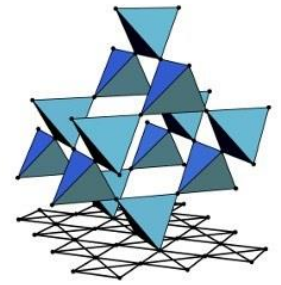




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concept



SFB 1143

Solid State Theory (SS2020)

Lecture 15: Semiclassical electron transport theory

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§ 4.1 Drude theory

- Assumptions of the classical theory:

{ electrons as classical particles
collisions with a relaxation time τ

DC conductivity:

$$\vec{j} = \sigma \vec{E}, \quad \sigma = \frac{ne^2\tau}{m}$$

§ 4.1 Drude theory

- Thermal conductivity:

Fourier's law:

$$\vec{j}^q = -\kappa \vec{\nabla} T$$

↪ thermal conductivity



1D case:

$$j^q = -\kappa \frac{dT}{dx}$$



Warmer region: larger electron (average) speed

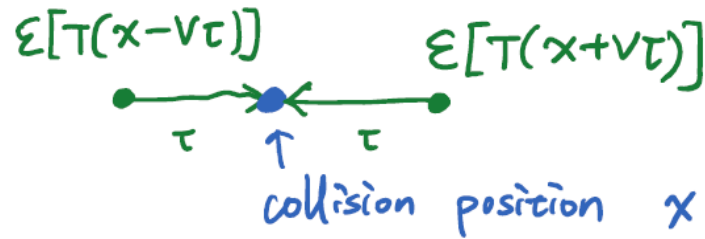
§ 4.1 Drude theory

thermal energy per electron (1D):

$$\mathcal{E}(T) = \frac{1}{2} m v^2 = \frac{1}{2} k_B T \quad (3D: \frac{1}{2} m v^2 = \frac{3}{2} k_B T)$$

→ T is now position dependent!

§ 4.1 Drude theory



Energy flow per electron:

$$\epsilon[T(x - v\tau)] - \epsilon[T(x + v\tau)]$$

length:

$$v\tau$$

$$v\tau$$

of electrons:

$$\frac{1}{2} n v \tau$$

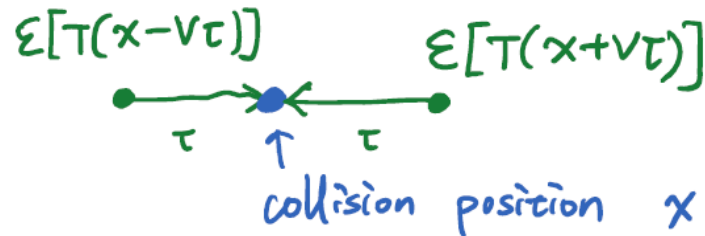
(from left)

$$\frac{1}{2} n v \tau$$

(from right)

electron density: $n = \frac{N_e}{L}$

§ 4.1 Drude theory



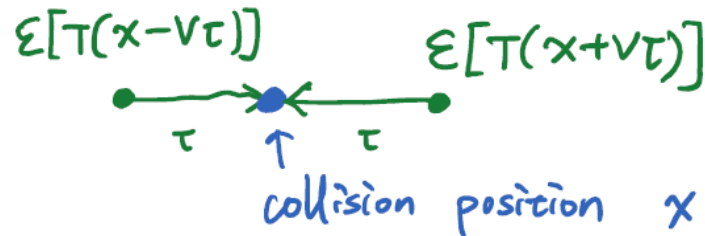
Energy flow per electron:

$$\epsilon[T(x-v\tau)] - \epsilon[T(x+v\tau)]$$

\Rightarrow energy transport within time τ :

$$E^g = \frac{n}{2} v \tau \left(\epsilon[T(x-v\tau)] - \epsilon[T(x+v\tau)] \right)$$

§ 4.1 Drude theory



Energy flow per electron:

$$\epsilon[T(x - vt)] - \epsilon[T(x + vt)]$$

\Rightarrow thermal current (energy transport per unit time):

$$J^q = \frac{E^q}{\tau} = \frac{n}{2} v \left(\epsilon[T(x - vt)] - \epsilon[T(x + vt)] \right)$$

