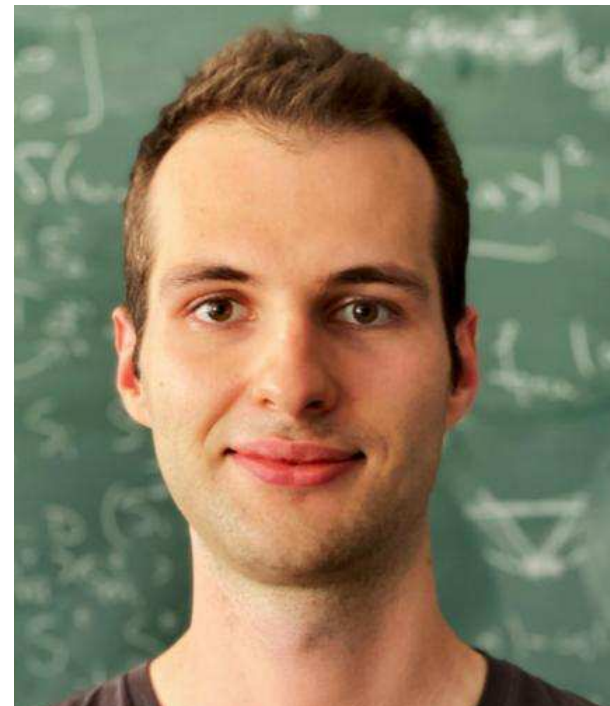
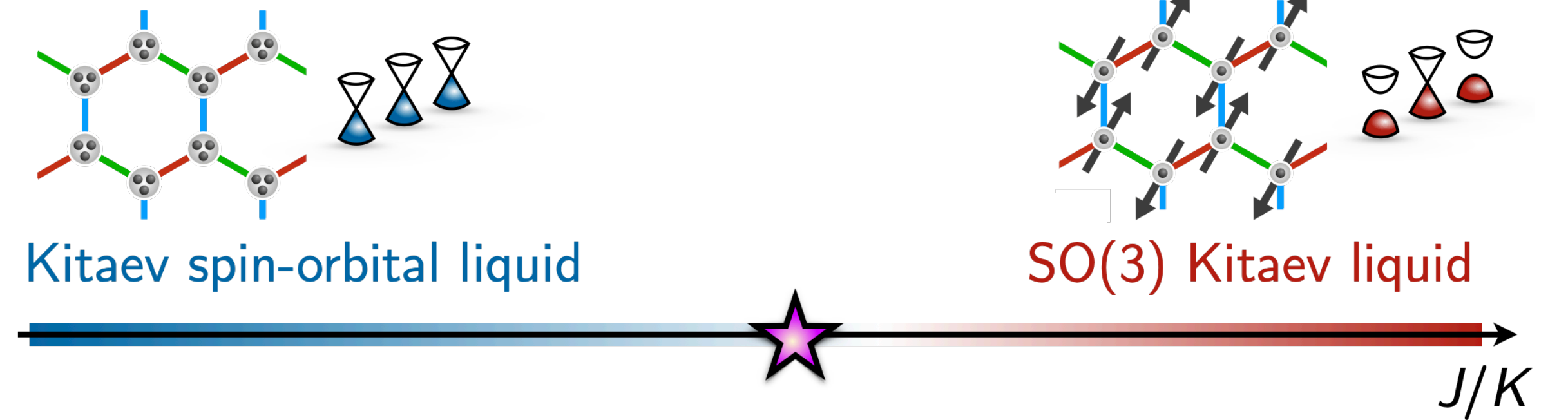


Fractionalized fermionic quantum criticality

Lukas Janssen
TU Dresden



Urban Seifert, Santa Barbara



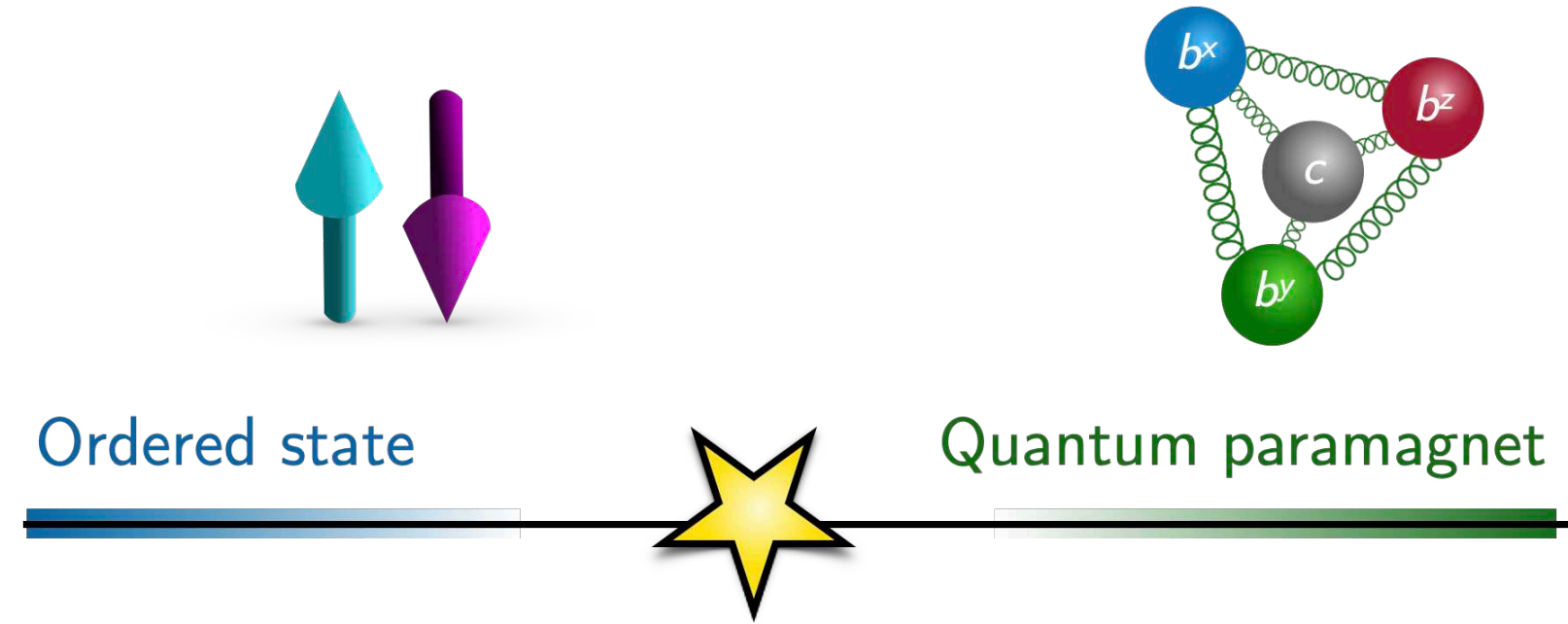
Zihong Liu, Würzburg

Fakher Assaad, Würzburg
Sreejith Chulliparambil, Dresden
Xiao-Yu Dong, Ghent
Hong-Hao Tu, Dresden
Matthias Vojtá, Dresden

