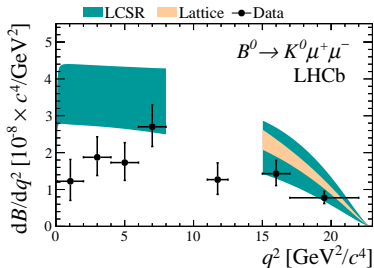
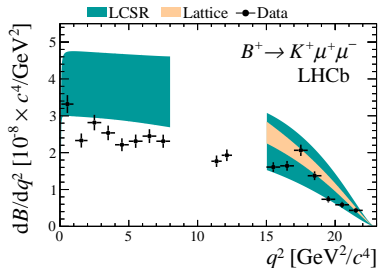
i.e. physics depends on q^2



$$\frac{d\mathcal{B}}{dq^2}$$

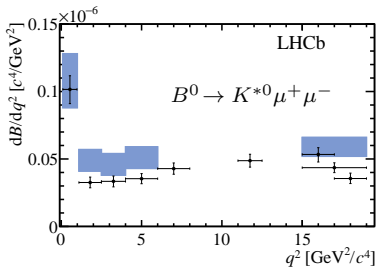
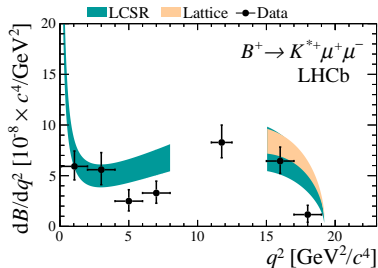
All results with 3 fb^{-1} (Run I)

$b \rightarrow sl^+l^-$ differential branching fractions



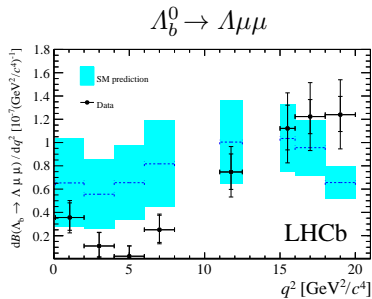
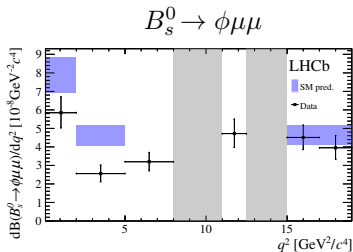
- Measure branching fractions as a function of q^2 .
 - Normalize $B \rightarrow K \mu^+ \mu^-$ to $B \rightarrow K J/\psi$
- For $B^0 \rightarrow K^0 \mu^+ \mu^-$ use decay $K_S^0 \rightarrow \pi^+ \pi^-$
- Measured values significantly below prediction for low q^2 .

$b \rightarrow sl^+l^-$ differential branching fractions



- Measure branching fractions as a function of q^2
 - Normalize $B \rightarrow K^* \mu^+ \mu^-$ to $B \rightarrow K^* J/\psi$
- For $B^+ \rightarrow K^{*+} \mu^+ \mu^-$ use decay $K^{*+} \rightarrow K_S^0 \pi^+$
- Measured values below prediction for low q^2 .

$b \rightarrow s l^+ l^-$ differential branching fractions

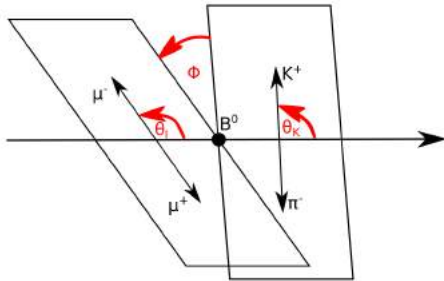


- Measure branching fractions as a function of q^2
 - Normalize $B_s^0 \rightarrow \phi \mu^+ \mu^-$ to $B_s^0 \rightarrow \phi J/\psi$
 - Normalize $\Lambda_b^0 \rightarrow \Lambda \mu^+ \mu^-$ to $\Lambda_b^0 \rightarrow \Lambda J/\psi$
- Basically all differential branching fractions are lower than their prediction for low values of q^2 .

