

Photons to search for new phenomena with the ATLAS detector

with an emphasis on the
[ATLAS-CONF-2022-018](#)

preliminary results, first presented at Moriond QCD 2022

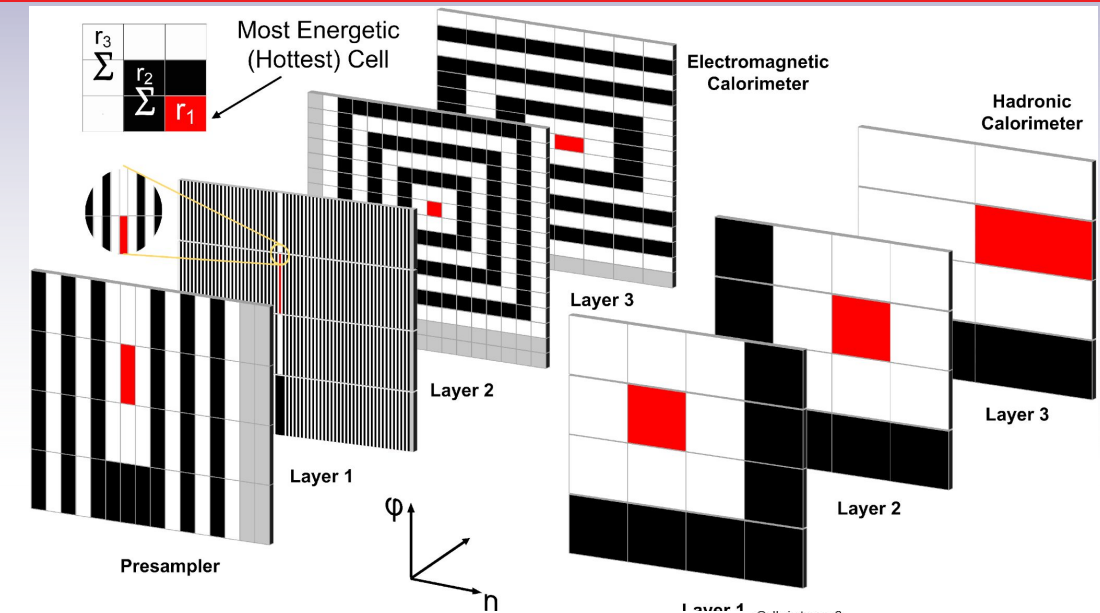
José Ocariz, LPNHE-Paris and Université Paris Cité

Sampling electromagnetic calorimeter

- measures energy of photons (and e)
- interleaves layers of lead (absorber) and liquid Argon
- longitudinally segmented into three layers
- provides an excellent photon energy resolution :

$$\frac{\sigma(E)}{E} = \frac{10\% \sqrt{GeV}}{\sqrt{E}} \oplus 0.7\%$$

crucial for the Higgs boson discovery in 2012 !

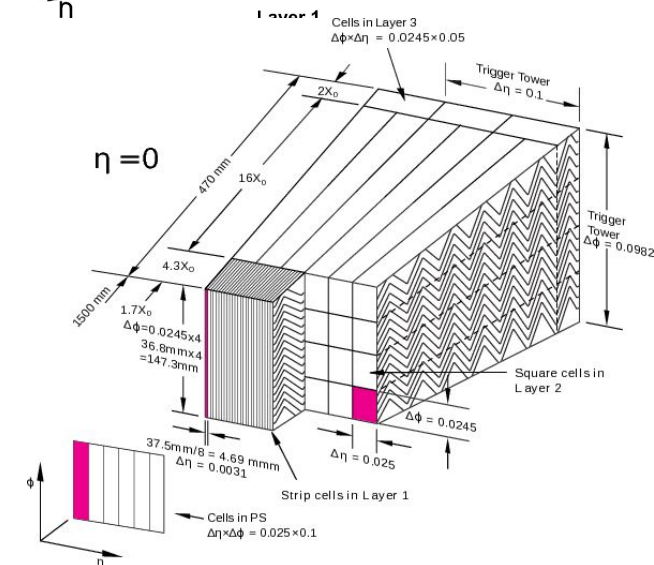


Hadronic calorimeters

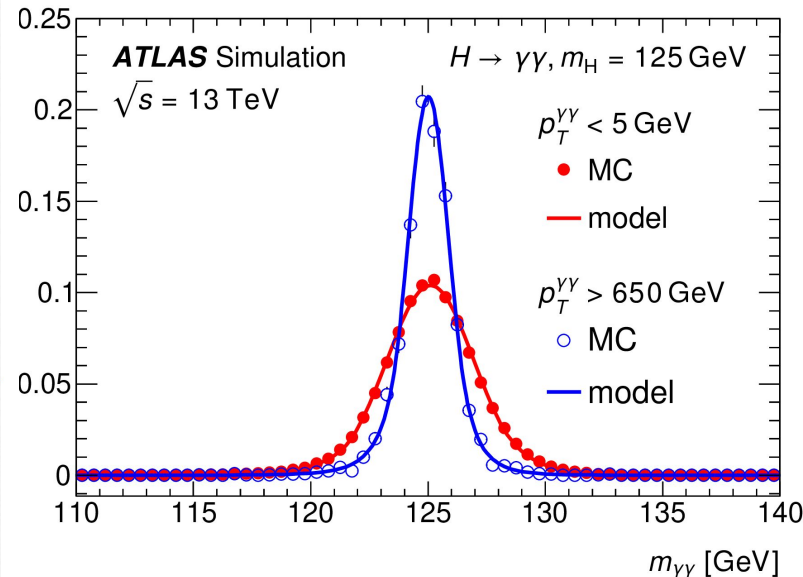
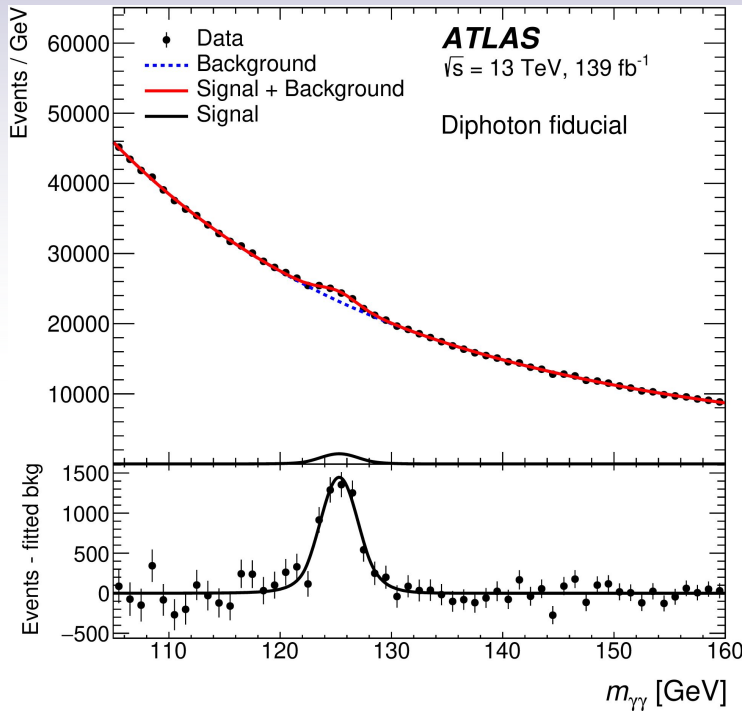
- complement energy measurements (jets and taus)

Energy deposits in calorimeter cells clustered together into “topoclusters”

- dynamical objects that span over all layers
- matching of tracks and clusters to
 - identify photons
 - discriminate from other objects in the calorimeters



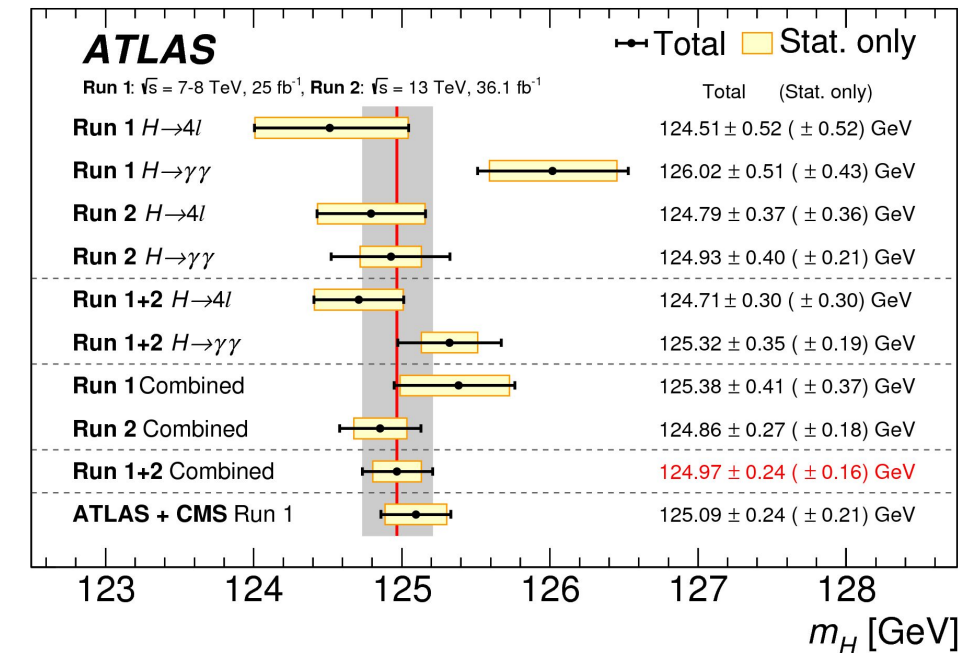
The $H \rightarrow \gamma\gamma$ channel provides several of the cleanest signatures for the study of Higgs boson properties

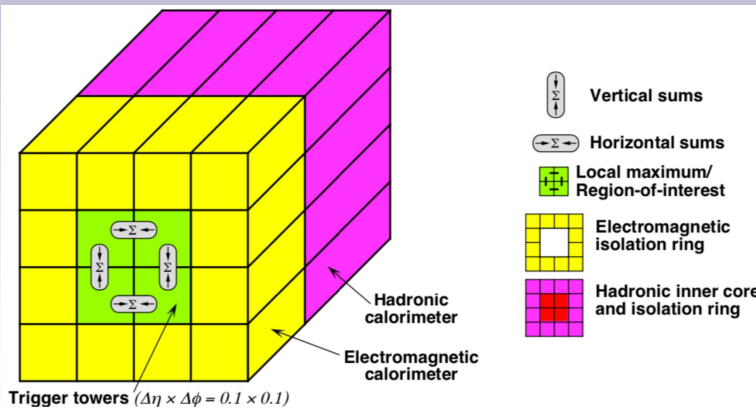


Higgs boson mass :

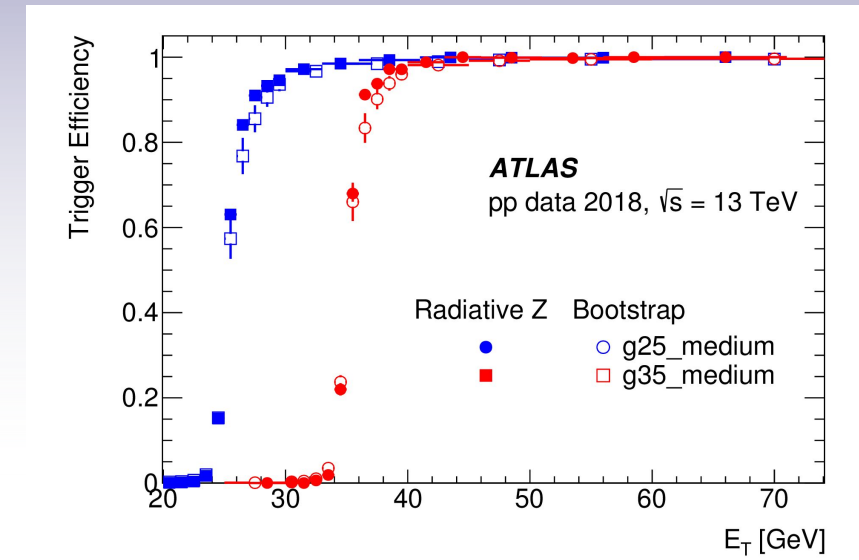
- precision measurement, at $\pm 0.17\%$ level
- $H \rightarrow \gamma\gamma$ precision comparable to $H \rightarrow 4l$ that benefits from a very clean S/B but suffers from a low signal rate