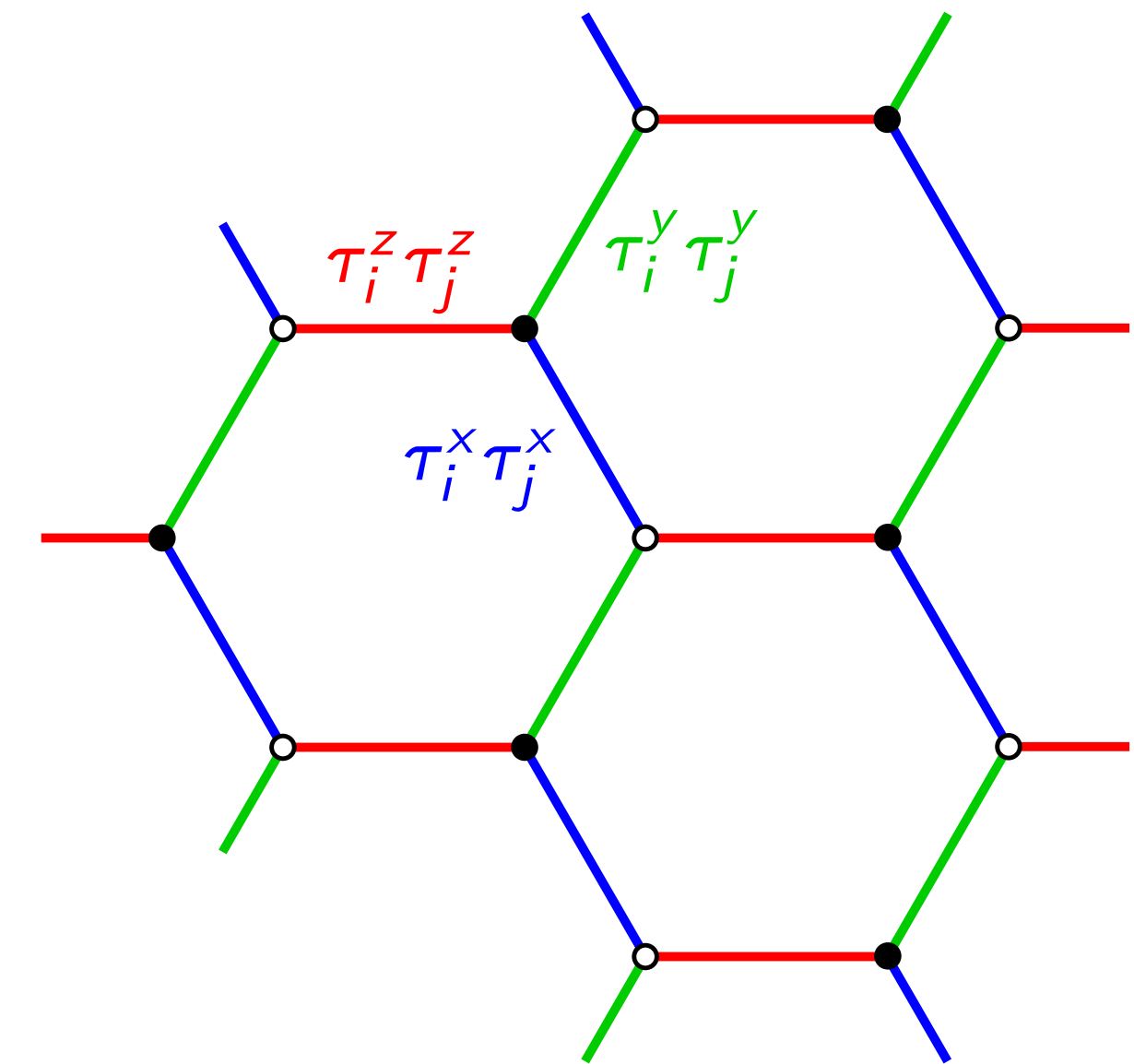


Outline

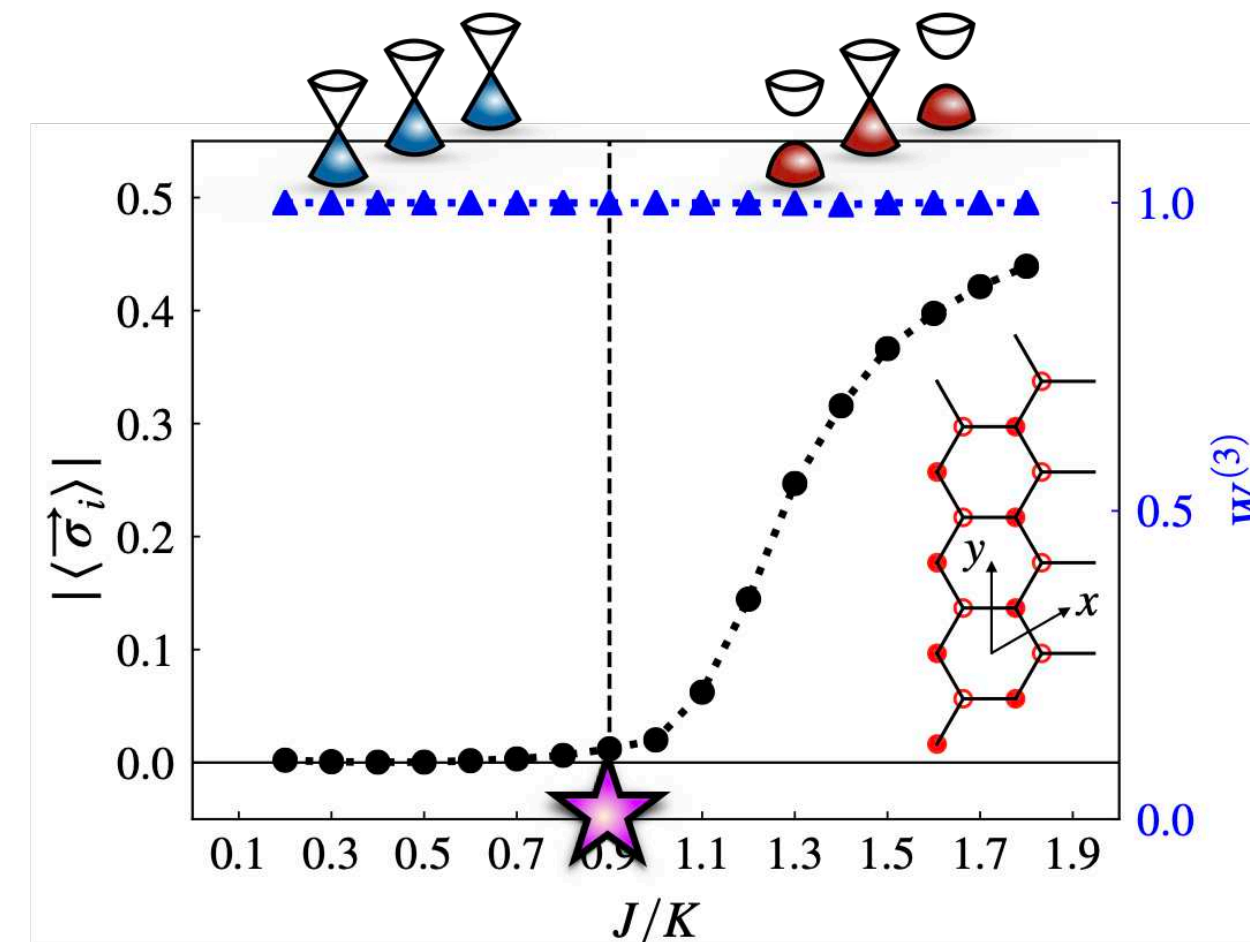
(1) Fractionalized quantum criticality



(2) From Kitaev to Kitaev-Kugel-Khomskii



(3) Kitaev-Heisenberg spin-orbital models



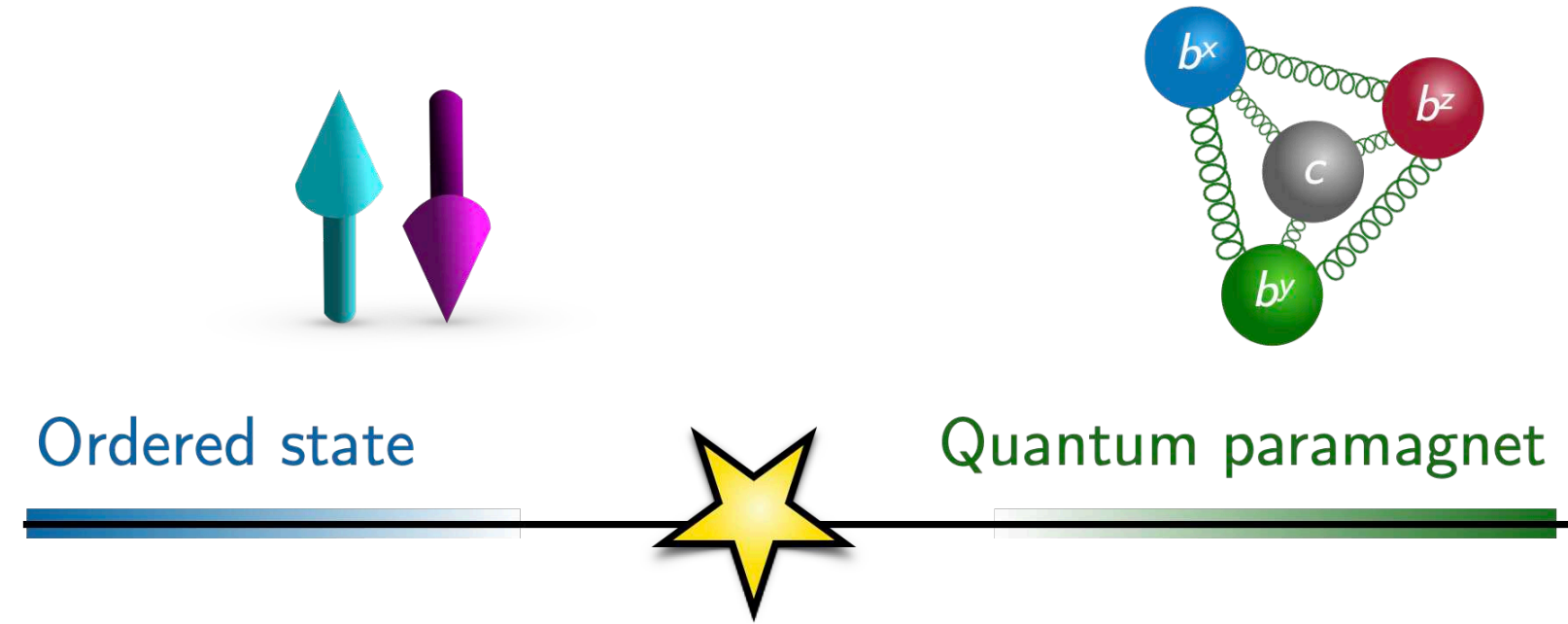
(4) Conclusions



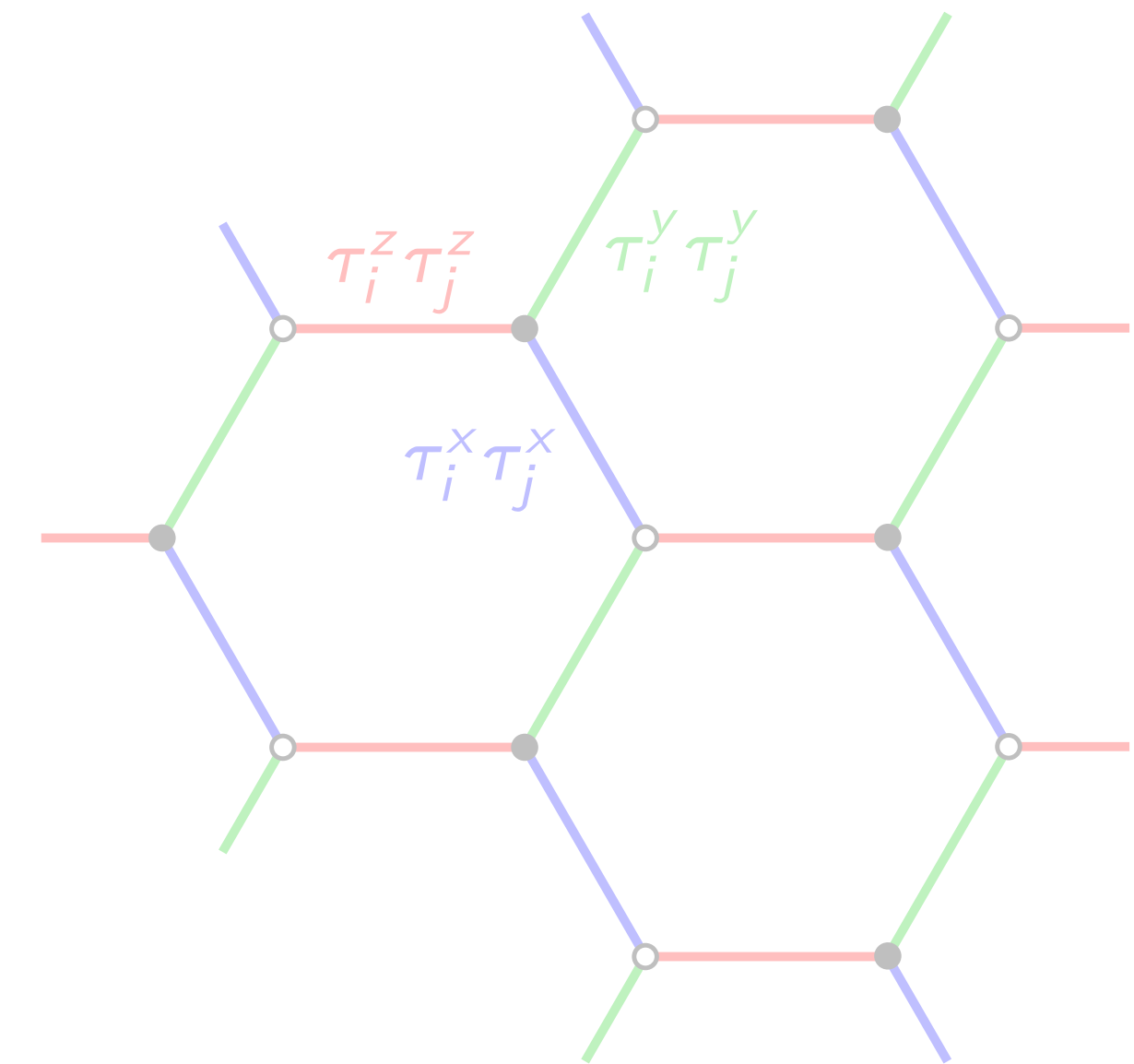