$$\frac{d\mathcal{B}}{dq^2 d\Omega}$$

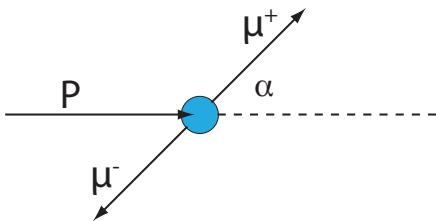
All results with 3 fb^{-1} (Run I)

Angular analysis of $B^0 \rightarrow K^{*0} \mu^+ \mu^-$



- $B^0 \rightarrow K^{*0} \mu^+ \mu^-$ has a very rich angular structure.
- Have 4 observables: $\theta_\ell, \theta_K, \phi$ and q^2
- Strategy for all following measurements: Measure angular distribution in intervals of q^2 .
- $$\frac{1}{\Gamma} \frac{d\Gamma}{d \cos \theta_\ell d \cos \theta_K d \phi} = \frac{9}{32\pi} \sum_i J_i(q^2) f(\cos \theta_\ell, \cos \theta_K, \phi)$$

A toy angular analysis

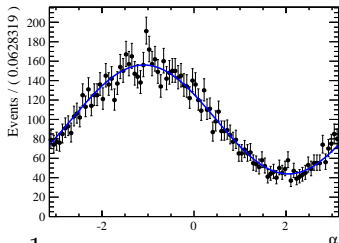


- A particle decays into two particles, with angle α .
- Suppose we can formulate the angular distribution as:

$$\frac{d\Gamma}{d\alpha} = \frac{1}{2\pi} [A \cos \alpha + B \sin \alpha + C] \quad \alpha \in [-\pi, \pi]$$

- The angular terms are given by kinematics / spin only.
- Remember: $\frac{d\sigma}{d\Omega}(e^+e^- \rightarrow \mu^+\mu^-) = \frac{\alpha^2}{4s} (1 + \cos^2 \theta)$

A toy angular analysis



$$\frac{d\Gamma}{d\alpha} = \frac{1}{2\pi} [A \cos \alpha + B \sin \alpha + C] \quad \alpha \in [-\pi, \pi]$$

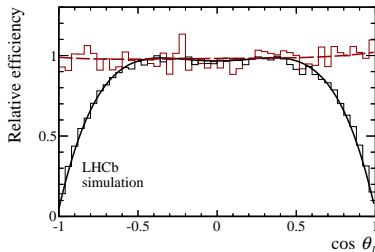
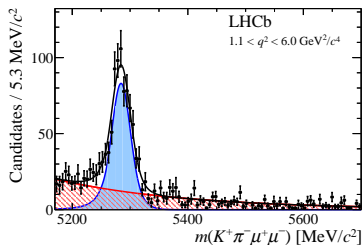
- The coefficients contain the physics-information we are interested in.
- Do: Run an experiment, collect data, select your decay, plot number of events as a function of α .
- Fit the angular distribution in collision data with the pdf and extract the coefficients.

Angular analysis of $B^0 \rightarrow K^{*0} \mu^+ \mu^-$

$$\frac{d^4(\Gamma + \bar{\Gamma})}{d \cos \theta_\ell d \cos \theta_K d\phi dq^2} = \frac{9}{32\pi} \left[\frac{3}{4}(1 - F_L) \sin^2 \theta_K + F_L \cos^2 \theta_K + \right. \\ \left. \frac{1}{4}(1 - F_L) \sin^2 \theta_K \cos 2\theta_\ell - F_L \cos^2 \theta_K \cos 2\theta_\ell + \right. \\ S_3 \sin^2 \theta_K \sin^2 \theta_\ell \cos 2\phi + S_4 \sin 2\theta_K \sin 2\theta_\ell \cos \phi + \\ S_5 \sin 2\theta_K \sin \theta_\ell \cos \phi + S_6 \sin^2 \theta_K \cos \theta_\ell + \\ S_7 \sin 2\theta_K \sin \theta_\ell \sin \phi + \\ \left. S_8 \sin 2\theta_K \sin 2\theta_\ell \sin \phi + S_9 \sin^2 \theta_K \sin^2 \theta_\ell \sin 2\phi \right]$$