§ 4.1 Drude theory

Assumption: mean free path $l = v\tau$ small

$$J^q = \frac{1}{2} n v \left(\varepsilon[T(x-v\tau)] - \varepsilon[T(x+v\tau)] \right)$$

$$= \frac{1}{2} n v \cdot (-2v\tau) \frac{d\varepsilon}{dx}$$

$$= \underbrace{-n v^2 \tau}_{\text{m}} \underbrace{\frac{d\varepsilon}{dT}}_{\text{w}} \cdot \frac{dT}{dx} \quad \leftarrow \quad n \frac{d\varepsilon}{dT} = \frac{N}{L} \frac{d\varepsilon}{dT} = \frac{1}{L} \frac{dE}{dT}$$

$$= C_v \text{ (heat capacity)}$$

§ 4.1 Drude theory

From 1D to 3D:

$$\frac{dT}{dx} \rightarrow \vec{\nabla} T$$

$$\langle V_x^2 \rangle + \langle V_y^2 \rangle + \langle V_z^2 \rangle = \frac{1}{3} V^2$$

$$\Rightarrow J^q = \frac{1}{3} v^2 \tau n \frac{d\varepsilon}{dT} \cdot (-\vec{\nabla} T)$$

$$\frac{N}{V} \frac{d\varepsilon}{dT} = \frac{1}{V} \frac{dE}{dV} = C_V \quad \begin{array}{l} \rightarrow E = N\varepsilon : \text{total} \\ \text{internal energy} \end{array}$$

$$\Rightarrow J^q = k (-\vec{\nabla} T) \quad \text{with} \quad k = \frac{1}{3} v^2 \tau C_V$$

§ 4.1 Drude theory

- Wiedemann-Franz law:

$$\frac{\kappa}{\sigma} = \frac{\frac{1}{3} v^2 C_V}{\frac{ne^2}{m}} = \frac{\frac{1}{3} m v^2 C_V}{ne^2}$$

$$C_V = \frac{3}{2} n k_B, \quad \frac{1}{2} m v^2 = \frac{3}{2} k_B T$$

classical! \Rightarrow $= \frac{3}{2} \left(\frac{k_B}{e}\right)^2 T$ ← universal!

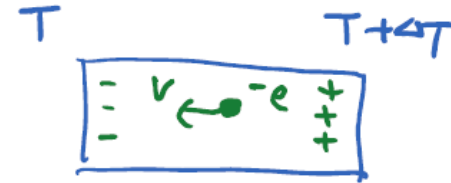
§ 4.1 Drude theory

- Thermoelectric effect:

Heat flow generates electric field!

$$\vec{E} = Q \vec{\nabla} T$$

↓
thermopower



§ 4.1 Drude theory

- “Net” velocity at position x (due to the temperature gradient):



$$\begin{aligned} v_d &= \frac{1}{2} [v(x - v\tau) - v(x + v\tau)] \\ &= \frac{1}{2} \cdot (-2v\tau) \frac{dv}{dx} \\ &= -v\tau \frac{dv^2}{dx} \\ &= -\frac{1}{2} \tau \frac{dv^2}{dT} \cdot \frac{dT}{dx} \end{aligned}$$

Constant (with an arrow pointing to the $\frac{dT}{dx}$ term)

§ 4.1 Drude theory

From 1D to 3D: $v_x^2 = v_y^2 = v_z^2 = \frac{1}{3} v^2$

$$\Rightarrow \vec{v}^g = -\frac{1}{6} \tau \frac{dv^2}{dT} (\vec{\nabla} T)$$

$$\vec{v}^E = -\frac{e\vec{E}\tau}{m}$$

Due to the acceleration from the thermally generated electric field \vec{E}

§ 4.1 Drude theory

From 1D to 3D: $v_x^2 = v_y^2 = v_z^2 = \frac{1}{3} v^2$

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Due to the acceleration from the thermally generated electric field \vec{E}

Steady state: $\vec{v}^g + \vec{v}^E = 0$

$$\frac{e\vec{E}\tau}{m} = -\frac{\tau}{6} \frac{dv^2}{dT} (\vec{\nabla} T)$$

§ 4.1 Drude theory

$$\begin{aligned}\vec{E} &= -\frac{1}{b} \frac{m}{e} \frac{dv^2}{dT} (\vec{\nabla}T) \\ &= -\frac{1}{3e} \frac{d}{dT} \left(\frac{3}{2} k_B T \right) (\vec{\nabla}T) \quad \leftarrow \text{use } \frac{1}{2} m v^2 = \frac{3}{2} k_B T \\ &= -\frac{C_V}{3ne} (\vec{\nabla}T) \quad \leftarrow C_V = \frac{3}{2} n k_B\end{aligned}$$

thermopower $Q = -\frac{C_V}{3ne}$

$$C_V = \frac{3}{2} n k_B \quad \Rightarrow \quad Q = -\frac{k_B}{2e} \sim 10^{-4} \text{ volt/K}$$

§ 4.1 Drude theory

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$$C_V = \frac{3}{2} n k_B \quad \Rightarrow \quad Q = -\frac{k_B}{2e} \sim 10^{-4} \text{ volt/K}$$

10^2 times larger compared to experiments!