- despite large backgrounds from QCD (di)photons
- clear, narrow $m_{\gamma\gamma}$ signal
- excellent invariant mass resolution
 - typically in the order of $\pm 1.5 \text{ GeV}$ to $\pm 2 \text{ GeV}$
 - can be as low as $\pm 1 \text{ GeV}$ for the hardest photon pairs





- Two-level photon triggers based on
 - readout of towered energy sums around ROIs at L1
 - close-to-complete reconstruction at HLT
- Thresholds on transverse energy are needed to cope with overwhelming rates
 - excellent efficiency plateaus above the thresholds
- Thresholds and criteria vary with luminosity and data-taking periods
- A diphoton trigger combines two separate single-photon triggers



Year	2015	2016 up to D3	2016 from D3	2017	2018
L1 item	L1_2EM10VH	L1_2EM15VH	L1_2EM15VH	L1_2EM15VHI	L1_2EM15VHI
HLT item	2g20_tight	2g20_tight	2g22_tight	2g20_tight_icalovloose	2g20_tight_icalovloose
luminosity [fb^{-1}]	3.2	11.5	21.5	43.6	58.5

Photons (and electrons) : very similar development of electromagnetic showers in a uniform calorimeter

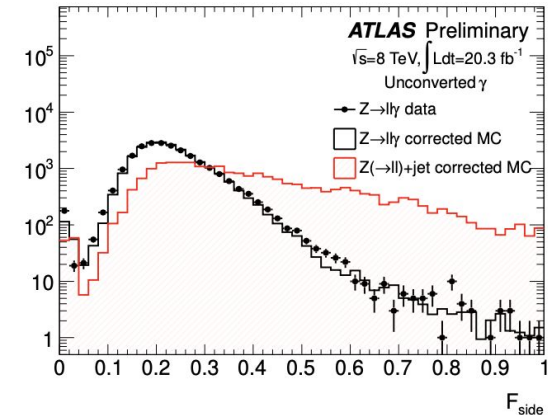
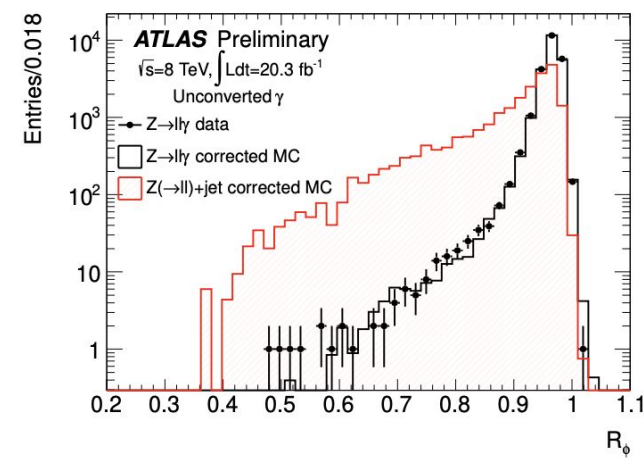
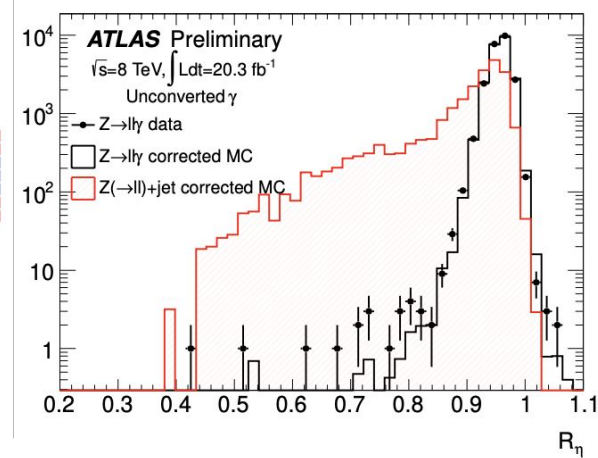
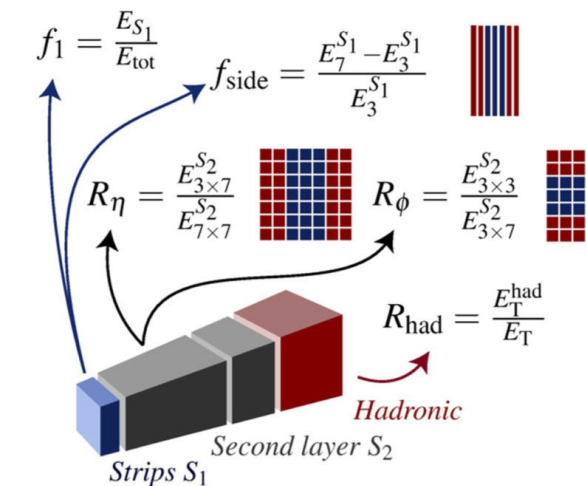
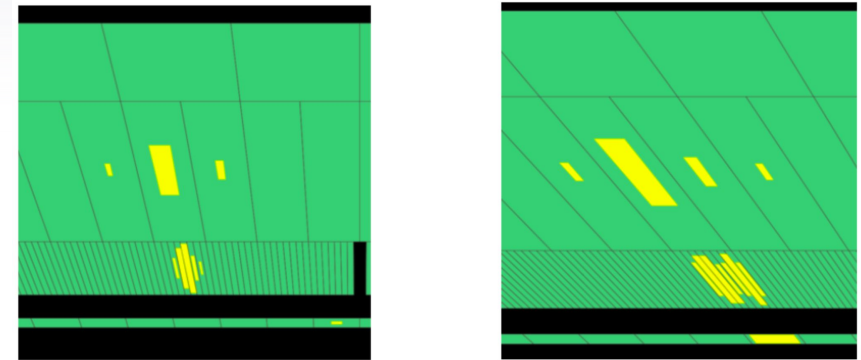
- (almost) invariant lateral and longitudinal developments

An e/gamma candidate is a “super-cluster” arrangement of topoclusters, compatible with an EM shape

- multivariate selection on a collection of correlated “shower-shape” variables
- for photons, selection is optimized to reject “fake” candidates (mostly neutral hadrons in jets)

Photon-ID:

- high-granularity of the EM calorimeter plays a crucial role

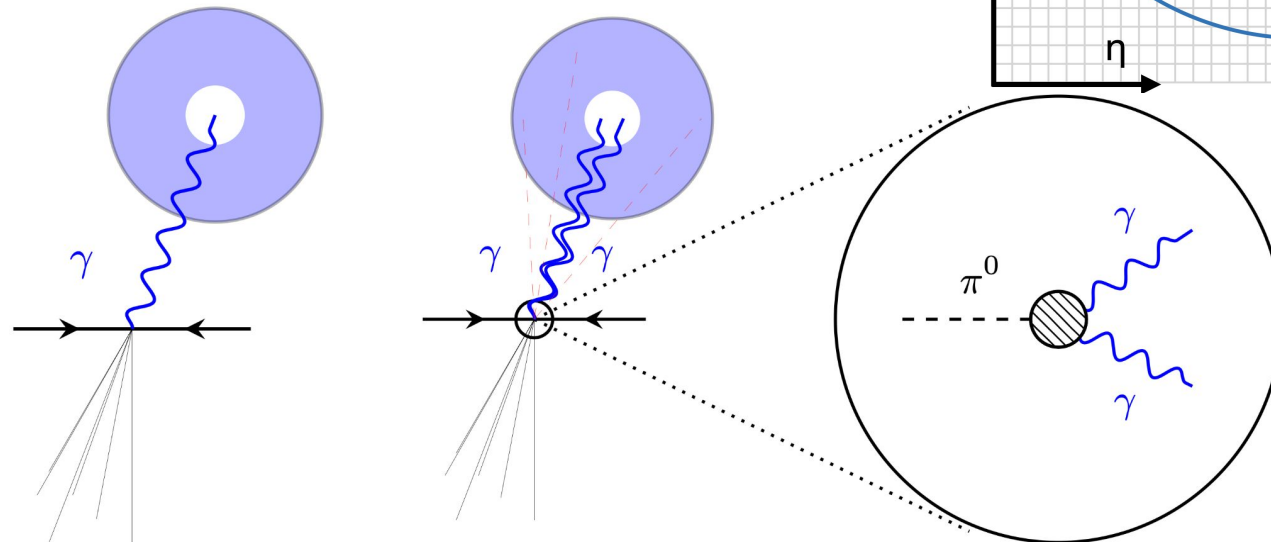
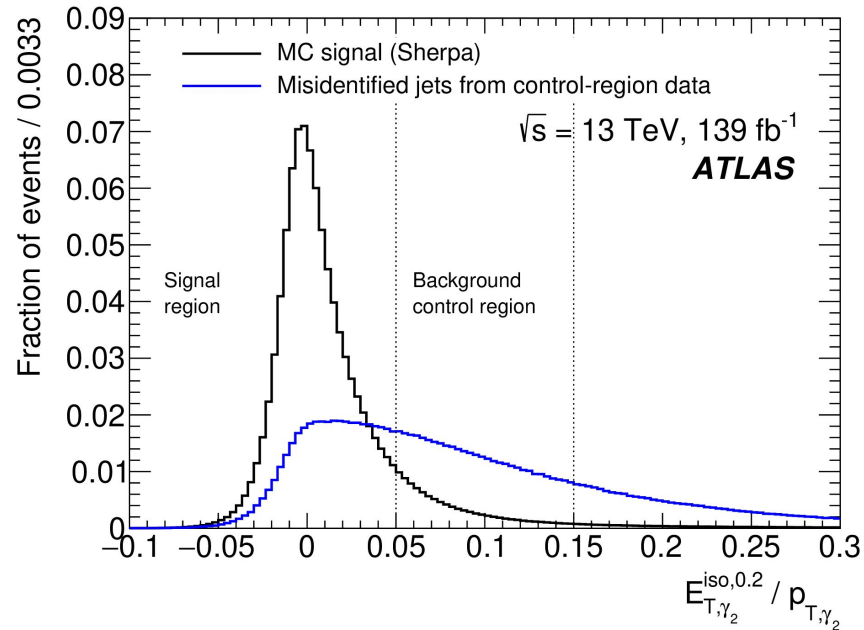
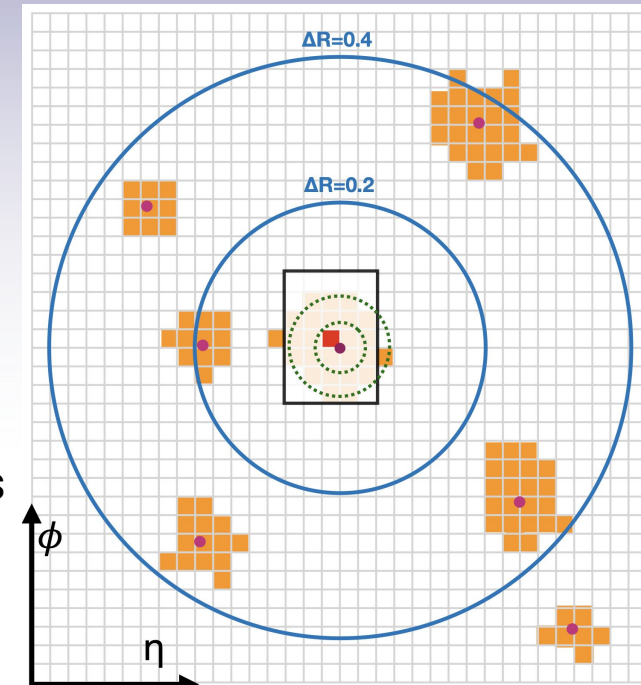


Photon isolation rejects further backgrounds on top of photon ID

- fakes are usually surrounded by additional particles
- calorimetric isolation energy :