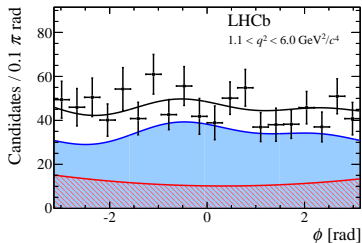
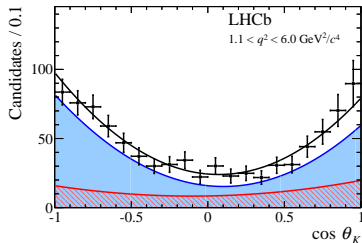
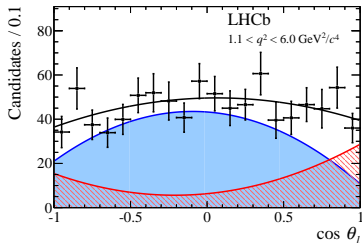
- Call the coefficient in front of the angular expressions "observable".
- Angular terms are (almost all) orthogonal.
- $S_i = f(A_0^{L,R}, A_\perp^{L,R}, A_\parallel^{L,R})$
- $S_6 = \frac{4}{3} A_{FB} = \frac{4}{3} \frac{\#\cos\theta_\ell > 0 - \#\cos\theta_\ell < 0}{\#\cos\theta_\ell > 0 + \#\cos\theta_\ell < 0}$: Forward-backward asymmetry of the leptons.

Angular analysis of $B^0 \rightarrow K^{*0} \mu^+ \mu^-$

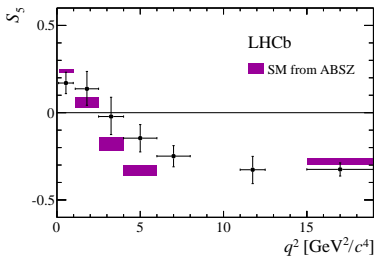
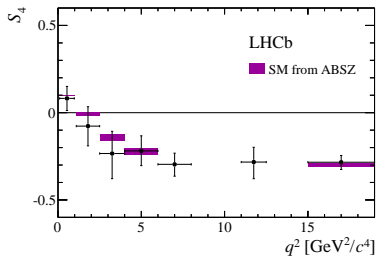
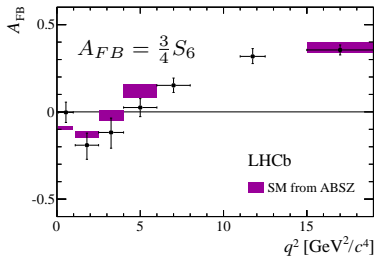
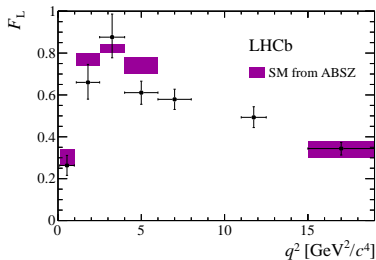


- Use a BDT to select the events, in total ≈ 2400 signal candidates.
- Use simulated sample of phase-space generated events to determine the effect of the acceptance and selection.