§ 4.2 Semiclassical approach

- Classical theory:

Electron velocity distribution in Drude's theory:

$$f_{\mathbf{B}}(\vec{v}) \propto e^{-\frac{1}{2}mv^2/k_B T}$$



Boltzmann distribution

(electrons treated as classical particles)

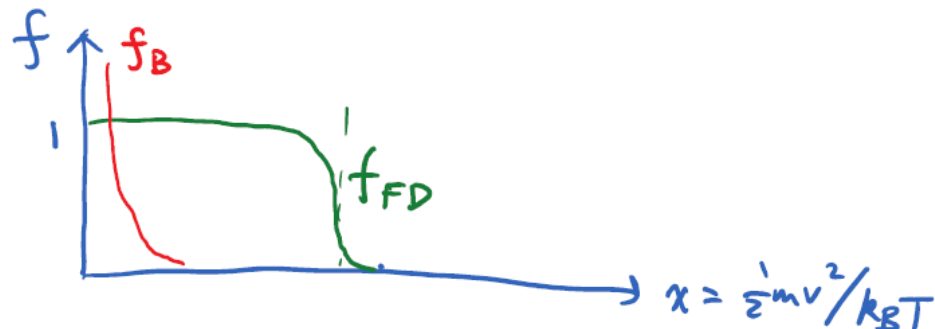
§ 4.2 Semiclassical approach

- Sommerfeld's theory:

free electron gas + Fermi-Dirac distribution

$$f_{\text{FD}}(\vec{v}) \propto \frac{1}{e^{(\frac{1}{2}mv^2 - k_B T_0)/k_B T} + 1}$$

Fermi-Dirac distribution



§ 4.2 Semiclassical approach

$$E_{\vec{k}} = \frac{\hbar^2 \vec{k}^2}{2m}$$

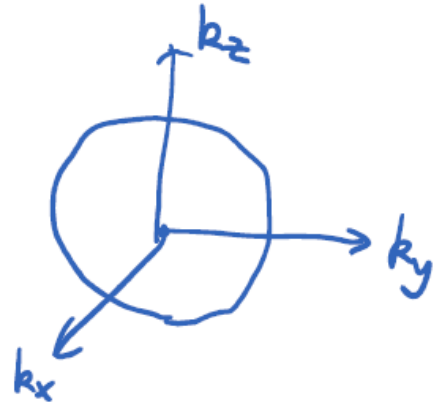
$$\psi_{\vec{k}}(\vec{r}) = \frac{1}{\sqrt{V}} e^{i\vec{k} \cdot \vec{r}}$$

momentum: $\vec{p} = \hbar \vec{\nabla}$

$$\Rightarrow \vec{p} \psi_{\vec{k}}(\vec{r}) = \hbar \vec{\nabla} \psi_{\vec{k}}(\vec{r}) = \hbar \vec{k} \psi_{\vec{k}}(\vec{r})$$

velocity: $\vec{v} = \frac{\vec{p}}{m}$

$$\vec{v} \psi_{\vec{k}}(\vec{r}) = \frac{\hbar \vec{k}}{m} \psi_{\vec{k}}(\vec{r})$$



§ 4.2 Semiclassical approach

Work out $f(\vec{v}) d\vec{v}$ (velocity distribution)

Each \vec{k} point has volume $\frac{(2\pi)^3}{V}$

DOS in \vec{k} -space: $2 \times \frac{1}{(2\pi)^3/V} d\vec{k}$ ($T=0$)

Finite T : $2 \times \frac{V}{(2\pi)^3} \int_{FD} f(\vec{k}) d\vec{k}$

§ 4.2 Semiclassical approach

Work out $f(\vec{v}) d\vec{v}$ (velocity distribution)

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Finite T : $2 \times \frac{V}{(2\pi)^3} \int_{FD} f(\vec{k}) d\vec{k}$


$$\vec{v} = \frac{\hbar \vec{k}}{m} \Rightarrow d\vec{k} = \left(\frac{m}{\hbar}\right)^3 d\vec{v}, \quad \epsilon_{\vec{k}} = \frac{\hbar^2 k^2}{2m} = \frac{1}{2} m \vec{v}^2$$

$$\Rightarrow 2 \times \frac{V}{(2\pi)^3} \frac{1}{e^{(\frac{1}{2} m \vec{v}^2 - \mu)/k_B T} + 1} \left(\frac{m}{\hbar}\right)^3 d\vec{v}$$

$$f(\vec{v}) = \frac{m^3 V}{4\pi^3 \hbar^3} \frac{1}{e^{(\frac{1}{2} m \vec{v}^2 - \mu)/k_B T} + 1}$$