Slides available on <https://tu-dresden.de/physik/qcm/vortraege>

Outline

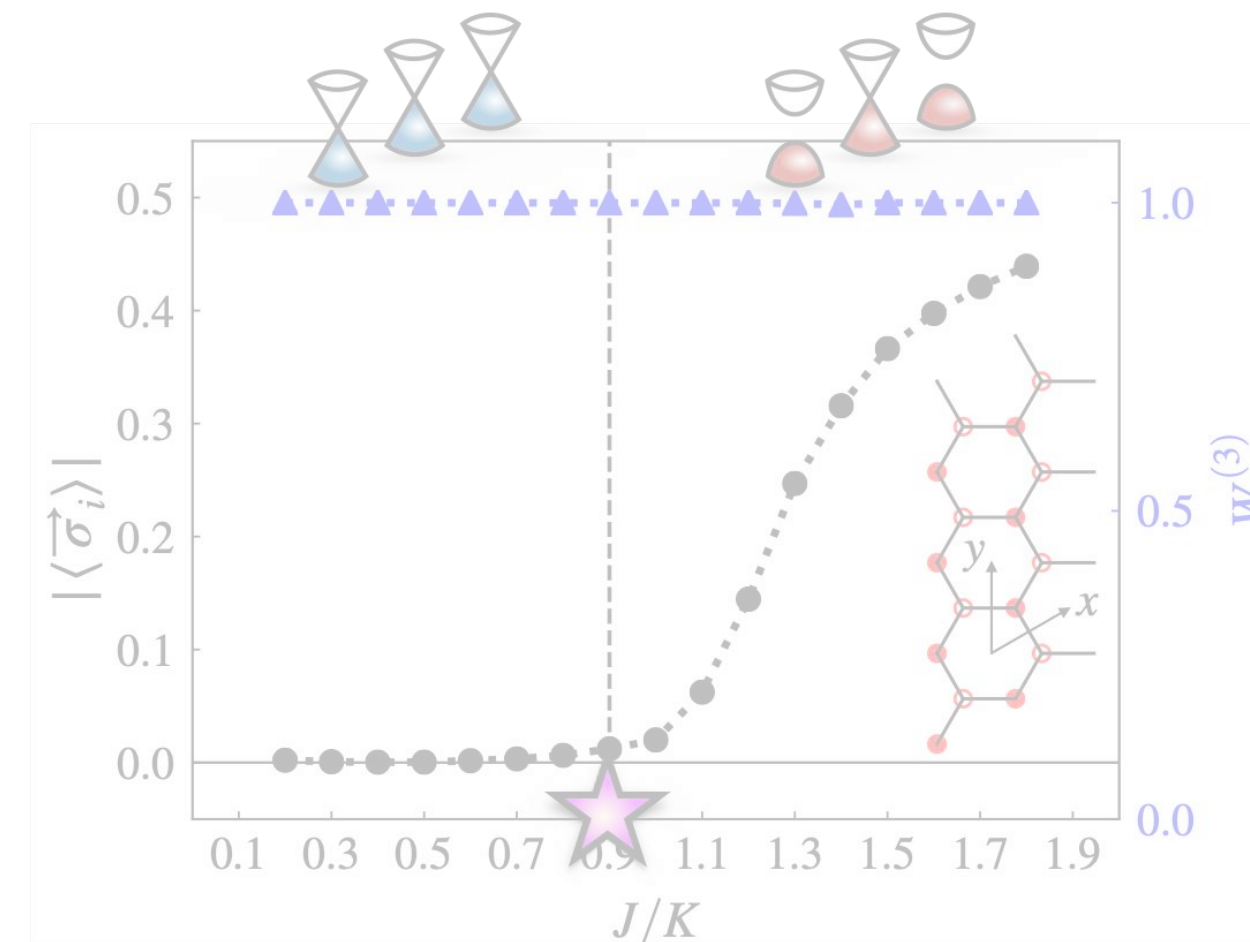
(1) Fractionalized quantum criticality



(2) From Kitaev to Kitaev-Kugel-Khomskii



(3) Kitaev-Heisenberg spin-orbital models



(4) Conclusions

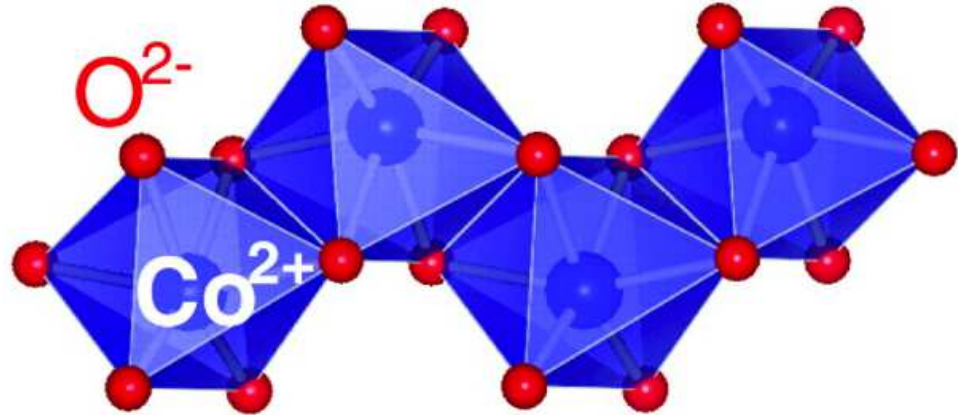
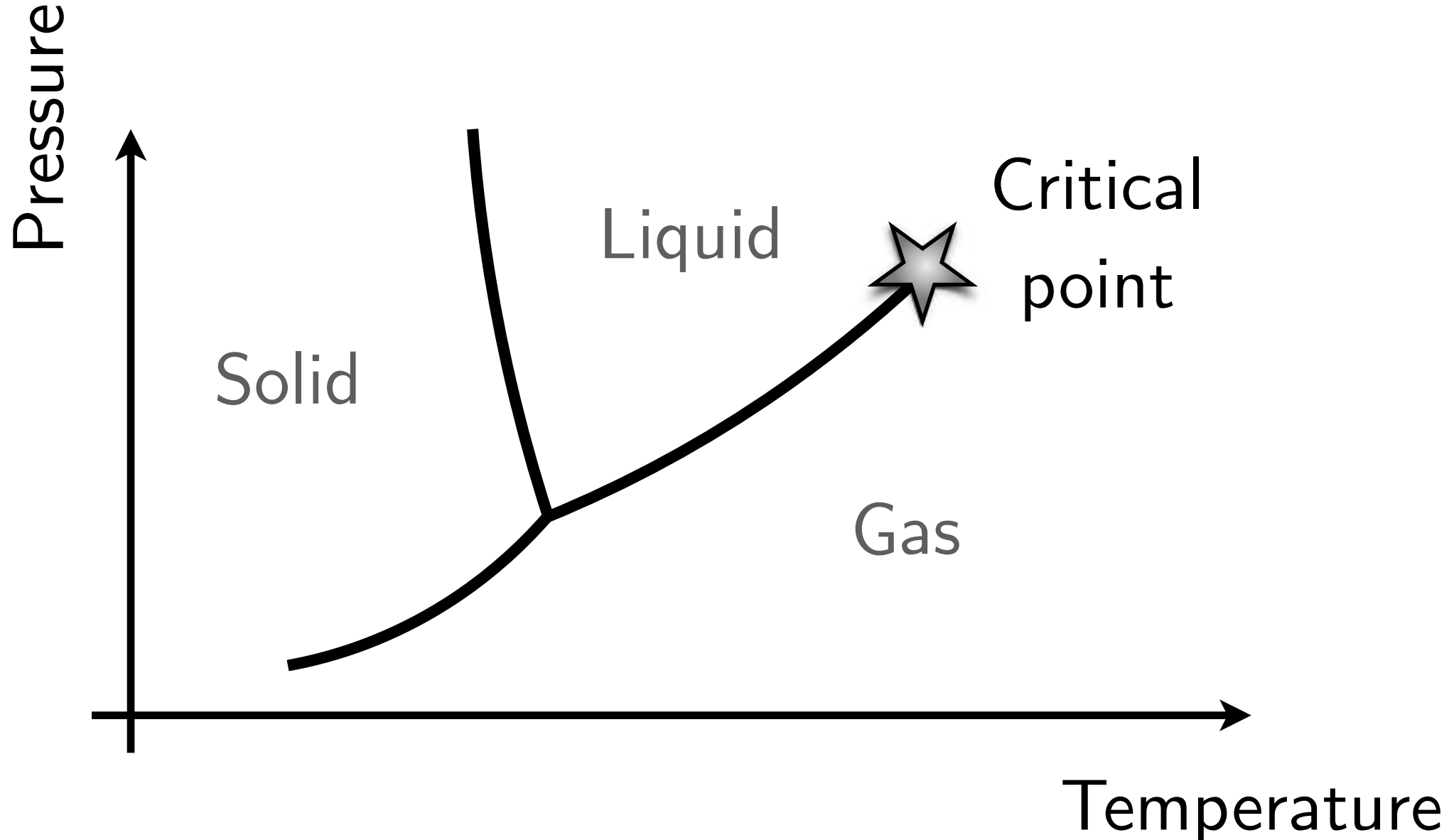