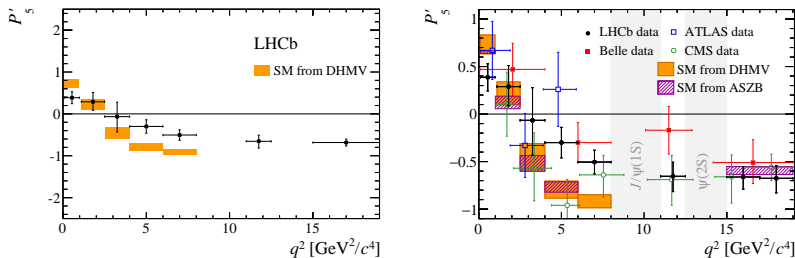
Angular analysis of $B^0 \rightarrow K^{*0} \mu^+ \mu^-$



Angular analysis of $B^0 \rightarrow K^{*0} \mu^+ \mu^-$

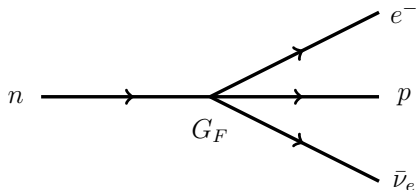


Angular analysis of $B^0 \rightarrow K^{*0} \mu^+ \mu^-$



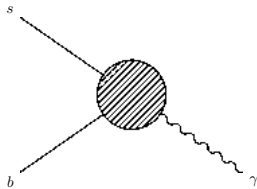
- $P'_5 = \frac{S_5}{\sqrt{1-F_L}}$
- The P'_i observables are less prone to hadronic form-factor uncertainties than the S_i ones (when using so-called "soft form-factors").
- Measurement also performed by Belle, CMS, ATLAS.
- Global significance is about 3.4σ from the SM (LHCb measurement alone).

Wilson coefficients

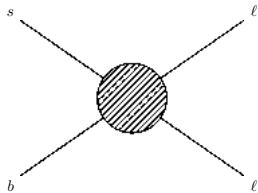


- Need a framework to describe all these different types of processes, with as little assumptions as possible.
- Fermi solved this problem already 85 years ago for the β decay by introducing a point interaction.
- G_F is a coupling constant that gives the strength of the interaction, as long as $E \ll m_W$.

Wilson coefficients



\mathcal{O}_7

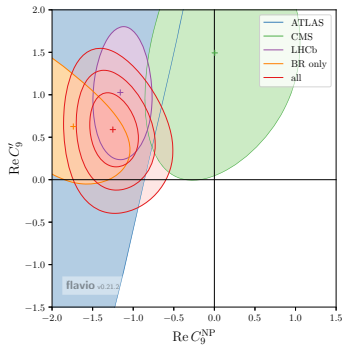
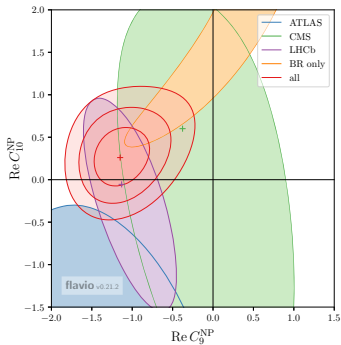


$\mathcal{O}_{9V}, \mathcal{O}_{10A}$

$$\mathcal{H}_{eff} = -4 \frac{G_F}{\sqrt{2}} V_{tb} V_{ts}^* \sum_i \mathcal{C}_i \mathcal{O}_i$$

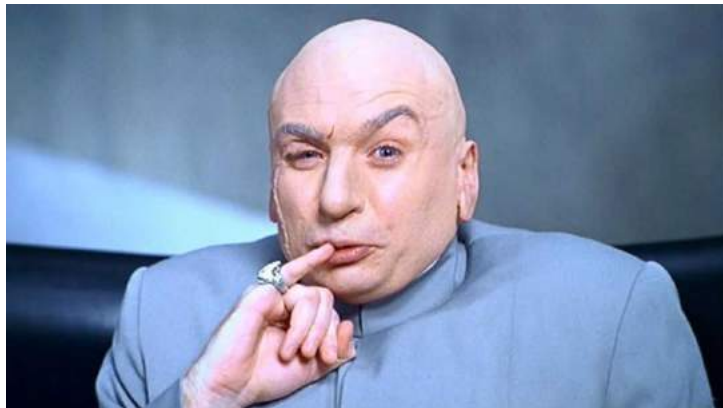
- Do the same thing for the 3 possible interactions in $b \rightarrow s \ell^+ \ell^-$ processes.
- G_F is Fermi constant, V_{tb} , V_{ts}^* CKM elements.
- \mathcal{C}_i are called Wilson coefficients, they are (complex) numbers.
- Derive \mathcal{C}_i from all measurements and combine them in global fits.