§ 4.2 Semiclassical approach

- Semiclassical approach: electrons treated as wave packets


$$v \approx \frac{\hbar k_F}{m}$$
$$\Delta p \ll \hbar k_F$$

only those electrons close to the Fermi energy are important.

uncertainty principle :

$$\Delta x \cdot \Delta p \sim \hbar$$

$$\Rightarrow \Delta x \sim \frac{\hbar}{\Delta p} \gg \frac{1}{k_F} \approx \underbrace{10^{-10} \text{ m}}_{\substack{\downarrow \\ \sim \text{size of a unit cell}}}$$

§ 4.2 Semiclassical approach

- Semiclassical approach: electrons treated as wave packets

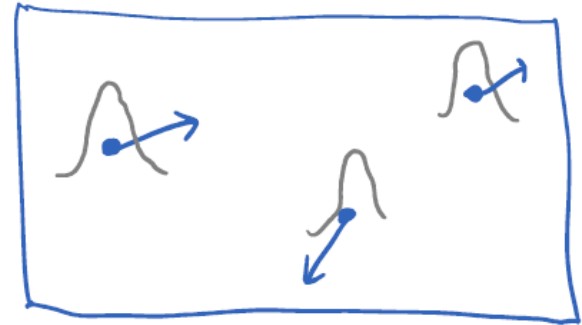
Mean free path:

$$l \approx v_F \tau$$

$\tau \sim 10^{-14} \text{ s}$

$l \sim 10^{-8} \text{ m}$

(room temperature)



§ 4.2 Semiclassical approach

- Semiclassical approach: electrons treated as wave packets

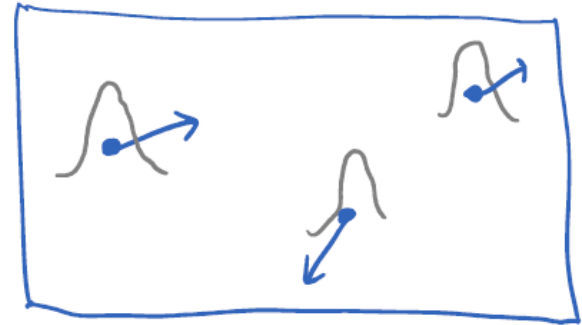
Mean free path :

$$l \approx v_F \tau$$

$\tau \sim 10^{-14} \text{ s}$

$$\sim 10^{-8} \text{ m}$$

(room temperature)



- 1) size of electron wave packet \gg size of a unit cell
- 2) mean free path \gg . . .



semiclassical treatment valid !

§ 4.2 Semiclassical approach

- Thermal conductivity:

$$k = \frac{1}{3} v^2 \tau C_V$$

$v_F^2 = \frac{2\varepsilon_F}{m}$ vs. Drude value $v^2 = \frac{3k_B T}{m}$

$\frac{\pi^2}{2} n k_B \left(\frac{k_B T}{\varepsilon_F} \right)$ vs. Drude value $\frac{3}{2} n k_B$

$$= \frac{\pi^2}{3} \frac{\tau}{m} n k_B^2 T$$

§ 4.2 Semiclassical approach

- Thermal conductivity:

$$\begin{aligned}
 \kappa &= \frac{1}{3} v^2 \tau C_V \\
 &\quad \downarrow \quad \swarrow \\
 & \quad v_F^2 = \frac{2\varepsilon_F}{m} \quad \rightarrow \quad \frac{\pi^2}{2} n k_B \left(\frac{k_B T}{\varepsilon_F} \right) \quad \text{vs. Drude value } \frac{3}{2} n k_B \\
 & \quad \text{vs. Drude value } v^2 = \frac{3k_B T}{m} \\
 &= \frac{\pi^2}{3} \frac{\tau}{m} n k_B^2 T
 \end{aligned}$$

- Wiedemann-Franz law:

$$\frac{\kappa}{\sigma} = \frac{\pi^2}{3} \left(\frac{k_B}{e} \right)^2 T$$

\swarrow
 (Drude value: $\frac{3}{2}$)

$$\sigma = \frac{ne^2 \tau}{m}$$

more accurate now!

§ 4.2 Semiclassical approach

- Thermopower:

$$Q = - \frac{C_V}{3ne} \quad \leftarrow \quad C_V = \frac{\pi^2}{2} n k_B \left(\frac{k_B T}{\epsilon_F} \right)$$
$$= - \frac{1}{6} \frac{k_B}{e} \left(\frac{k_B T}{\epsilon_F} \right)$$
$$\sim - \left(\frac{T}{T_F} \right) \times 10^{-4} \text{ volt/K} \quad \checkmark$$

(Handwritten annotations: 10^2 K above T , 10^4 K below T_F)