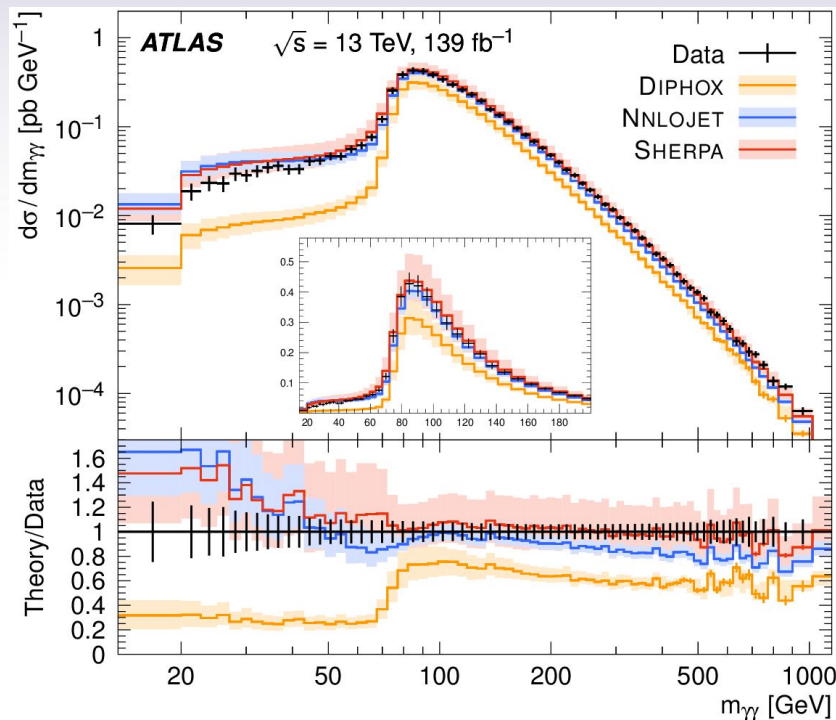
$$E_T^{\text{iso}}(R) = \left(\sum_{r_i < R}^{\text{clusters}} E_T^i \right) - E_T^{\text{corr}}(E_T^\gamma, \eta^\gamma, \langle \mu \rangle)$$

- sum of topocluster transverse energies within a ΔR cone around the candidate
 - corrections to account for various effects (pileup...)
 - the ΔR size impacts diphoton selections
- complemented with a track-based definition for charged particles around candidates

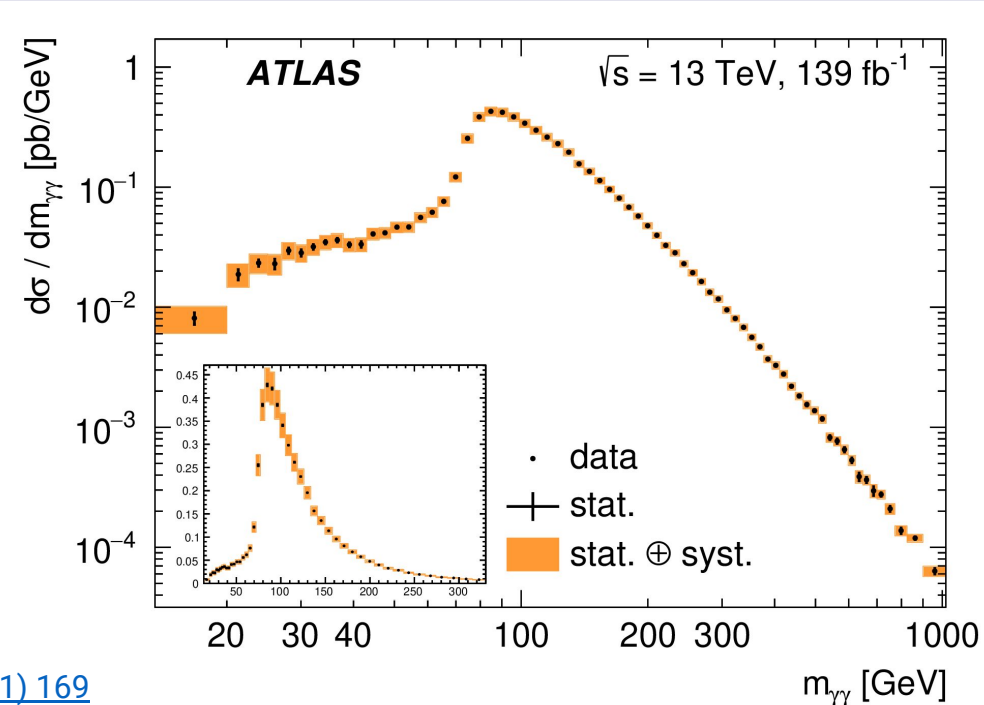


(made@ITKP) study of diphoton production in pp collisions

- differential cross-sections, as function of $m_{\gamma\gamma}$, $p_{\gamma\gamma}^T$ and angular distributions of the photon pair
- precision tests of both the SM predictions and phenomenological models



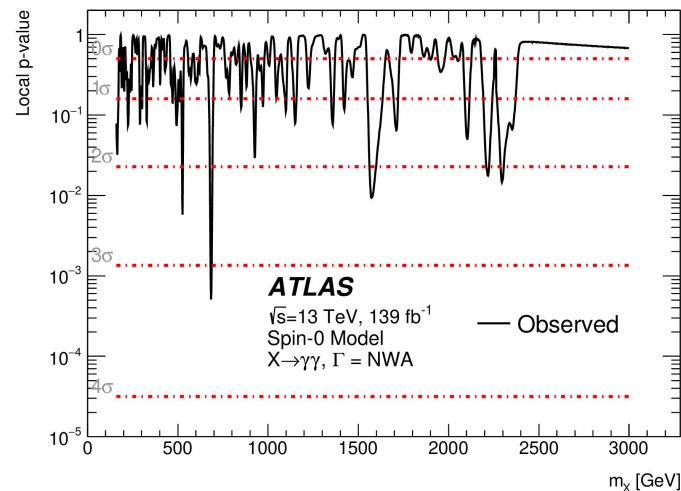
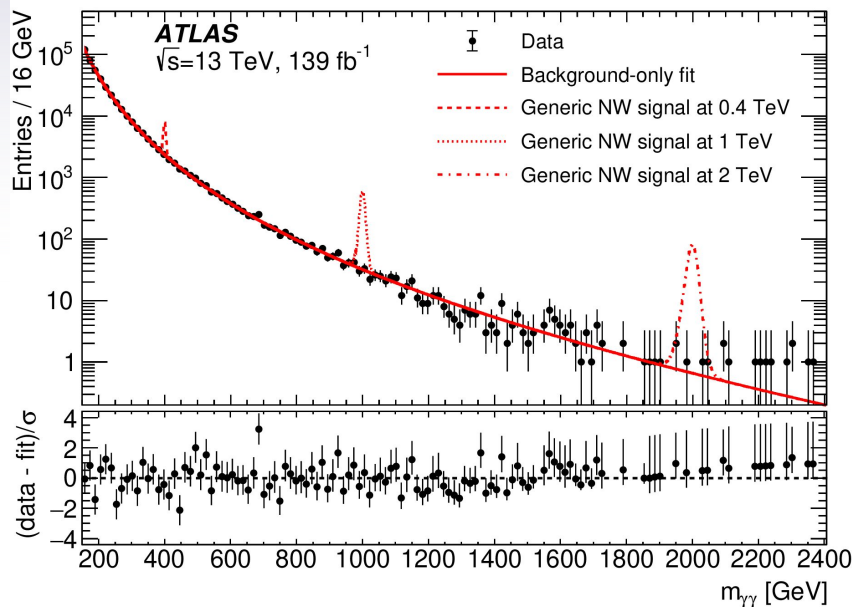
[JHEP 11 \(2021\) 169](#)



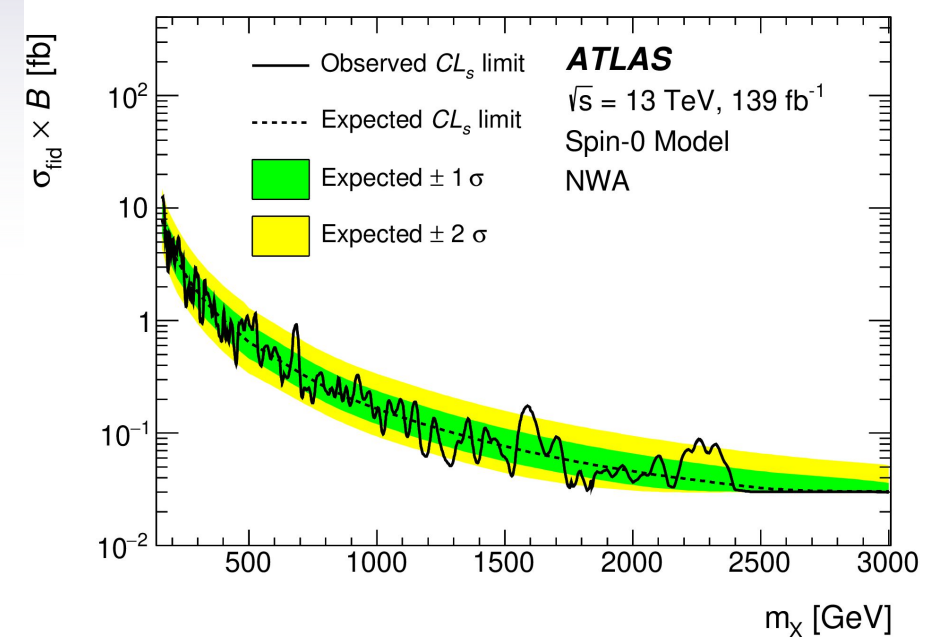
The SM predicts only one particle decaying into pair of photons in the ≥ 10 GeV mass range : the Higgs boson :-)

- many such new particles appear in a number of models beyond the Standard Model
 - among others, extended Higgs sectors lead to the prediction of additional neutral Higgses
- could there be other resonant peaks above the SM QCD diphoton spectrum ?

- ATLAS and CMS have searched for diphoton resonances, in mass ranges above the Higgs mass
 - analysis strategies very similar to the ones optimised for the 125 GeV Higgs boson
 - a smooth, monotonically-decreasing background, dominated by QCD diphotons



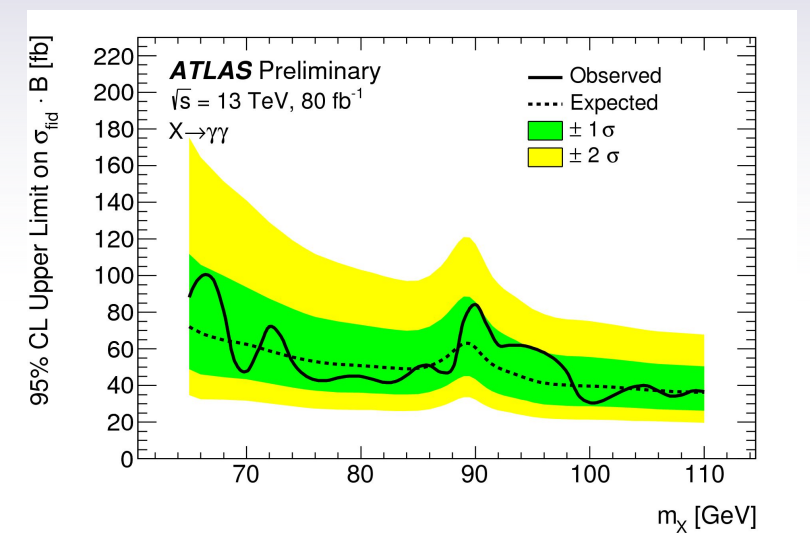
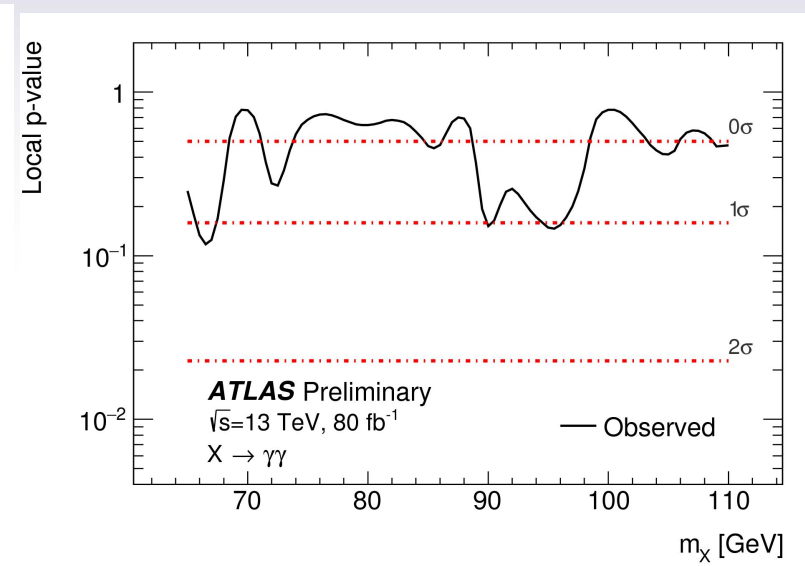
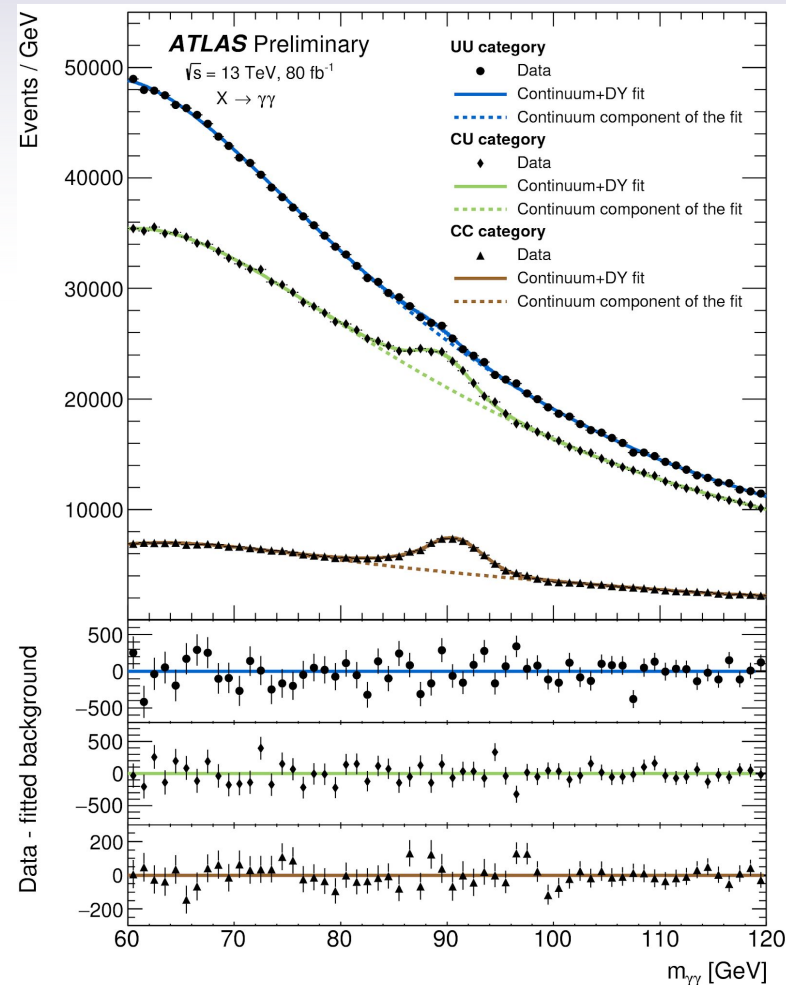
[Phys. Lett. B 822 \(2021\) 136651](#)