Global fits (part I)

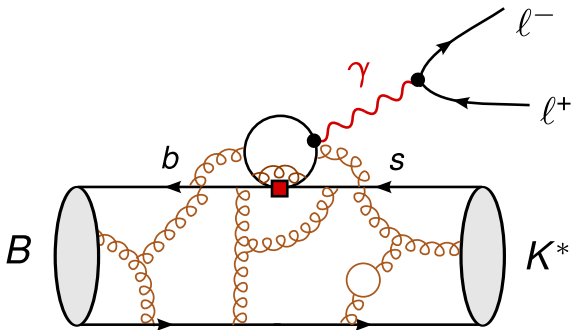


- Only consider C_9 and C_{10} .
- 0.0 / 0.0 is the Standard Model.
- Does not include lepton universality measurements (see later).
- $> 3\sigma$ away from Standard Model. Is it new physics?

The villain



The villain

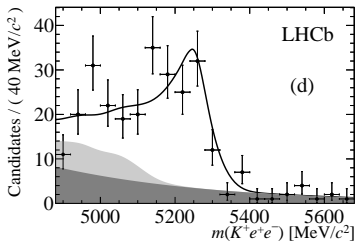
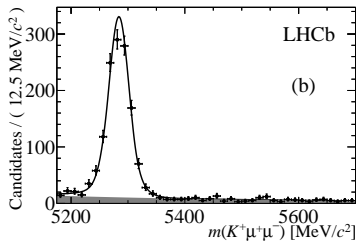


- This could mimic a new physics effect in C_9 , and is not included in the uncertainties of the hadronic form-factors.
- One could measure effect of charm-loops by a precise analysis of the $\mu\mu$ invariant mass (\rightarrow backup).

LFU

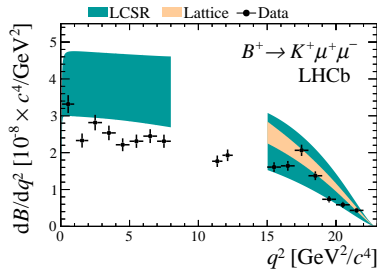
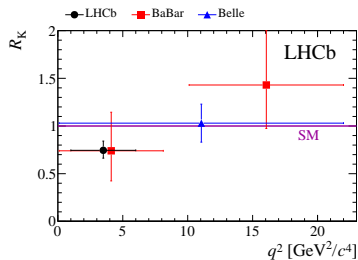
All results with 3 fb^{-1} (Run I)

Lepton Flavour universality in loop decays



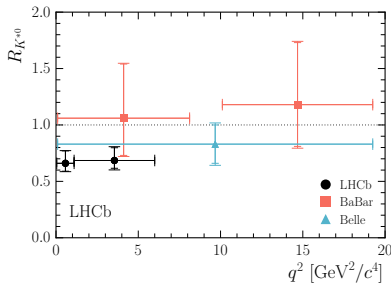
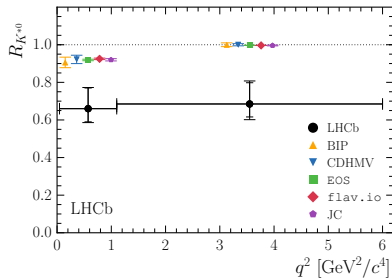
- Consider $R_K = \frac{\mathcal{B}(B^+ \rightarrow K^+ \mu^+ \mu^-)}{\mathcal{B}(B^+ \rightarrow K^+ e^+ e^-)} \stackrel{\text{theo}}{\approx} 1$
- All hadronic uncertainties cancel in the ratio, a very clean measurement.
- $B^+ \rightarrow K^+ \mu^+ \mu^-$ clean and high statistics,
 $B^+ \rightarrow K^+ e^+ e^-$ about $5 \times$ less events due to lower trigger efficiency.
- Bremsstrahlung worsens resolution.
- Only perform measurement in $1 \text{ GeV}^2/c^4 < q^2 < 6 \text{ GeV}^2/c^4$