Slides available on <https://tu-dresden.de/physik/qcm/vortraege>

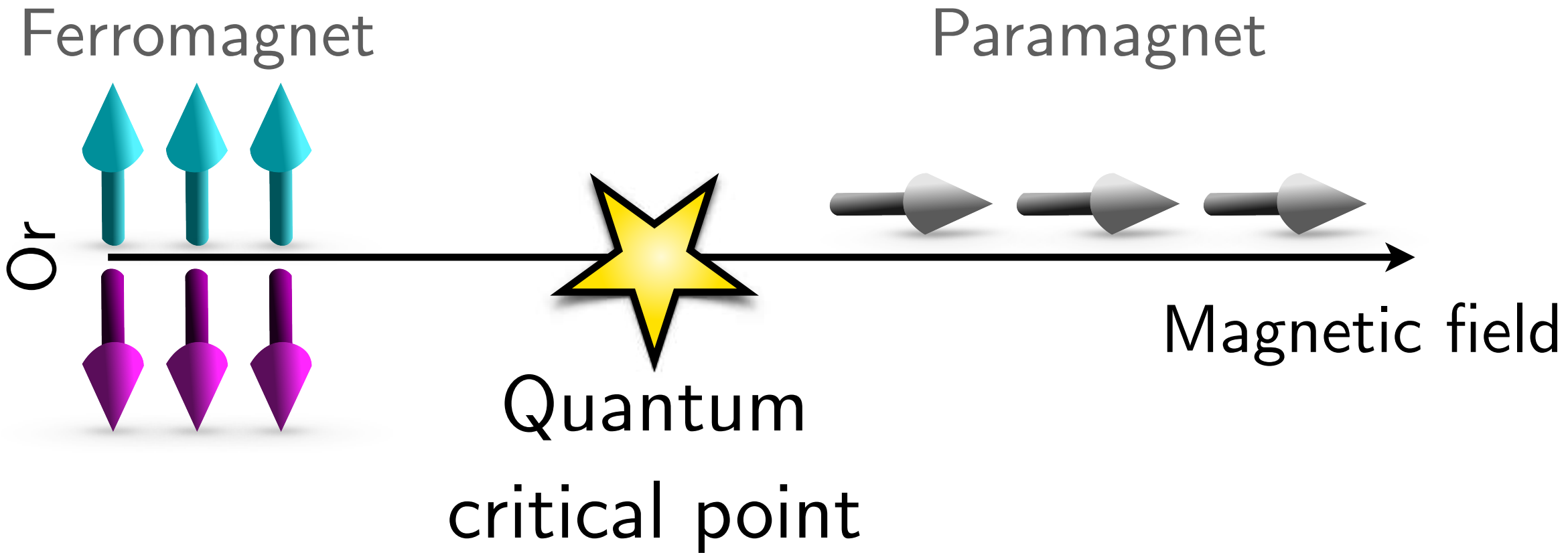
Classical vs quantum criticality



H_2O $T > 0$



CoNb_2O_6 $T \rightarrow 0$



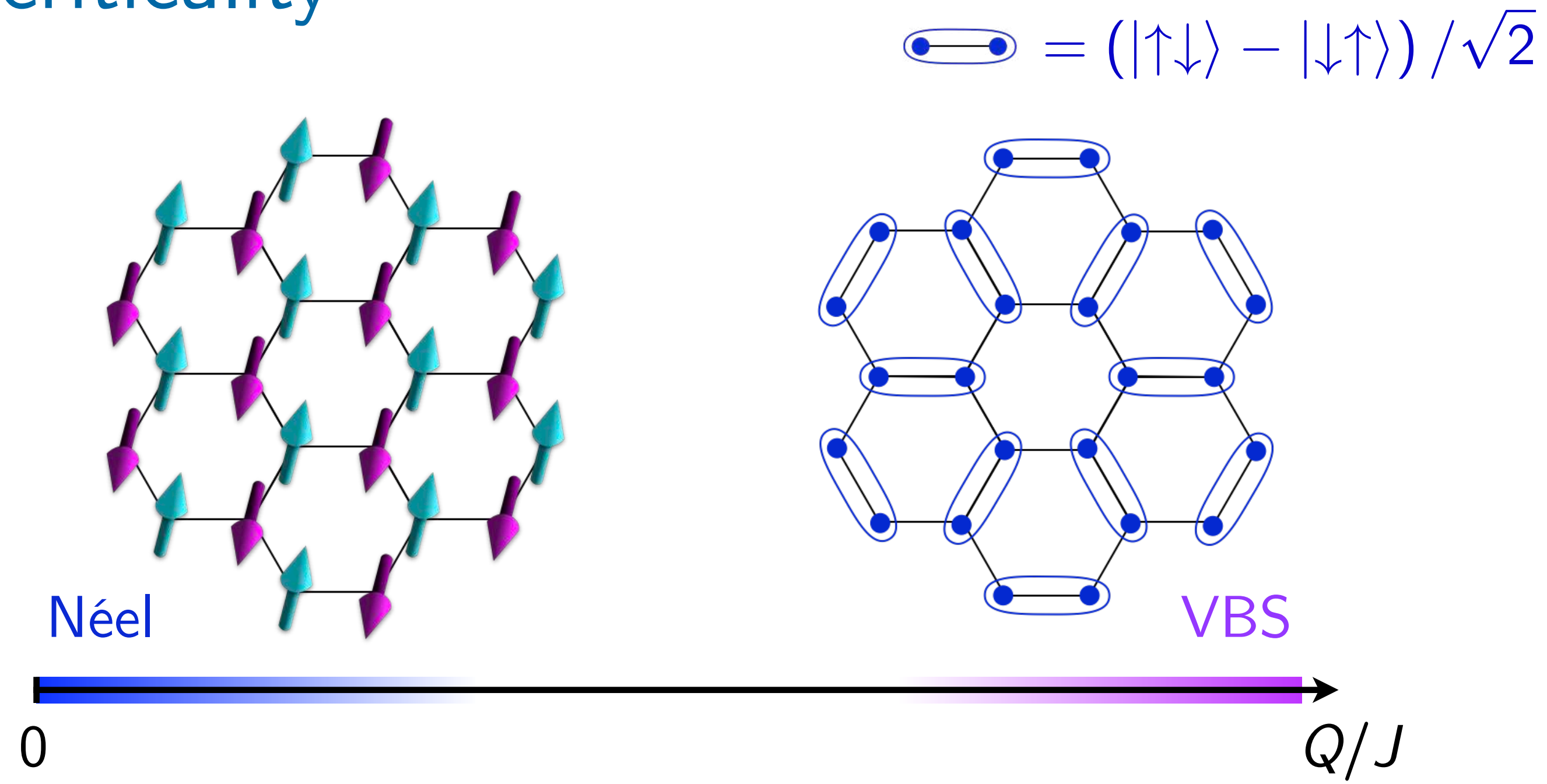
[Coldea *et al.*, Science '10]

[Kinross *et al.*, PRX '14]

[Morris *et al.*, Kaul, Armitage, Nat. Phys. '21]

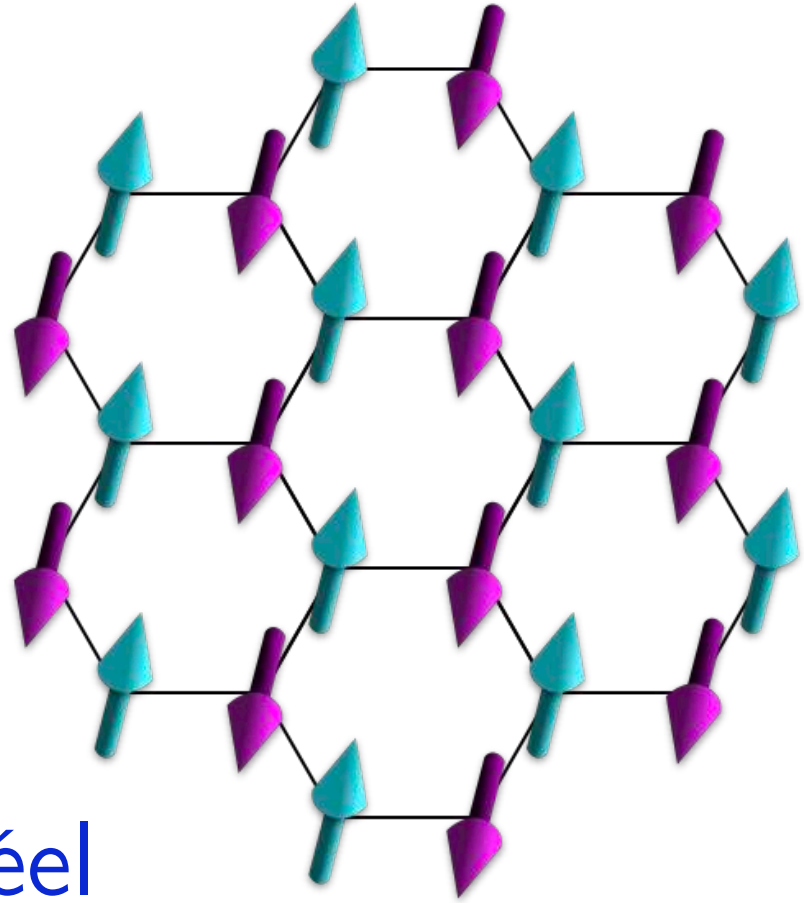
...