A large number of statistically independent mass subranges are explored

- owing to the excellent performances on photon performances, in particular energy calibration
- the rate of local fluctuations is compatible with the global expectations
 - beware of over-interpretations :-)
- results are interpreted as limits in the production and decay rate of a hypothetical, new, heavy, “X” particle
 - the exclusion limits span from 200 GeV to ~2.5 TeV

- ATLAS and CMS have also searched for diphoton resonances, in mass ranges below the Higgs mass
 - analysis strategies again very similar to the ones optimised for the 125 GeV Higgs boson
 - splitting the sample into photon conversion categories important to account for the Z background



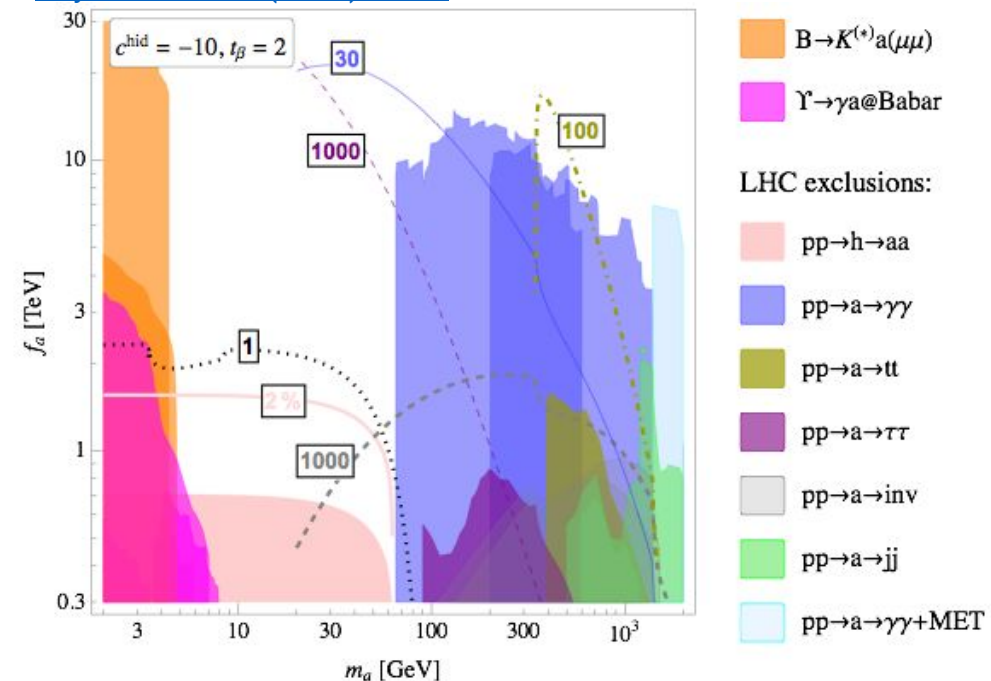
- Sensitivity strongly impacted by the difficulty in controlling the background shape
- the infamous “spurious signal” issue...
 - shape below 60 GeV dominated by the trigger “turn-on” effect

A different approach is needed to reach lower masses!

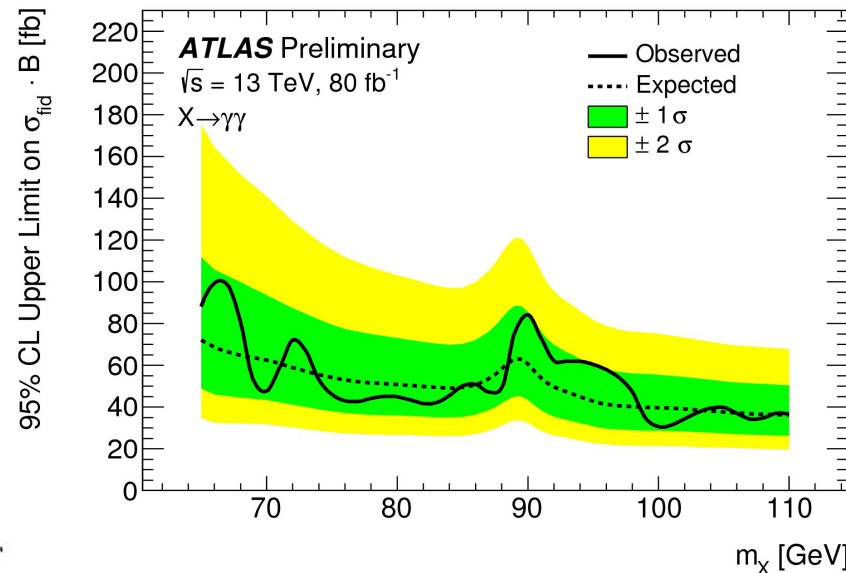
- Several proposals for resonant Axion-Like Particles within the LHC mass reach
 - pNGBs associated to a spontaneously broken approximate symmetry above the TeV scale
- main interest as a possible DM mediator due to its weakly interacting nature
- ALPs below the Higgs mass would couple predominantly to gluons and photons
 - LHC is the natural laboratory to search for them !
- existing search gap in $\gamma\gamma$ channel resonance searches
- goal: push the current 65 GeV limit towards lower masses, close the gap as much as possible !

[Phys.Rev.Lett. 119 \(2017\) 14, 141804](#)

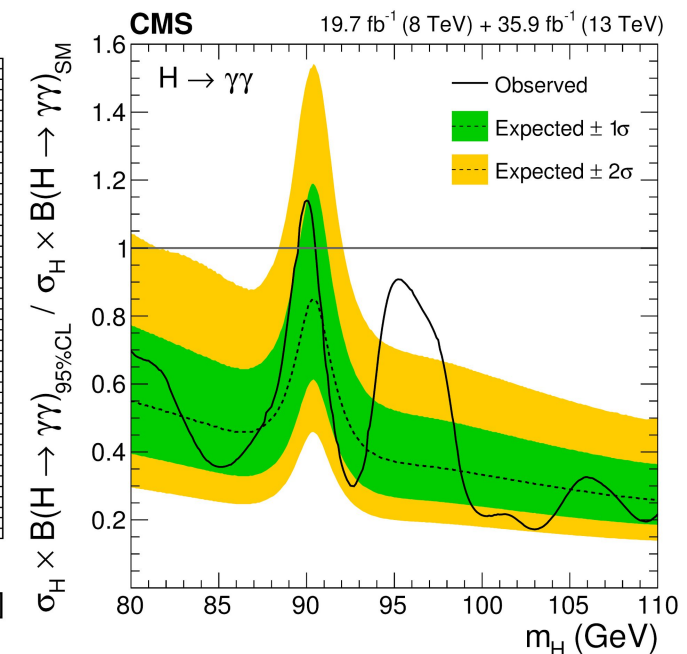
[Phys.Lett.B 783 \(2018\) 13-18](#)



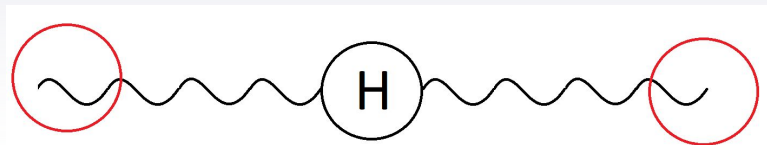
[ATLAS-CONF-2018-025](#)



[Phys.Lett.B 793 \(2019\) 320-347](#)

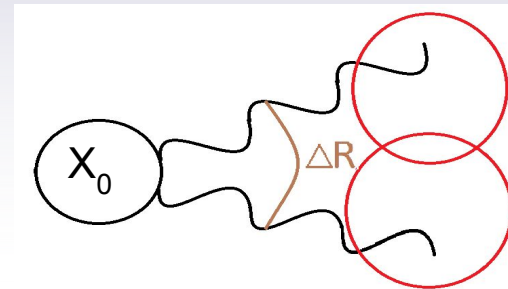


The “standard” ATLAS diphoton analysis has been mostly optimized for the study of the Higgs boson



$$m_{\gamma\gamma} = \sqrt{2E_{T,1}^{\gamma}E_{T,2}^{\gamma}[\cosh(\eta_2 - \eta_1) - \cos(\Phi_2 - \Phi_1)]}$$

- Most photons are emitted back-to-back
- Their E^T spectrum are largely within the trigger plateaus
- Standard photon isolation has very good performances
- A diphoton candidate is just the combination of two single-photon candidates



$$m_{\gamma\gamma} \approx \sqrt{p_{T,1}^{\gamma}p_{T,2}^{\gamma}\Delta R}$$

To attain small X_0 masses, two conflicting possibilities :

- reach lower E^T values
 - (but the trigger...)
- reach smaller angular separations
 - (but the isolation...)