Lepton Flavour universality in loop decays



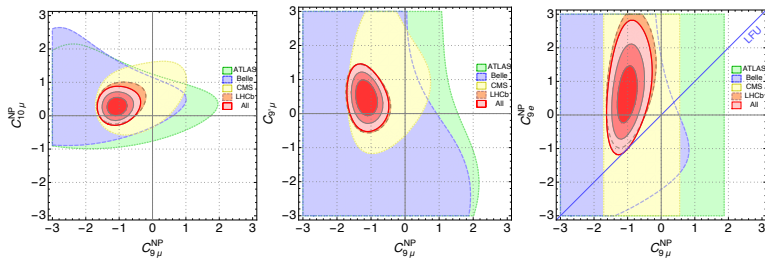
- 2.6σ from the standard model prediction.
- $\mathcal{B}(B^+ \rightarrow K^+ e^+ e^-)$ alone is measured to be compatible with the SM.
- Hm...

Lepton Flavour universality in loop decays



- Consider $R_{K^*} = \frac{\mathcal{B}(B^0 \rightarrow K^{*0} \mu^+ \mu^-)}{\mathcal{B}(B^0 \rightarrow K^{*0} e^+ e^-)} \stackrel{\text{theo}}{\approx} 1$
- All hadronic uncertainties cancel in the ratio, a very clean measurement.
- $B^0 \rightarrow K^{*0} e^+ e^-$ suffers from bremsstrahlung.
- 2.1 - 2.5 σ from the standard model prediction.

Global fits (part II)



- Can introduce a different Wilson coefficient C_9 for muons and electrons and redo global fits.
- Compatible with only deviations in the muon channels and not in the electron channels.

Models

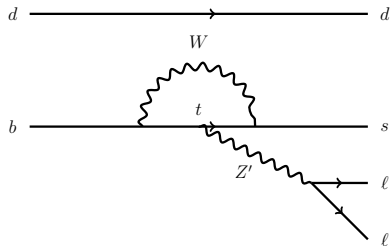
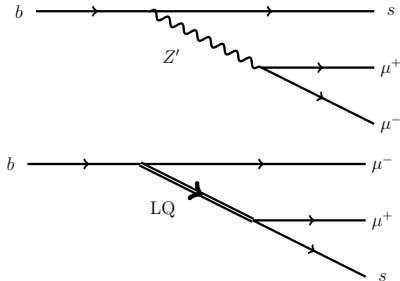
- There are a plethora of models that try to explain all anomalies simultaneously.
- Have to explain: Difference in branching fractions, angular distributions, lepton flavour on tree* and loop level.
- Possible masses of BSM particles:

tree level, unsuppressed $\sim 30 \text{ TeV}$	loop level, unsuppressed $\sim 2.5 \text{ TeV}$
tree level, MFV $\sim 6 \text{ TeV}$	loop level, MFV $\sim 0.5 \text{ TeV}$

* $R(D^*)$, see backup

MFV = Minimal flavour violation, i.e. CKM structure holds BSM

Z' and Leptoquarks



- Can introduce a Z' that causes a flavour-changing neutral current on tree level or loop level.
- Leptoquarks can simultaneously explain $R_{D^{*0}}$ (tree-level leptoquarks), R_K (loop-level leptoquarks) and muon $g - 2$ (PRL116, 141802 (2016))

Optimist's point of view



- Many measurements show a deviation, and when combined, it is significant.
- The pattern is somewhat consistent, as shown by global fits.
- The effects are observed by several (independent) measurements and experiments.
- No large uncertainty in the theoretical prediction has been discovered.