Deconfined quantum criticality

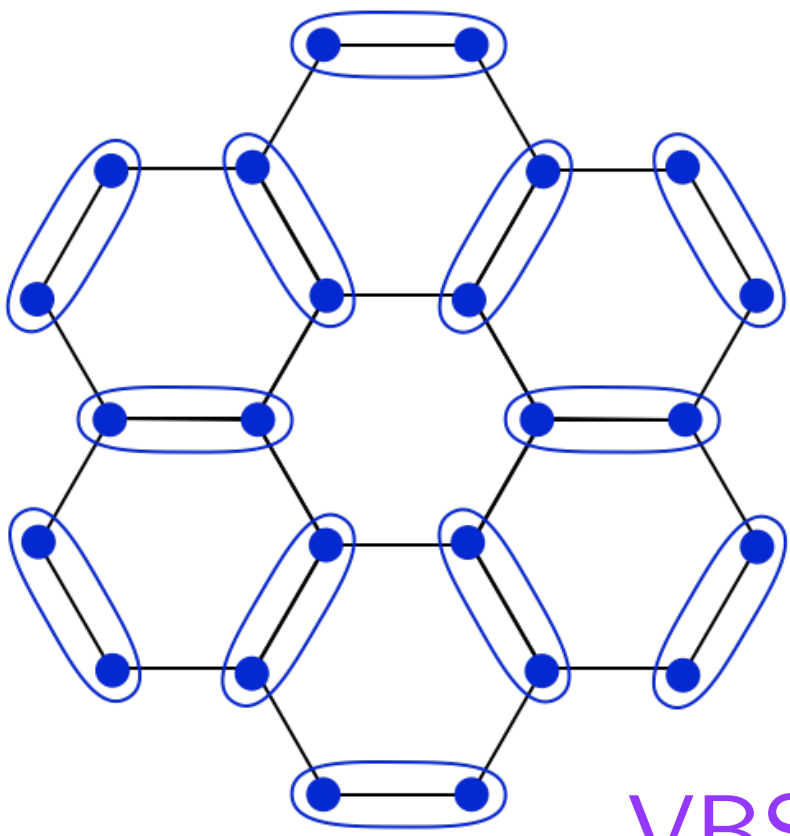


Deconfined quantum criticality

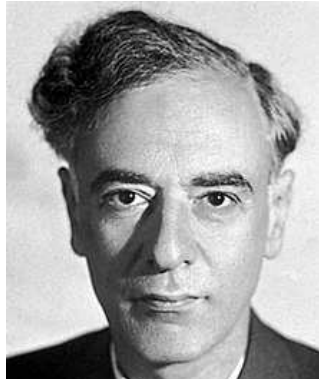
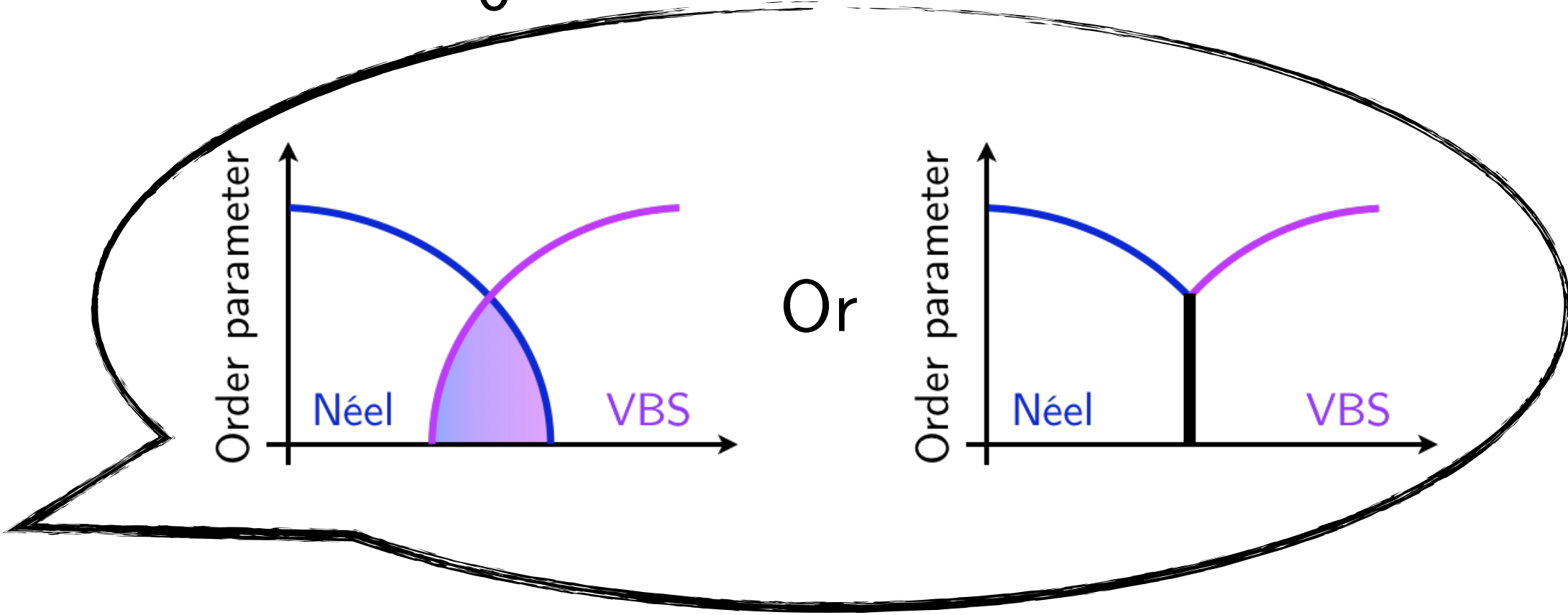
$$\text{---} = (|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle) / \sqrt{2}$$



Néel



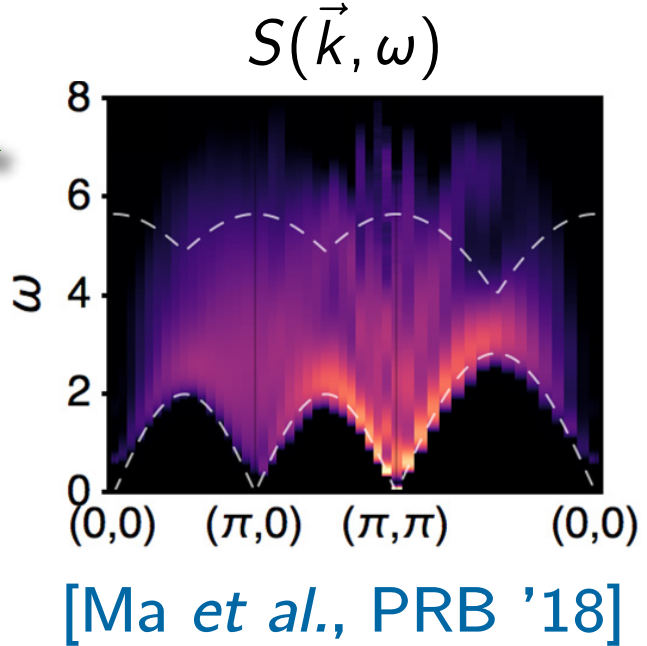
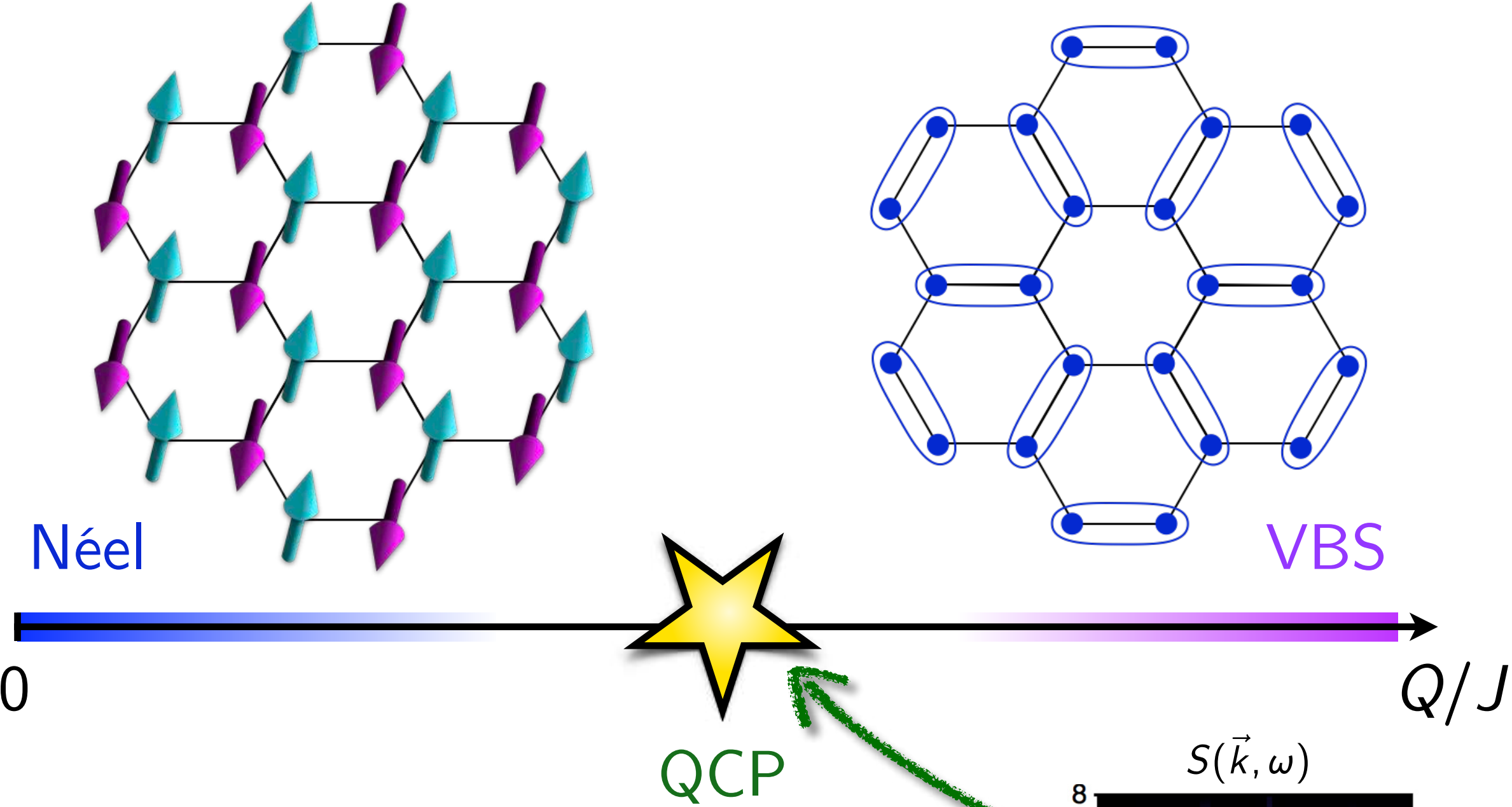
VBS