Main limiting factors to reach diphoton masses below 65 GeV :

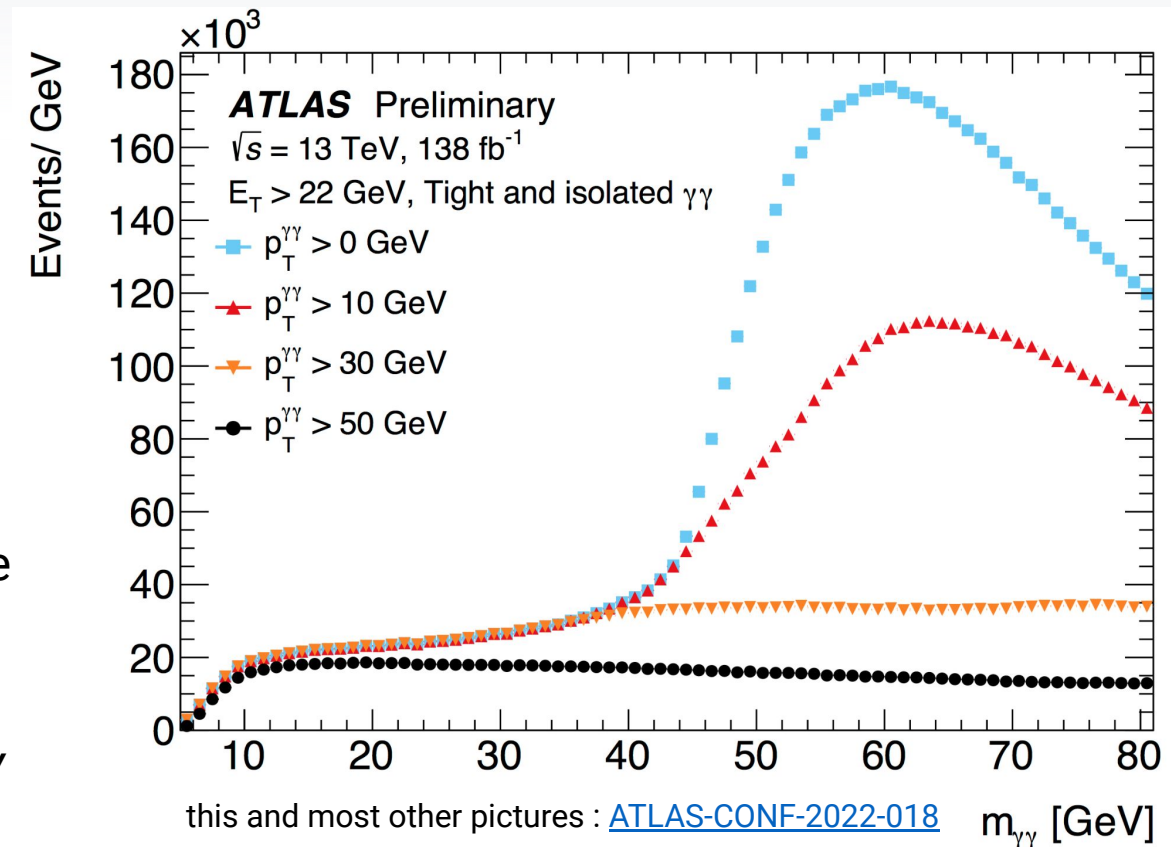
- diphoton trigger E_T thresholds at 20 GeV
 - reduces signal acceptance
 - limits background modelling with analytical functions due to steep trigger turn-on
- decrease in photon identification and isolation efficiencies for low-ET photons

This analysis follows standard ATLAS diphoton selections

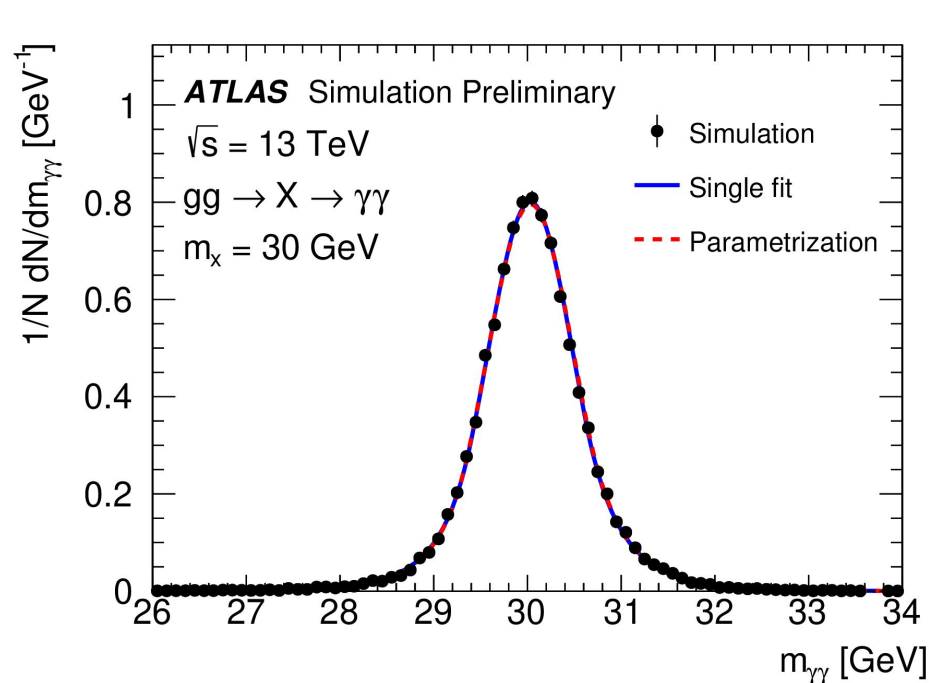
- data recorded with unrescaled diphoton triggers
- trigger thresholds and criteria evolved during Run-2
 - 20 GeV E_T thresholds for most data (except for 21.6 fb^{-1} in 2016 with a 22 GeV threshold)
 - additional trigger-level isolation criteria in 2017+2018
- two reconstructed photon candidates with $E_T > 22 \text{ GeV}$
 - within the $|\eta|$ acceptance
 - passing *tight identification* criteria
 - passing *tight isolation* criteria (calorimetric+track)
- isolation computed in a $\Delta R < 0.2$ cone around the candidate

Events at low mass have large transverse momentum $p_T^{\gamma\gamma}$

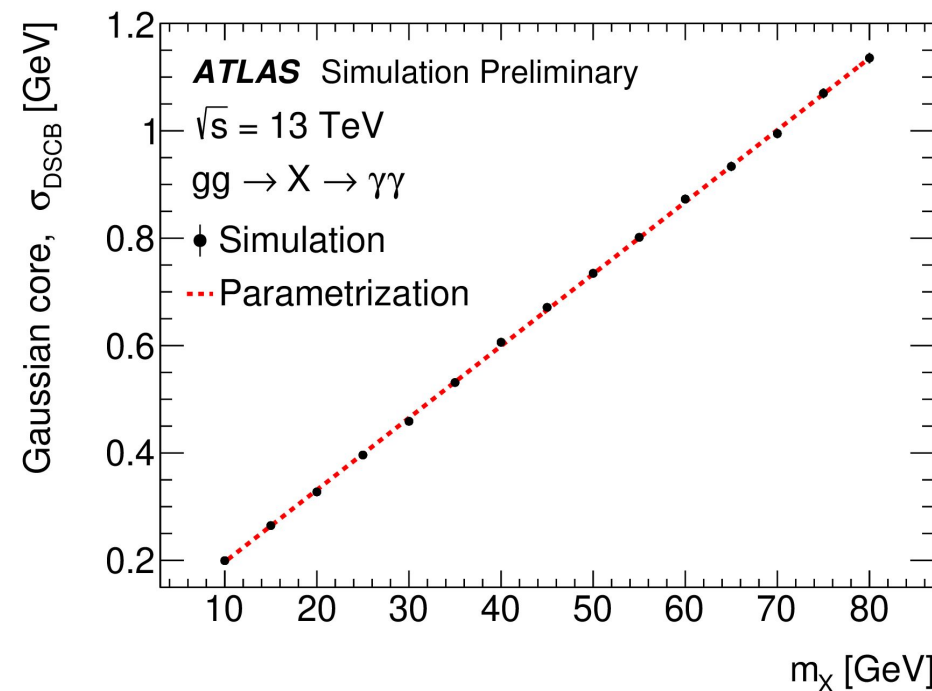
- add a boosted diphoton selection : $p_T^{\gamma\gamma} > 50 \text{ GeV}$
- results in a smooth background spectrum down to 10 GeV



- Signal MC control samples :
 - EFT framework: scalar “Higgs-like” resonance
 - gluon-fusion production only
 - generated with MadGraph at LO+0,1,2 jets
- Invariant diphoton mass resolution described with a Double Sided Crystal Ball (DSCB) function
 - narrow-width approximation (fixed $\Gamma = 4.07 \text{ MeV}$)
 - DCSB parameters are linear functions of the mass point being tested
 - biases on fitted signal yields below the $\pm 1\%$ level on the full mass range



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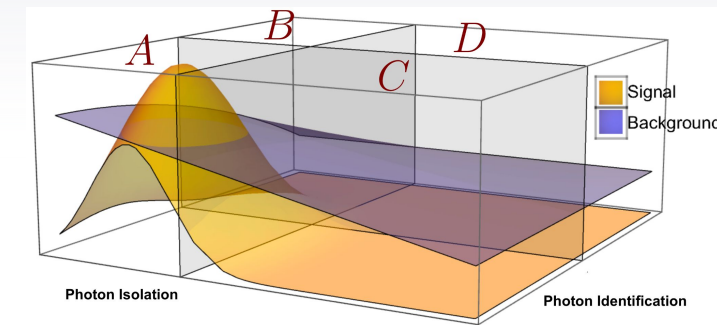
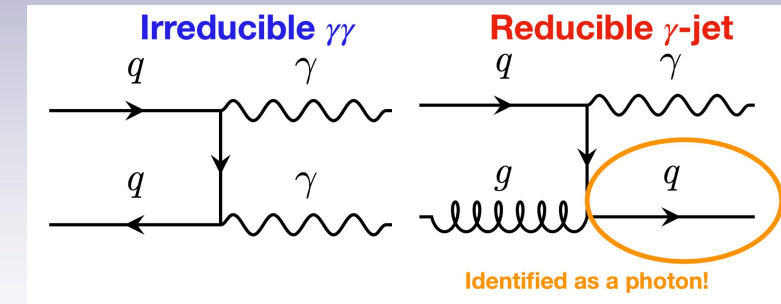


Non-resonant backgrounds:

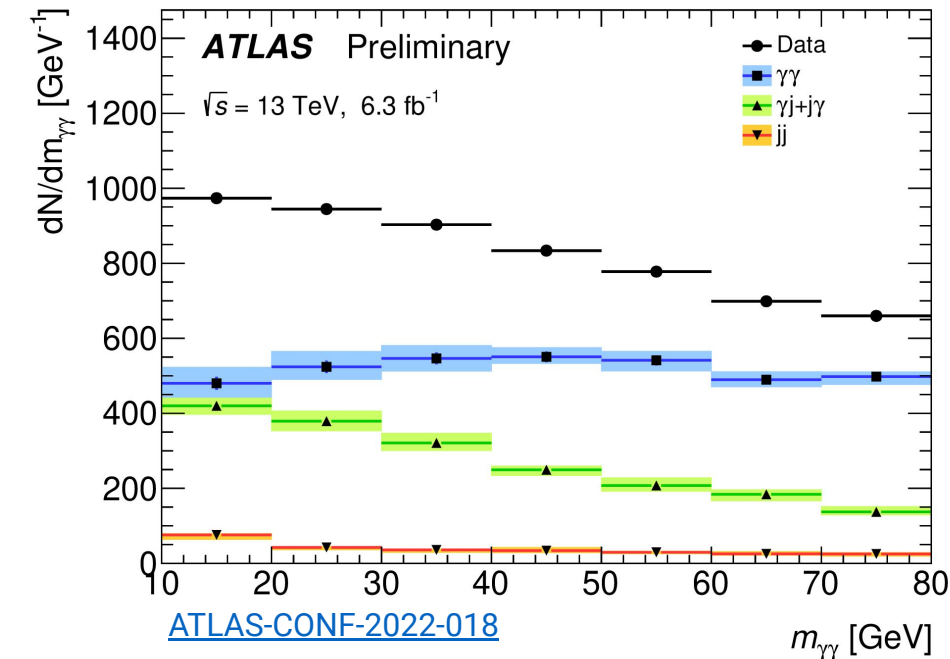
- irreducible ($\gamma\gamma$) from QCD diphoton production
- reducible ($\gamma j + j\gamma + jj$) from QCD with 1 or 2 jets misidentified as photon
- other backgrounds (i.e. from electrons) found to be negligible

Extract composition from double-ABCD method (aka 2x2D)

- using Isolation and Identification on each photon
- irreducible shapes extracted from Sherpa QCD diphoton
- reducible shapes extracted from control regions in data
 - photon candidates failing a subset of identification criteria