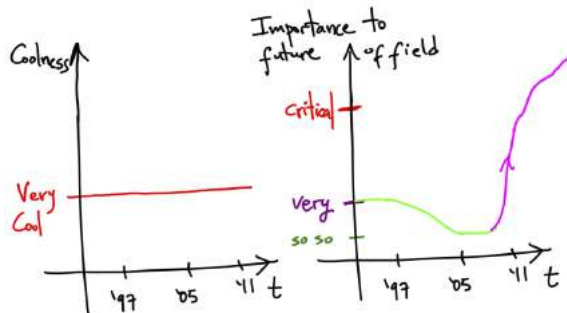


Pessimist's point of view



- "The effect is of a magnitude that remains close to the limit of detectability, or many measurements are necessary because of the very low statistical significance of the results."
I. Langmuir on pathological science.

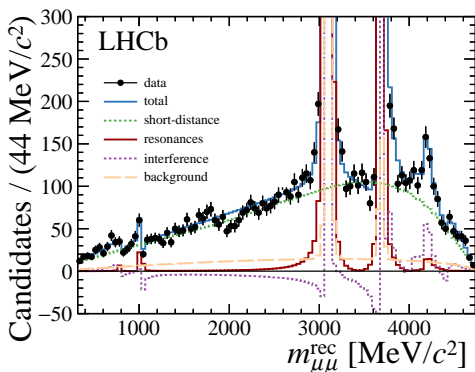
Conclusion



- Rare B decays are an exciting field of research.
- Several intriguing deviations from the SM have shown up in flavour-changing neutral currents.
- The combination could hint to a deviation from the Standard Model.
- Nature of these anomalies will hopefully soon be resolved.

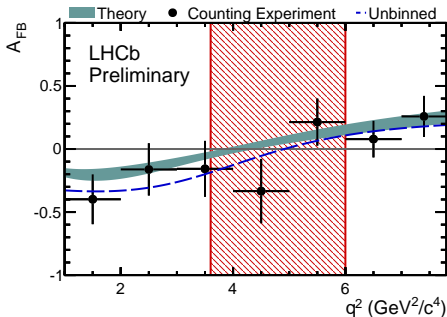
Backup

Short- and long-distance effects in $B \rightarrow K \mu^+ \mu^-$



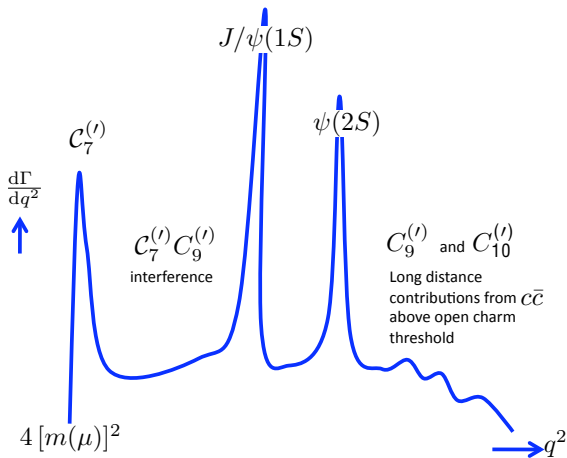
- Need to perform same analysis for $B^0 \rightarrow K^{*0} \mu^+ \mu^-$ to understand effect of charm-loops.