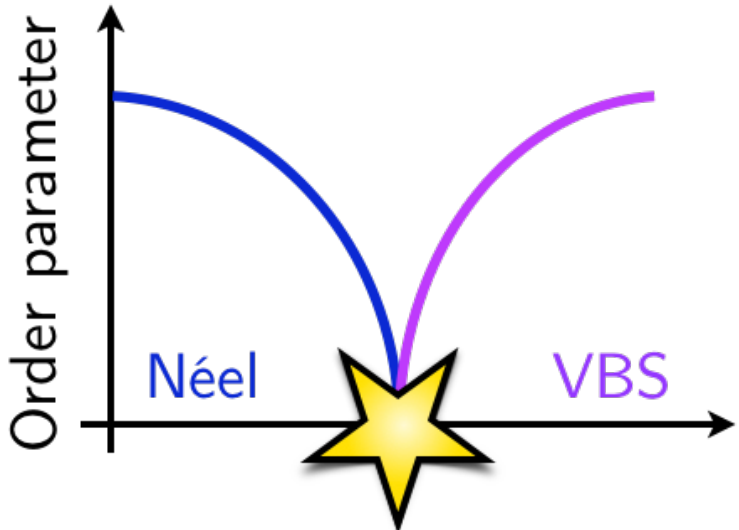
Landau

Deconfined quantum criticality

$$\text{---} = (|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle) / \sqrt{2}$$

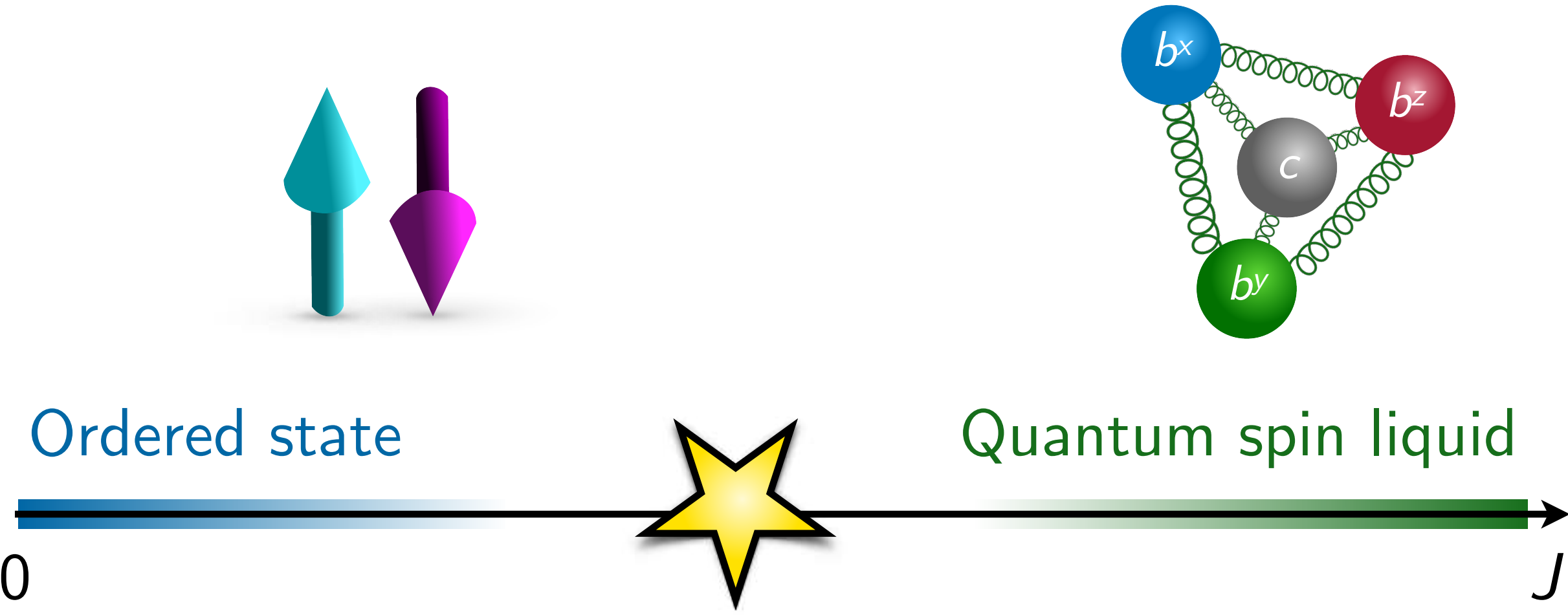


“Deconfined” quasiparticles B and \bar{B}

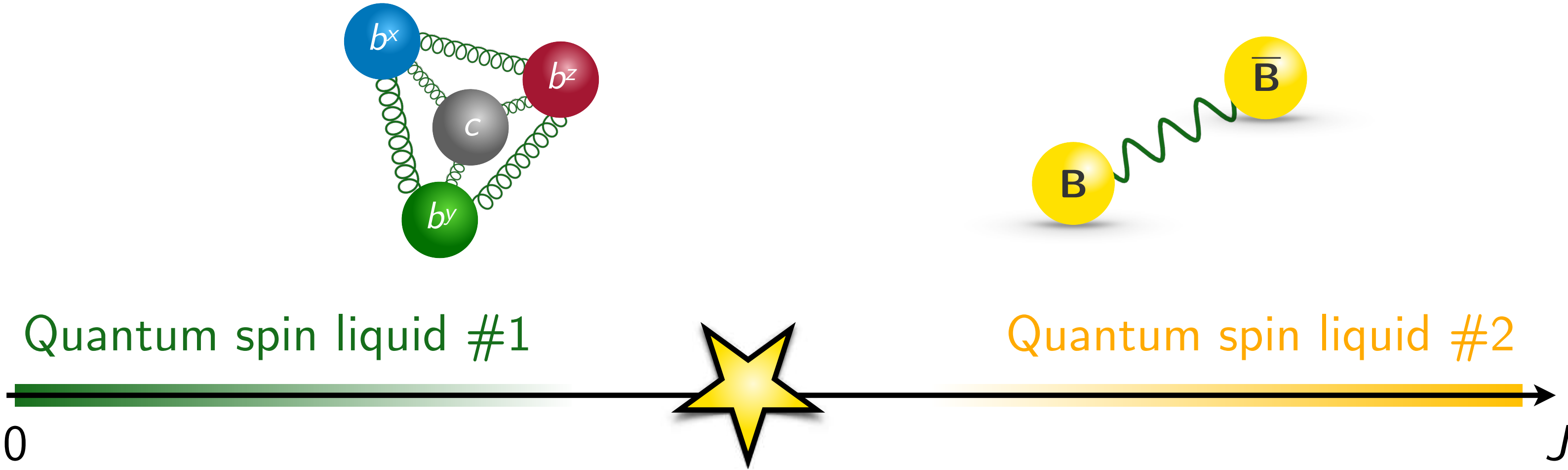


[Senthil *et al.*, Science '04]
 [Pujari, Damle, Alet, PRL '13]
 [Block, Melko, Kaul, PRL '13]
 [Shao, Guo, Sandvik, Science '16]
 ...

Fractionalized quantum criticality



[Isakov, Melko, Hastings, Science '12]
 [Assaad & Grover, PRX '16]
 [LJ, Wang, Scherer, Meng, Xu, PRB '20]
 ...