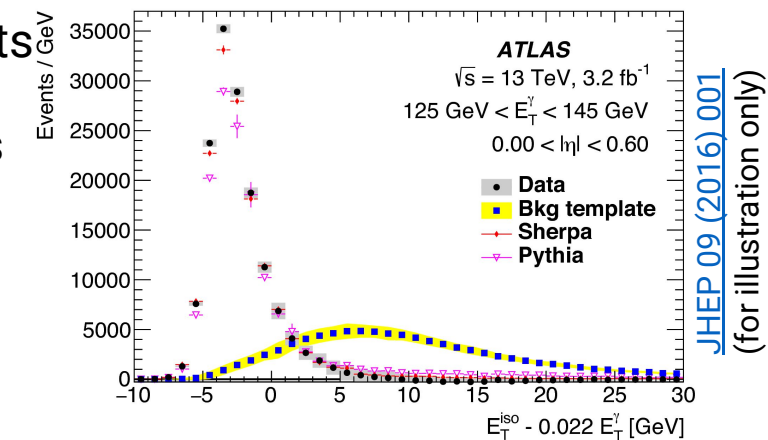
(borrowed from [Phys.Rev.D 103 \(2021\) 3, 035021](https://arxiv.org/abs/2103.03502))



2x2D composition on a control sample passing a (prescaled) diphoton trigger with looser trigger-level ID requirements

large diphoton purity at higher masses

reducible background becomes dominant at very-low masses



Background shape qualitatively divided into two regions:

- fast turn-on region for masses below ~ 20 GeV
 - described with an exponentially-saturating function (“Flat”)
- slowly decreasing region above, with a mild change in curvature between the mid- and higher- mass regions
 - described with the product of a power-law (“PowLaw”) times an “Activation” function

$$f(m_{\gamma\gamma}, \vec{\theta}) = \text{Flat} + [\text{PowLaw} \times \text{Activation}] =$$

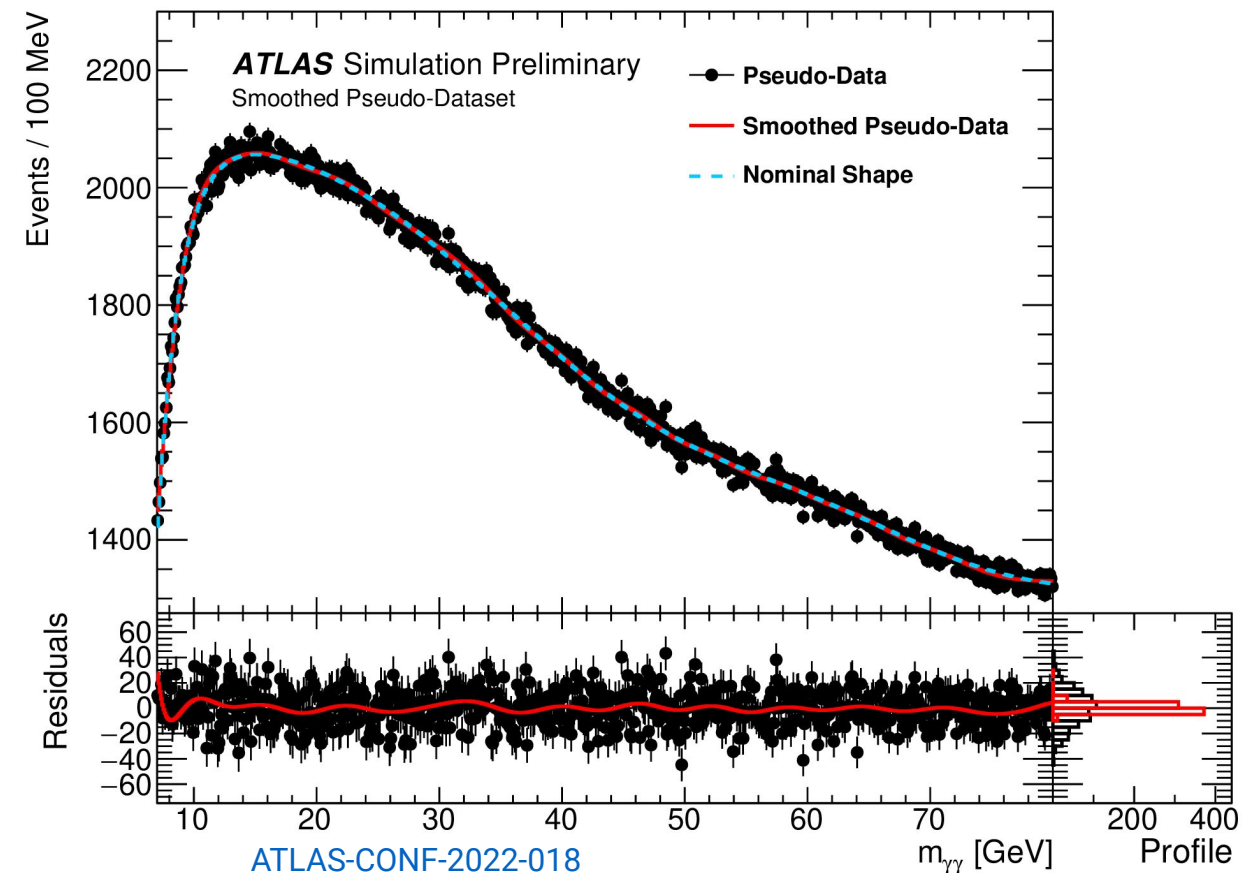
$$= \left[1 - (1 - f_0) e^{-\frac{m_{\gamma\gamma} - 10}{\tau_{\text{flat}}}} \right] \text{frac} + (1 - \text{frac}) \left[\underbrace{\left(1 - \left(\frac{m_{\gamma\gamma}}{c_1} \right)^{a_0} \right)^{c_0}}_{\text{PowLaw}} \underbrace{\left(1 + \frac{e^{\frac{m_{\gamma\gamma} - \delta_{\text{tail}}}{\tau_{\text{tail}}}}}{1 + e^{-\frac{m_{\gamma\gamma} - \delta_{\text{thresh}}}{\tau_{\text{thresh}}}}} \right)}_{\text{Activation function}} \right]$$

Variations on the template shape are used to evaluate the flexibility of the background modelling :

- different control region definitions
- different $\gamma\gamma$ purity
- varied $p_T^{\gamma\gamma}$ threshold

All variations lead to sizable changes in the spectrum

- all well described with the nominal model



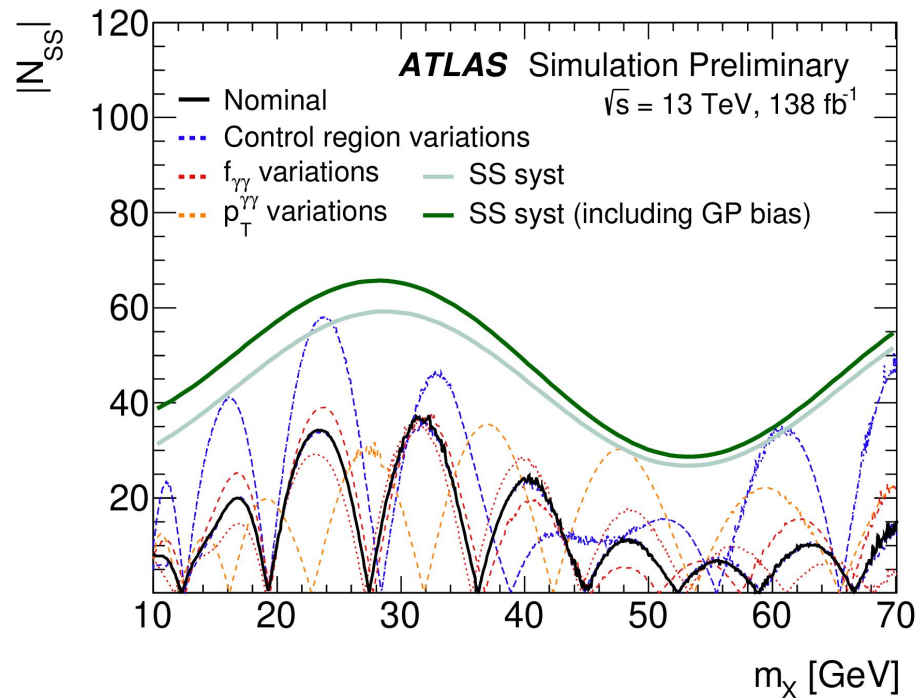
[ATLAS-CONF-2022-018](#)

Estimation of bias arising from the choice of the background model :

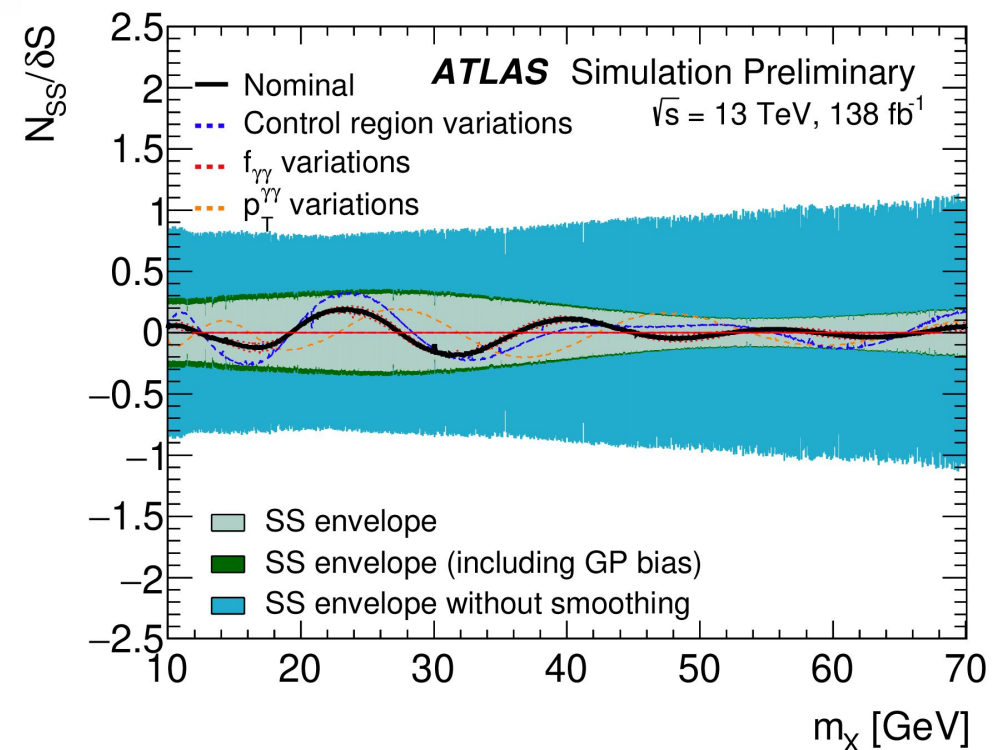
- signal-plus-background fits to background-only templates
- any fitted signal yield is denoted “spurious signal” (SS) and is a systematic uncertainty

Background templates are affected from low statistics :

- the Gaussian Process Regression (GPR) method
 - mitigates statistical fluctuations on the background shape
 - GPR decreases the SS systematics uncertainty
 - the bias from the GPR methodology is also accounted for



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Source	Uncertainty
On $\sigma_{\text{fid}} \cdot \mathcal{B}(X \rightarrow \gamma\gamma)$ [%]	
Pile-up modeling	± 3.5 (at 10 GeV) – ± 2 (beyond 15 GeV), mass dependent
Photon energy resolution	± 2.5 – ± 2.7 , mass dependent
Scale and PDFs uncertainties	± 2.5 – ± 0.5 , mass dependent
Trigger on close-by photons	± 2 (at 10 GeV) – < 0.1 (beyond 35 GeV), mass dependent
Photon identification	± 2.0
Isolation efficiency	± 2.0
Luminosity (2015–2018)	± 1.7
Trigger	± 1.0
Signal shape modeling	< 1
Photon energy scale	negligible
<i>Background modeling</i>	
Spurious signal (relative to δS)	30-65 events (10-30 %), mass dependent

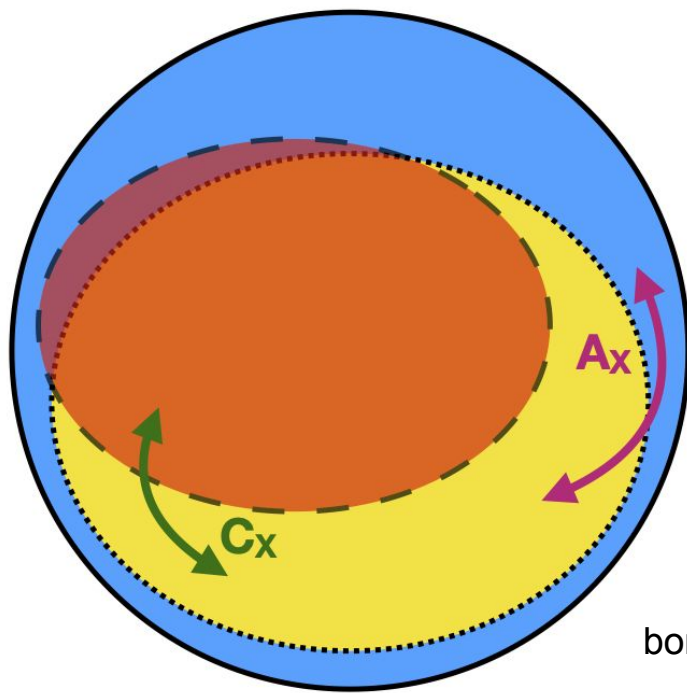
Most systematic uncertainties are percent-level or smaller