Zero-crossing point in $B^0 \rightarrow K^{*0} \mu^+ \mu^-$

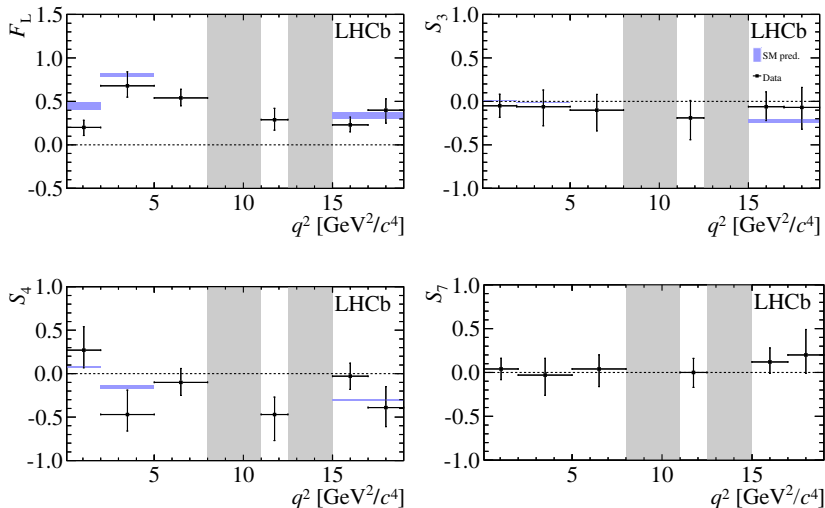


- Zero-crossing point of A_{FB} theoretically clean.
- For example: $S_6 = \frac{4}{3} A_{FB} = \frac{4}{3} \frac{\#\cos\theta_\ell > 0 - \#\cos\theta_\ell < 0}{\#\cos\theta_\ell > 0 + \#\cos\theta_\ell < 0}$.
Forward-backward asymmetry of the leptons.

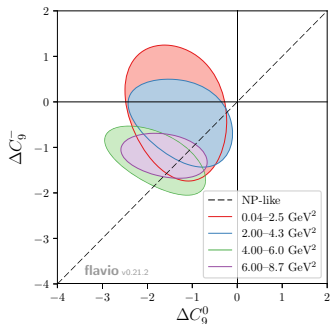
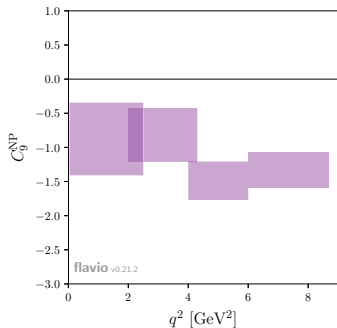
q^2 spectrum for $B^0 \rightarrow K^{*0} \mu^+ \mu^-$



Angular analysis of $B_s^0 \rightarrow \phi \mu^+ \mu^-$

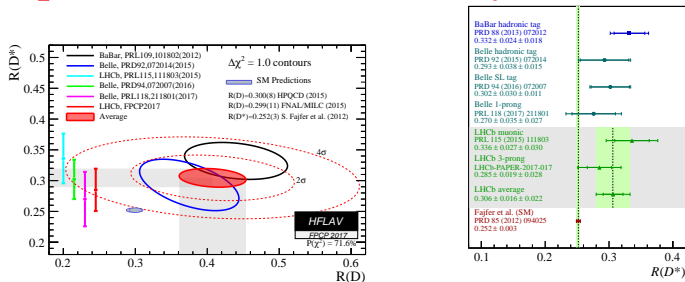


q^2 dependence of C_9



- Check shift of C_9 as a function of q^2 . Should be constant (assuming the new physics is heavy enough).
- That's a hint, but not a confirmation for non-hadronic BSM effects.

Lepton Flavour universality in tree decays



- Can also measure lepton flavour universality in semileptonic (tree) decays.
- Measure $\frac{\mathcal{B}(\bar{B}^0 \rightarrow D^{*+} \tau^- \nu)}{\mathcal{B}(\bar{B}^0 \rightarrow D^{*+} \mu^- \nu)}$
- Either with $\tau^- \rightarrow \mu^- \nu \nu$ (PRL 115, 111803 (2015))
- Or with $\tau^- \rightarrow \pi^- \pi^+ \pi^- \nu$ (arxiv:1708.08556)
- LHCb measurements consistent with B factories, combination about 4σ from the SM.