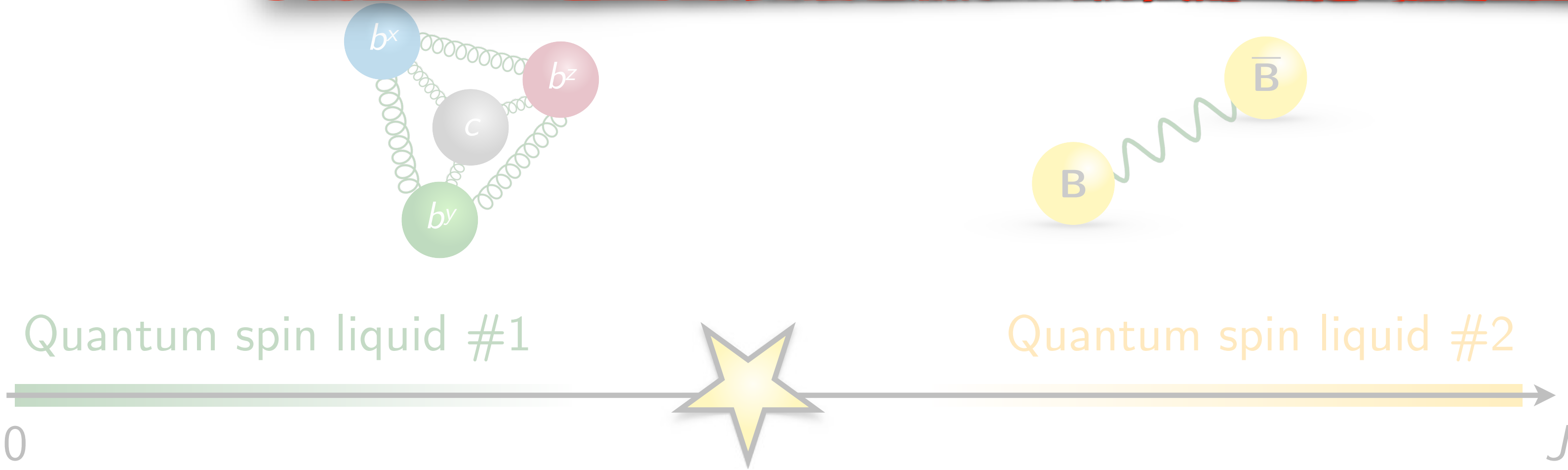
[Metlitski, Mross, Sachdev, Senthil, PRB '15]
 [LJ & He, PRB '17]
 [Boyack, Lin, Zerf, Rayyan, Maciejko, PRB '18]
 ...

Fractionalized quantum criticality



[Isakov, Melko, Hastings, Science '12]
[Assaad & Grover, PRX '16]
[LJ, Wang, Scherer, Meng, Xu, PRB '20]
...

Goal: Fractionalized transition in microscopic model



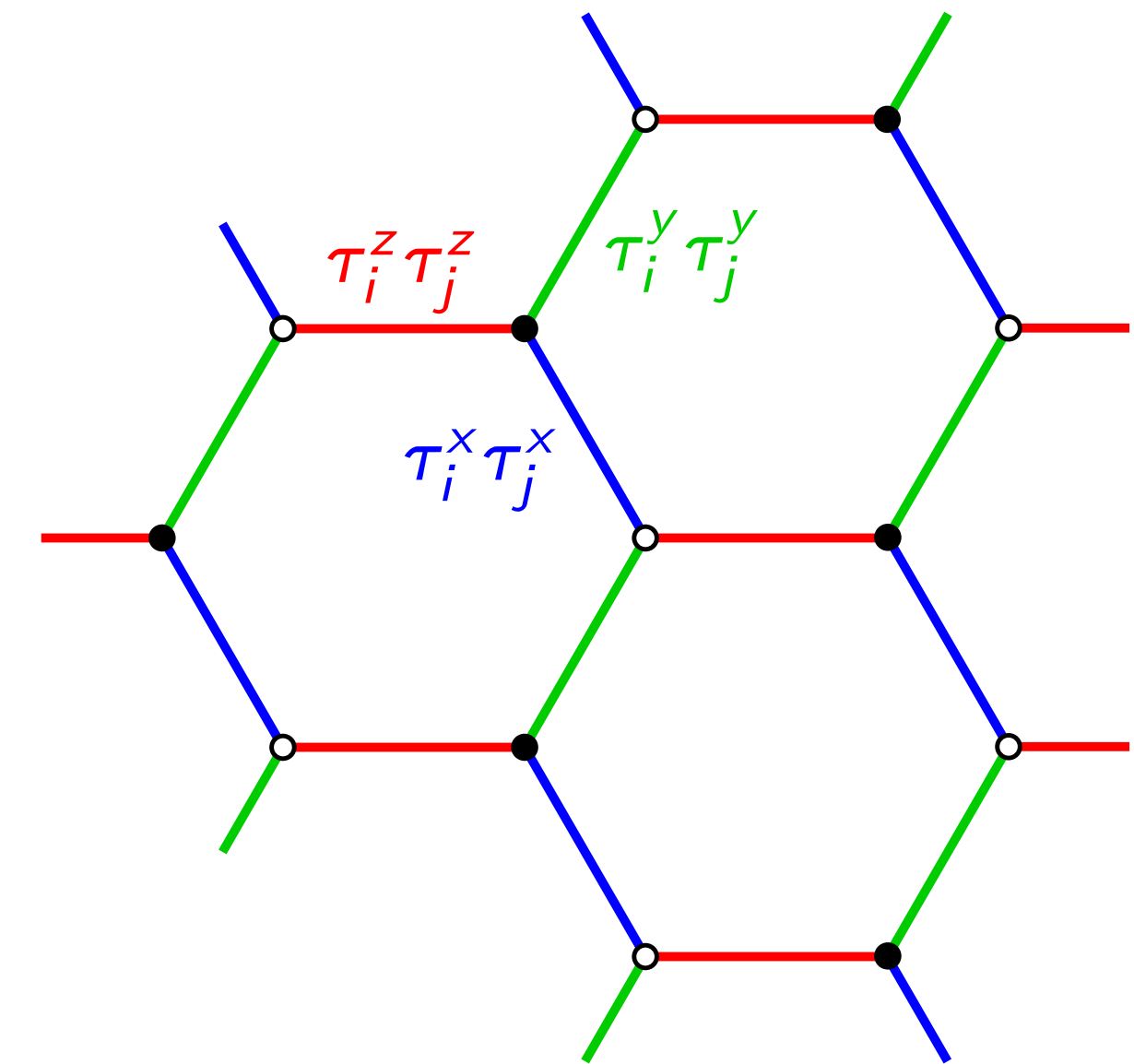
[Metlitski, Mross, Sachdev, Senthil, PRB '15]
[LJ & He, PRB '17]
[Boyack, Lin, Zerf, Rayyan, Maciejko, PRB '18]
...