- the dominant systematics arises from the background modelling uncertainties
 - spurious signal (SS) and GPR bias combined

The measurement of the signal production cross-section is performed in a fiducial volume:

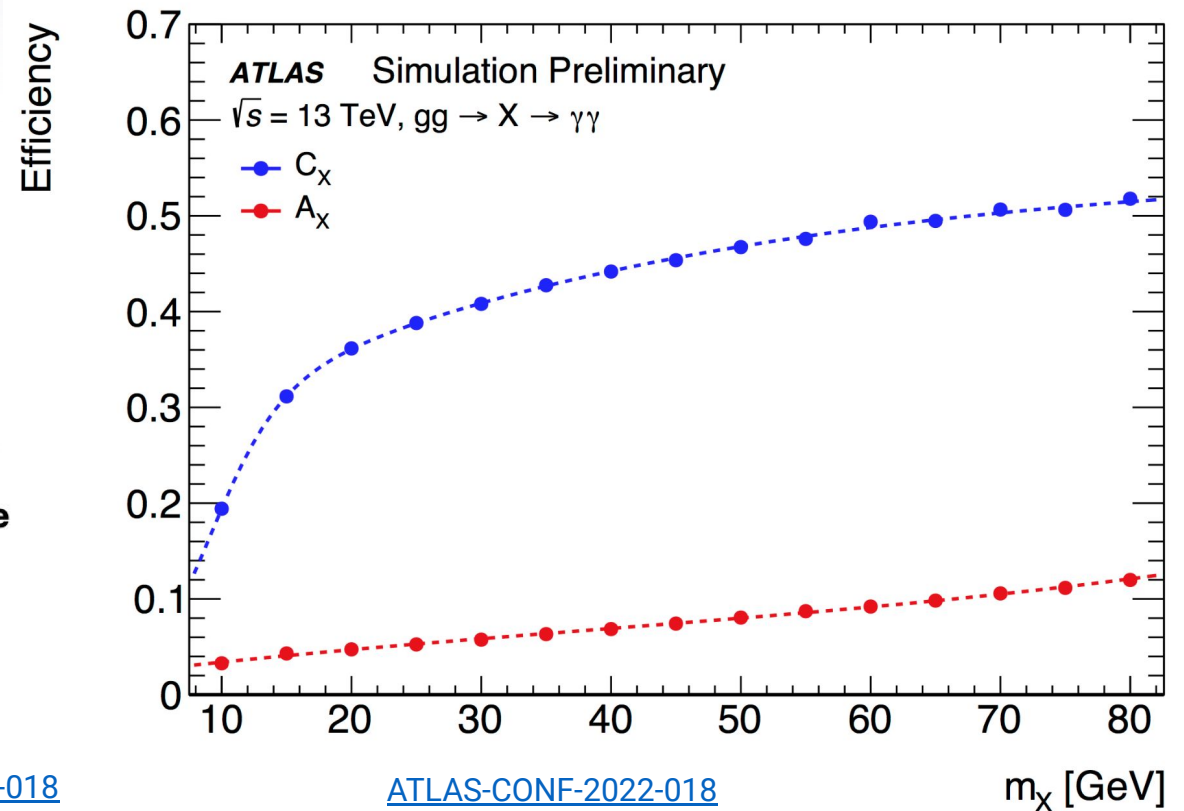
- defined at generator level from the signal MC simulation samples
 - $p_T^{\gamma\gamma} > 50 \text{ GeV}$
 - $E_T > 22 \text{ GeV}$ for each photon
 - particle-level isolation $E_T^{iso} < 0.05 \times E_T$ within a $\Delta R < 0.2$ cone

$$\sigma_{\text{fid}} \cdot \mathcal{B}(X \rightarrow \gamma\gamma) = \frac{N_S}{C_X \mathcal{L}}, \text{ with } C_X = \frac{N_{\text{MC}}^{\text{det}}}{N_{\text{MC}}^{\text{fid}}}$$



■ Full phase space
■ Fiducial phase space
■ Detector phase space
 $A_X := \text{Acceptance}$
 $C_X := \text{Unfolding}$

borrowed from [CERN-THESIS-2021-018](https://cds.cern.ch/record/2711111/files/CERN-THESIS-2021-018)

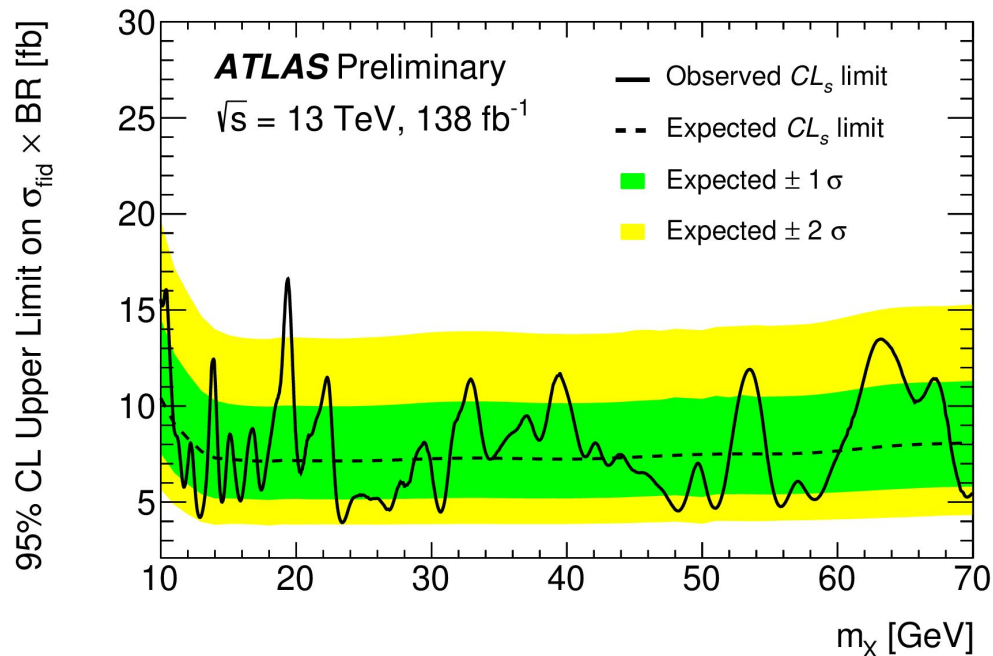


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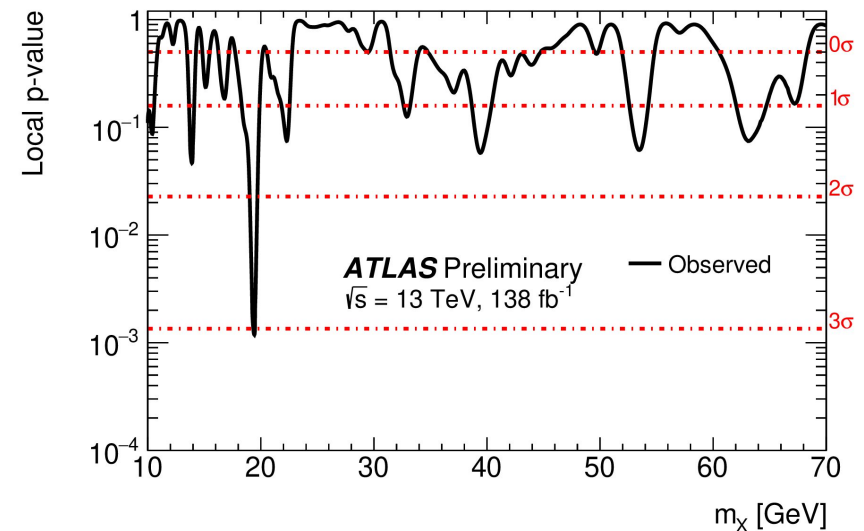
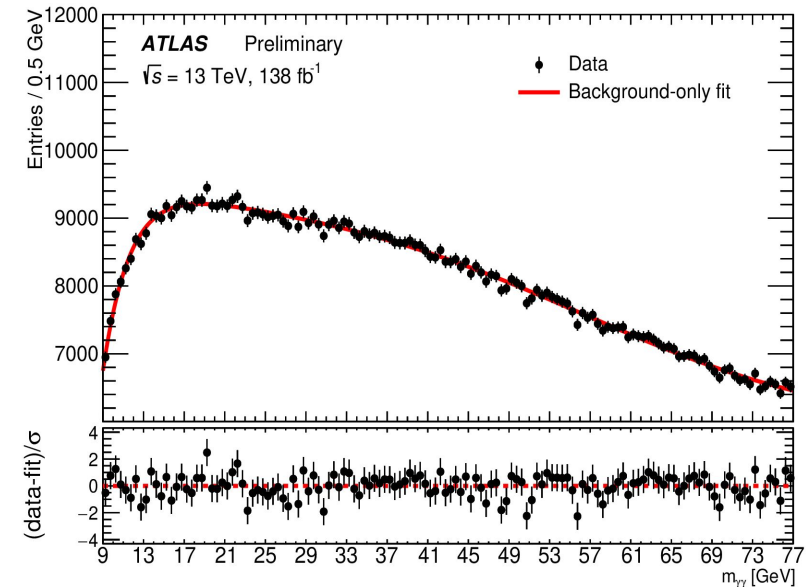
Search performed in the $[10,70]$ GeV mass range

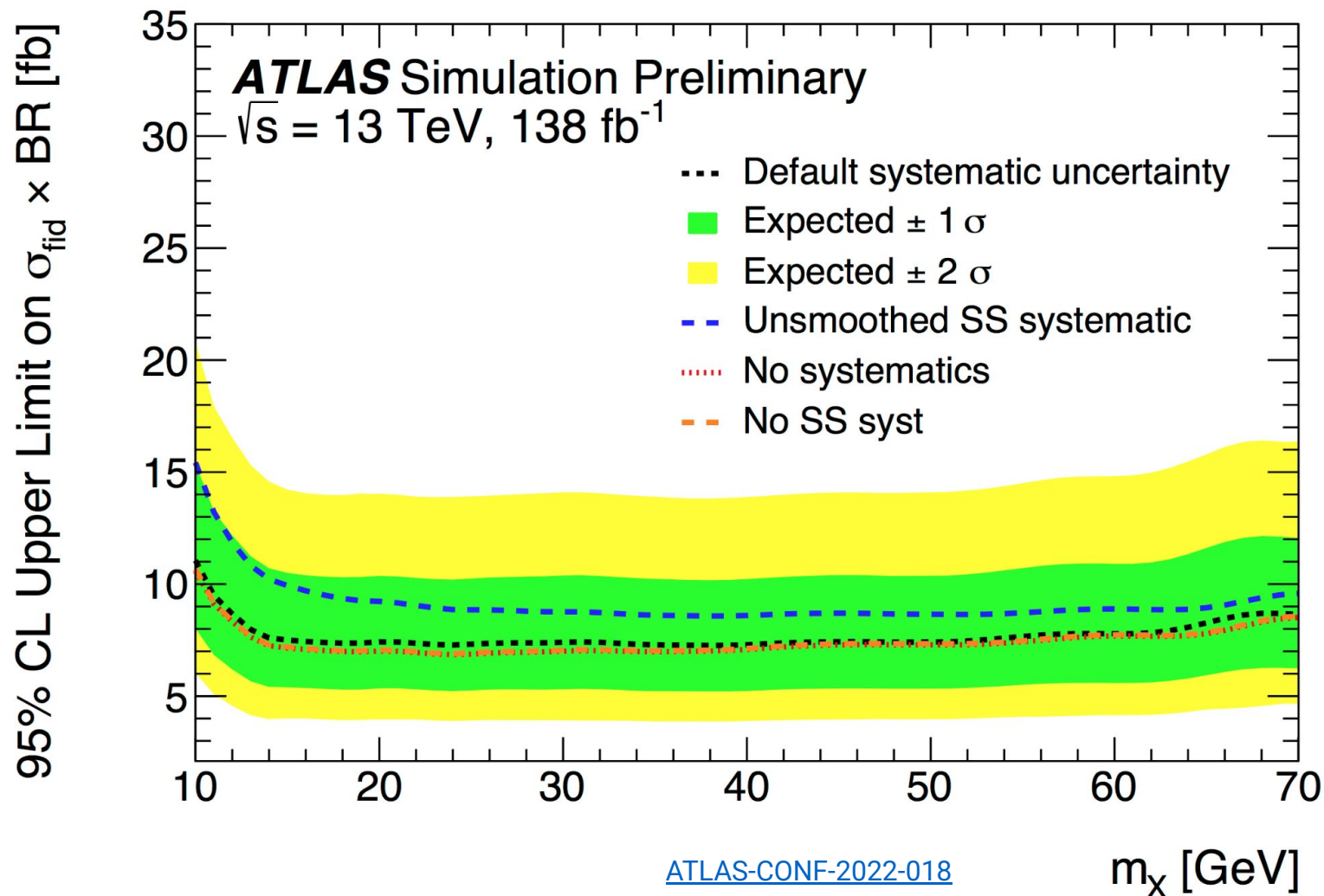
- binned likelihood fit in the $[9,77]$ range
 - (at least 5σ lever-arm from edges)
- parameter of interest: $\sigma_{fid} \times BR(X \rightarrow \gamma\gamma)$
- good description of the data with the background model
 - no significant deviation wrt the SM
 - largest deviation at 19.4 GeV, with 3.1σ local significance
 - $(1.48 \pm 0.02)\sigma$ global significance, evaluated with pseudo-data

Limits set on the $\sigma_{fid} \times BR(X \rightarrow \gamma\gamma)$ of a resonance decaying to two photons



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New Physics scenario being considered:

- all heavy states are beyond the reach of the LHC
 - no deviations from the SM behavior are expected in the TeV range
- a scalar a , singlet of the SM gauge group, naturally lighter than the EW scale exists
 - a is abundantly produced in proton-proton collisions
 - a decays promptly into a pair of SM particles with a narrow width
- a “KSVZ-ALP” model is considered, inspired by the simplest QCD axion model of the scalar a :

$$\mathcal{L}_{\text{int}} = \frac{a}{4\pi f_a} \left[\alpha_3 c_3 G^a \tilde{G}^a + \alpha_2 c_2 W^i \tilde{W}^i + \alpha_1 c_1 B \tilde{B} \right]$$

- barring a huge hierarchy among the anomaly coefficients:
 - for $m_a \lesssim m_Z$, the relevant two-body decays of a are to photons and to jets
 - the width into gluons dominates over the one into photons
 - the total width is dominated by its coupling to gluons and is always small compared to its mass

Limits recast into the ALP parameter space

- strongest limits on a hypothetical resonance produced in gluon fusion that decays to two photons

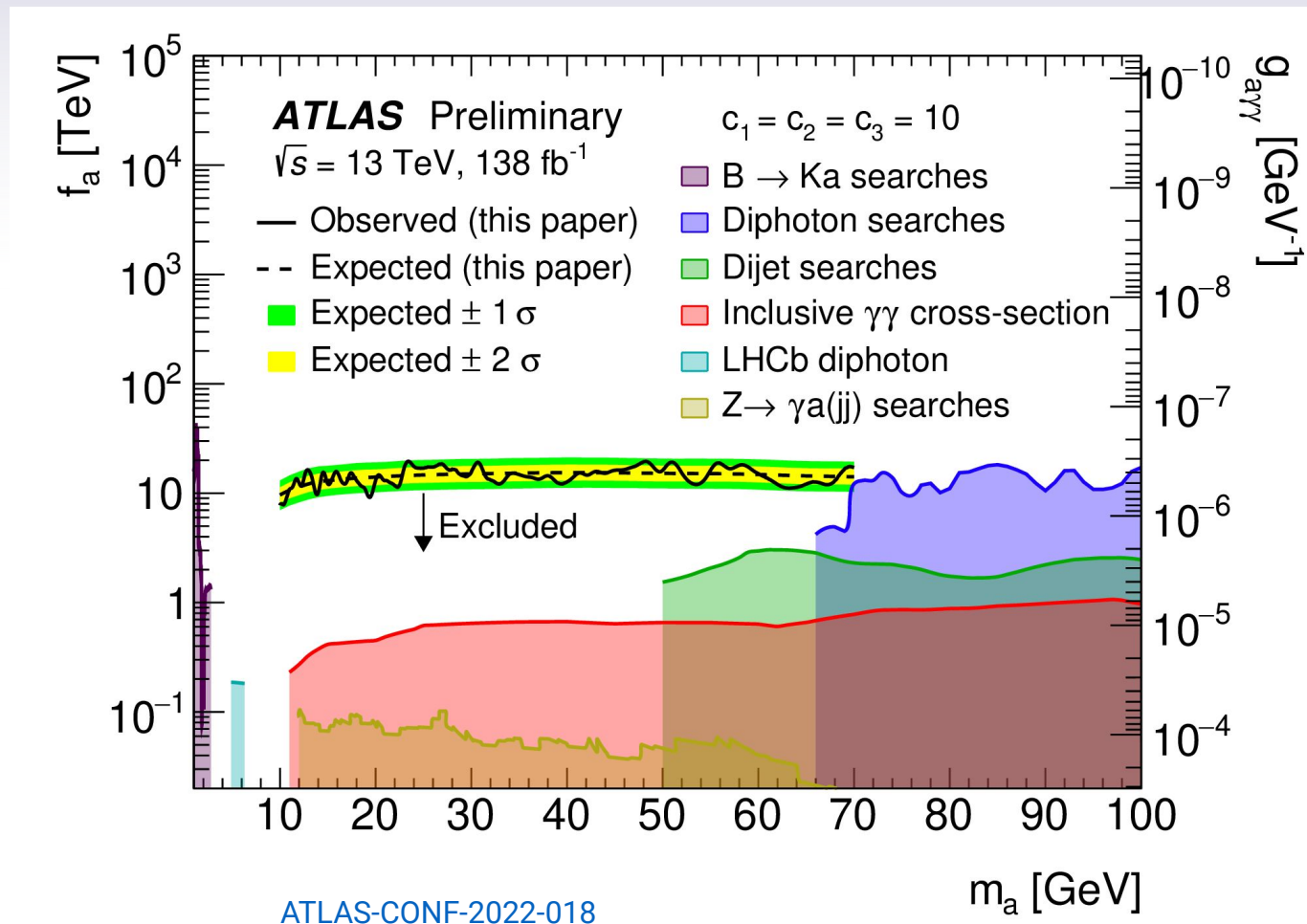
Other searches probing the same mass range:

- light-by-light scattering in heavy ion collisions significantly limited by the production mechanism
- dijet searches disfavoured by mass resolution

Other diphoton searches in proton proton collisions:

- CMS 13 dominates down to 70 GeV (35.9 fb^{-1})
- ATLAS extends the limit down to 65 GeV (80 fb^{-1})

A large piece of the $\gamma\gamma$ gap is now covered !



ATLAS searched for boosted resonances in the diphoton channel, with masses in the 10 to 70 GeV range

Analysis strategy:

- strongly relies on the excellent performance of the EM calorimeter
- novel selection of boosted diphoton pairs to reach masses below the trigger turn-on
- observed data in agreement with the SM-only (no excess) hypothesis
- largest deviation found at 19.4 GeV
 - corresponding to a 3.1σ (1.5σ) local (global) significance
- limits on $\sigma_{fid} \times BR(X \rightarrow \gamma\gamma)$ from 4 fb to 17 fb
- the total uncertainty is dominated by statistics
 - impact of background modelling mitigated by GPR

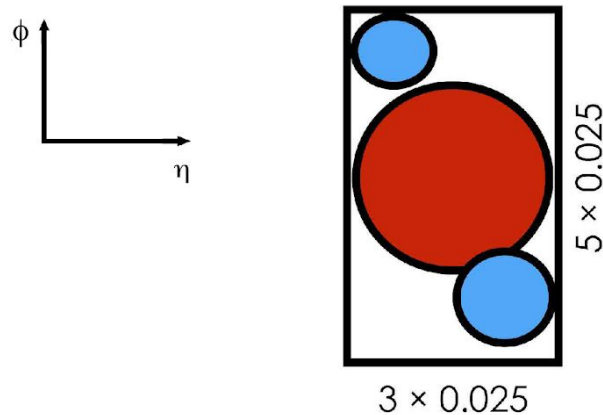
This analysis provides the strongest upper limits up to date using pp collisions:

- on the cross-section times branching ratio of a resonance that decays to two photons
- in the mass range below 65 GeV, and down to 10 GeV
- and in the ALP parameter space in that same mass range

<https://atlas.web.cern.ch/Atlas/GROUPS/PHYSICS/CONFNOTES/ATLAS-CONF-2022-018/>

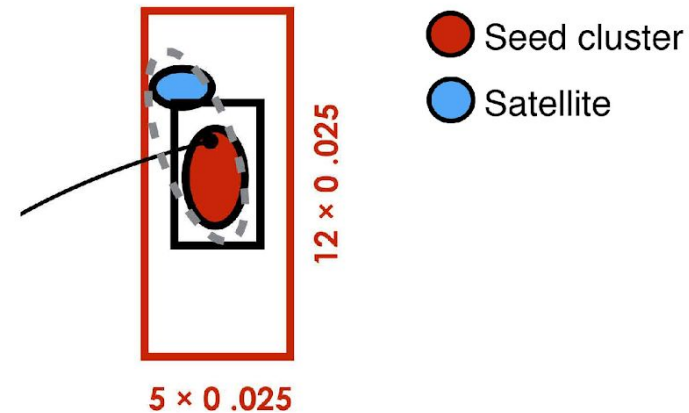
All $e^\pm, \gamma$:

Add all clusters within 3×5 window around seed cluster.



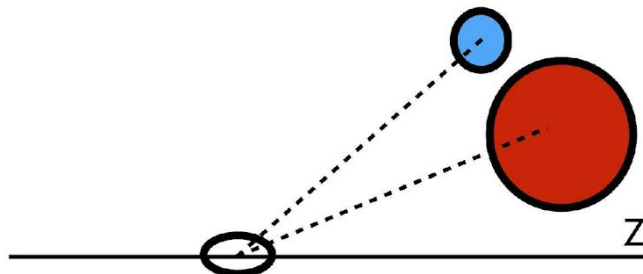
Electrons only:

Seed, secondary cluster **match the same track**.

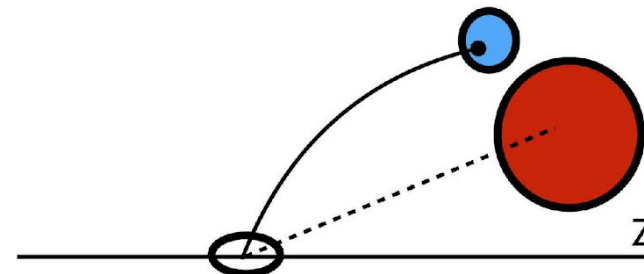


Converted photons only:

Add topo-clusters that have the **same conversion vertex** matched as the seed cluster.



Add topo-clusters with a **track match** that is **part of the conversion vertex** matched to the seed cluster.



New Physics scenario being considered:

- all heavy states are beyond the reach of the LHC
 - no deviations from the SM behavior are expected in the TeV range
- a scalar a , singlet of the SM gauge group, naturally lighter than the EW scale exists
 - a is abundantly produced in proton-proton collisions
 - a decays promptly into a pair of SM particles with a narrow width
- a “KSVZ-ALP” model is considered, inspired by the simplest QCD axion model of the scalar a :

$$\mathcal{L}_{\text{int}} = \frac{a}{4\pi f_a} \left[\alpha_3 c_3 G^a \tilde{G}^a + \alpha_2 c_2 W^i \tilde{W}^i + \alpha_1 c_1 B \tilde{B} \right]$$

- barring a huge hierarchy among the anomaly coefficients:
 - for $m_a \lesssim m_Z$, the relevant two-body decays of a are to photons and to jets
 - the width into gluons dominates over the one into photons
 - the total width is dominated by its coupling to gluons and is always small compared to its mass