Outline

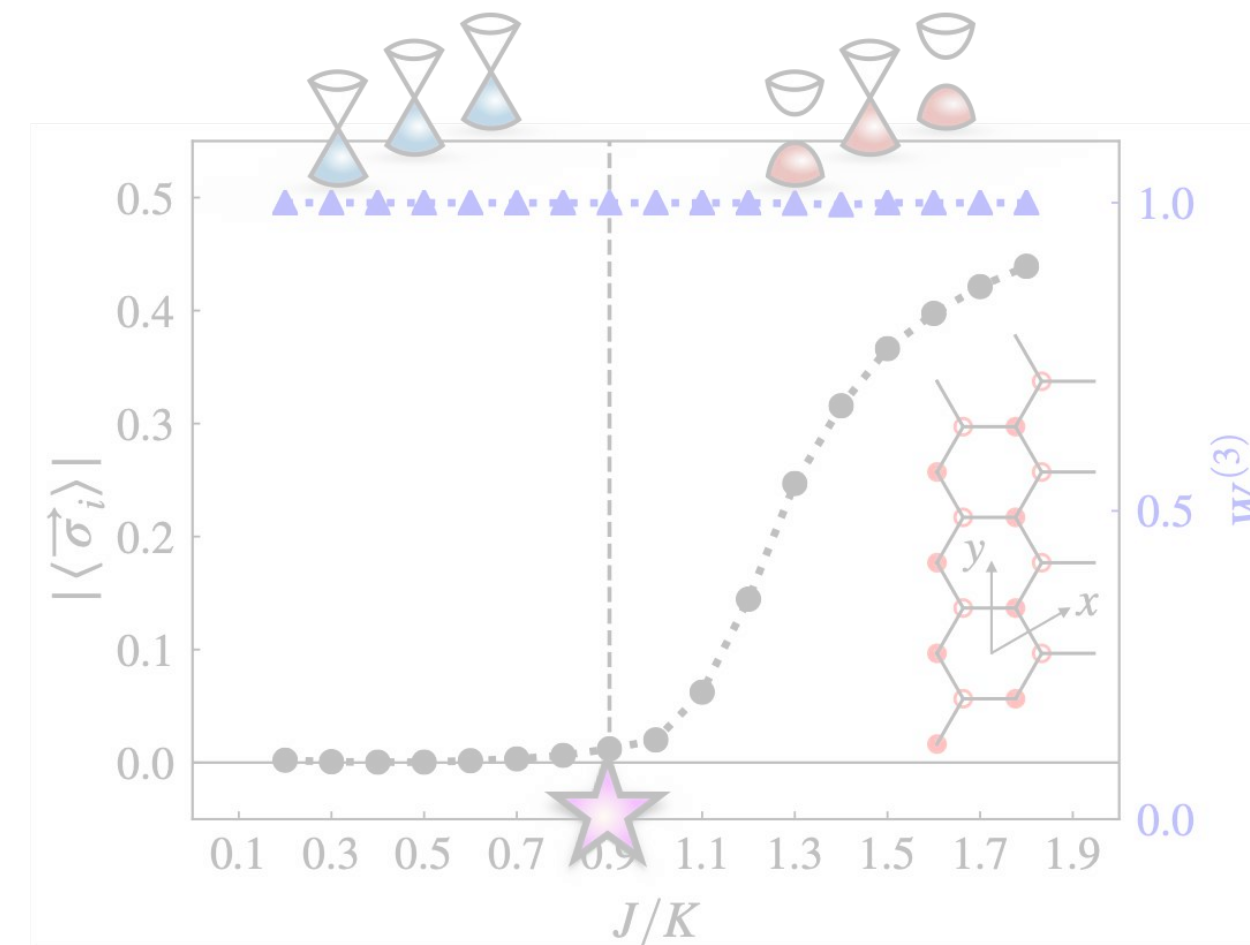
(1) Fractionalized quantum criticality



(2) From Kitaev to Kitaev-Kugel-Khomskii



(3) Kitaev-Heisenberg spin-orbital models

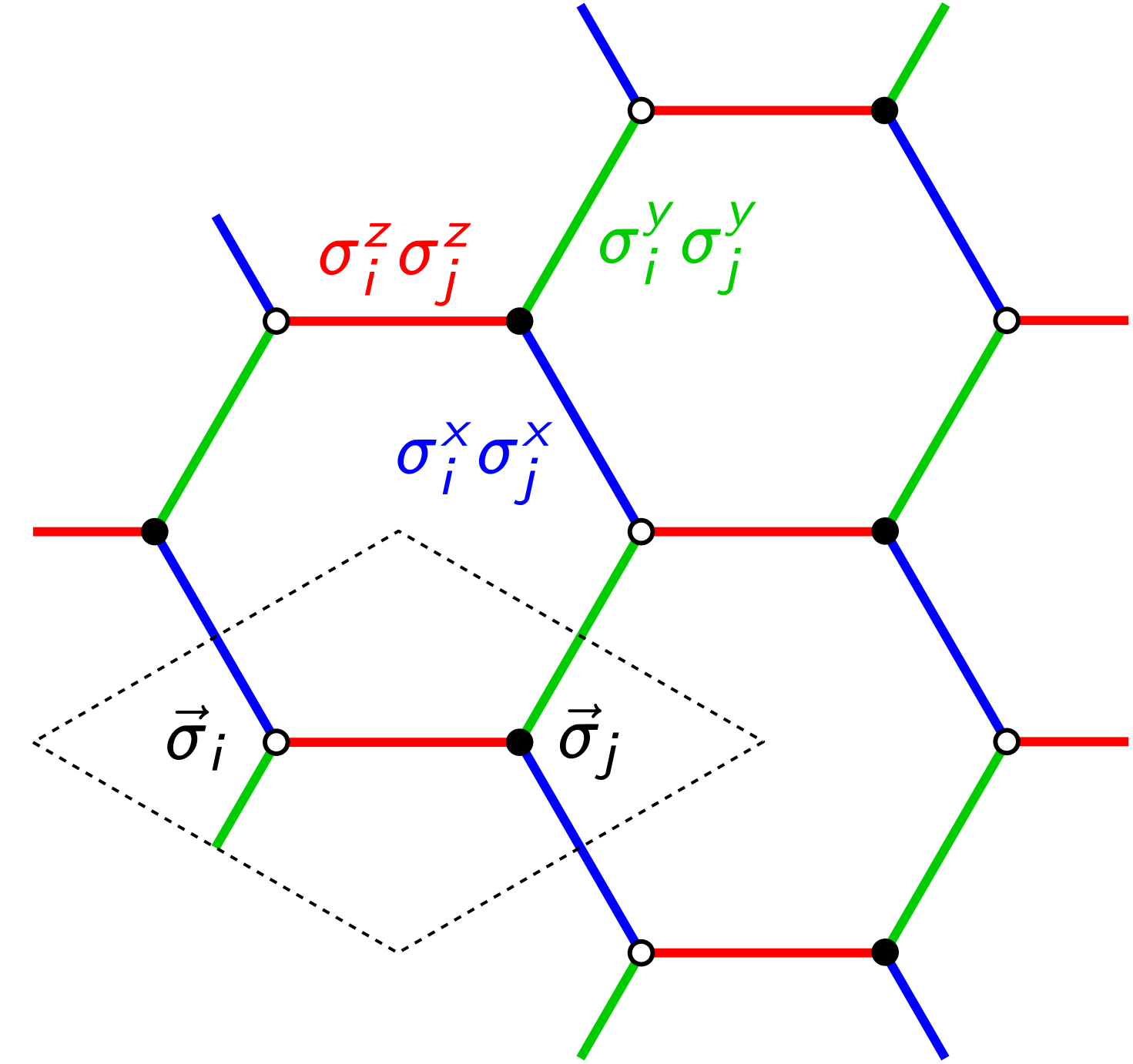


(4) Conclusions

Kitaev spin-1/2 model

Hamiltonian:

$$\mathcal{H} = K \sum_{\text{blue links}} \sigma_i^x \sigma_j^x + K \sum_{\text{green links}} \sigma_i^y \sigma_j^y + K \sum_{\text{red links}} \sigma_i^z \sigma_j^z$$



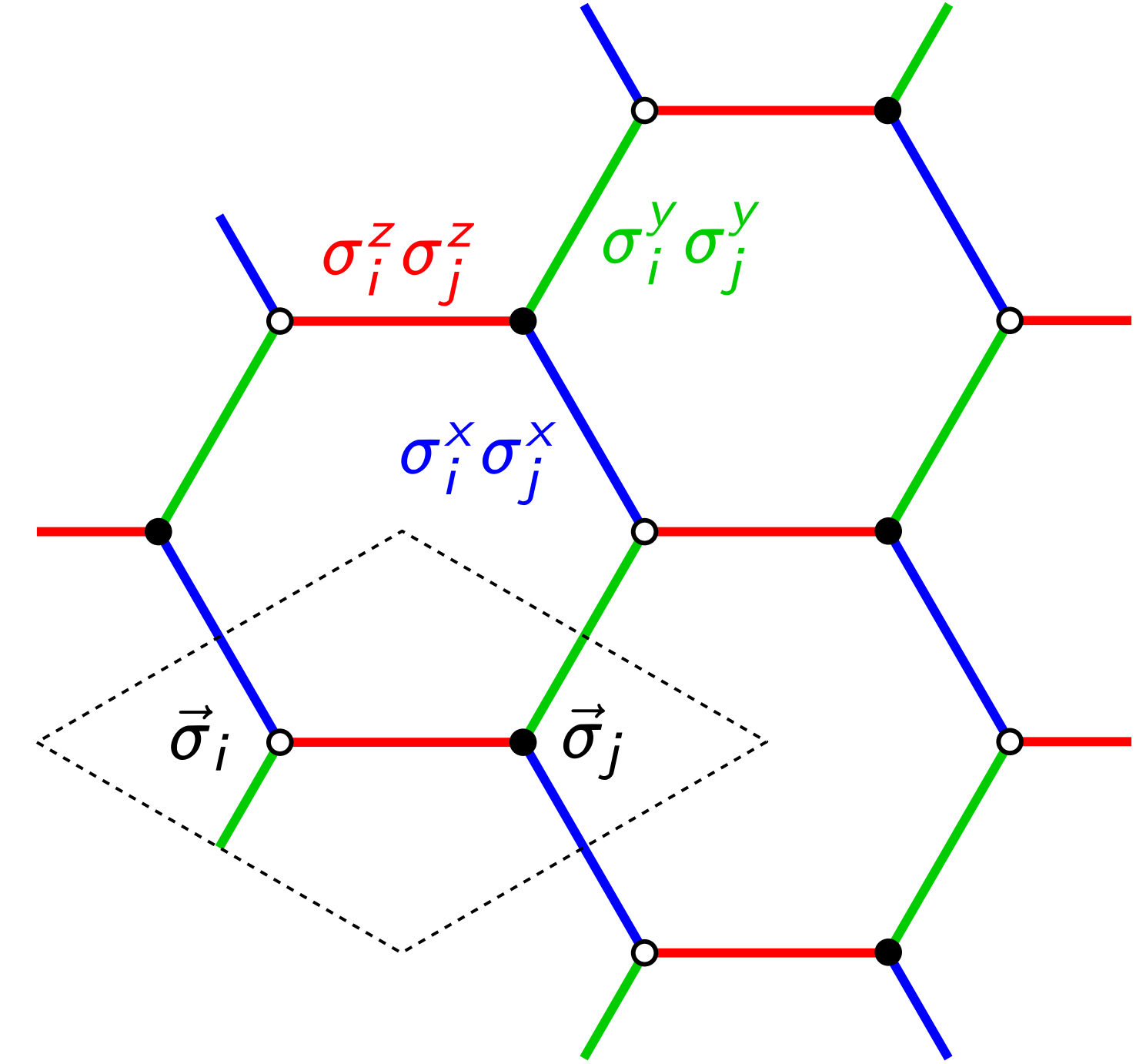
→ talks by N. Perkins, Y. Matsuda,
R. Valentí, ...

[Kitaev, Ann. Phys. '06]

Kitaev spin-1/2 model

Hamiltonian:

$$\mathcal{H} = K \sum_{\text{blue links}} \sigma_i^x \sigma_j^x + K \sum_{\text{green links}} \sigma_i^y \sigma_j^y + K \sum_{\text{red links}} \sigma_i^z \sigma_j^z$$

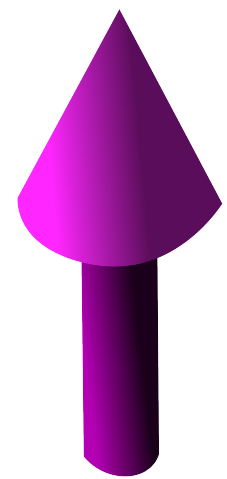


Majorana representation:

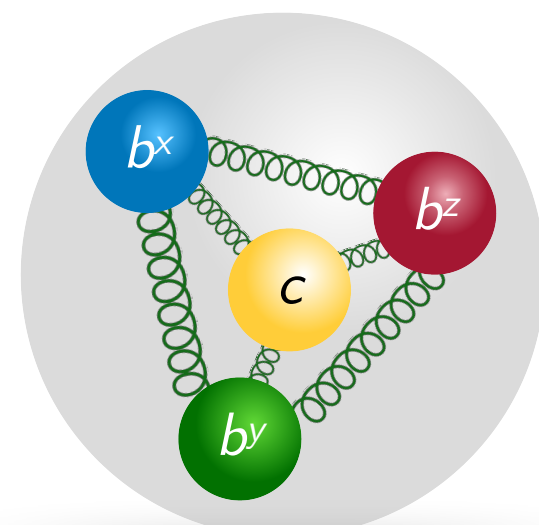
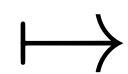
$$\sigma^x \mapsto \tilde{\sigma}^x = i b^x c$$

$$\sigma^y \mapsto \tilde{\sigma}^y = i b^y c$$

$$\sigma^z \mapsto \tilde{\sigma}^z = i b^z c$$



1 spin



4 Majoranas
with gauge constraint

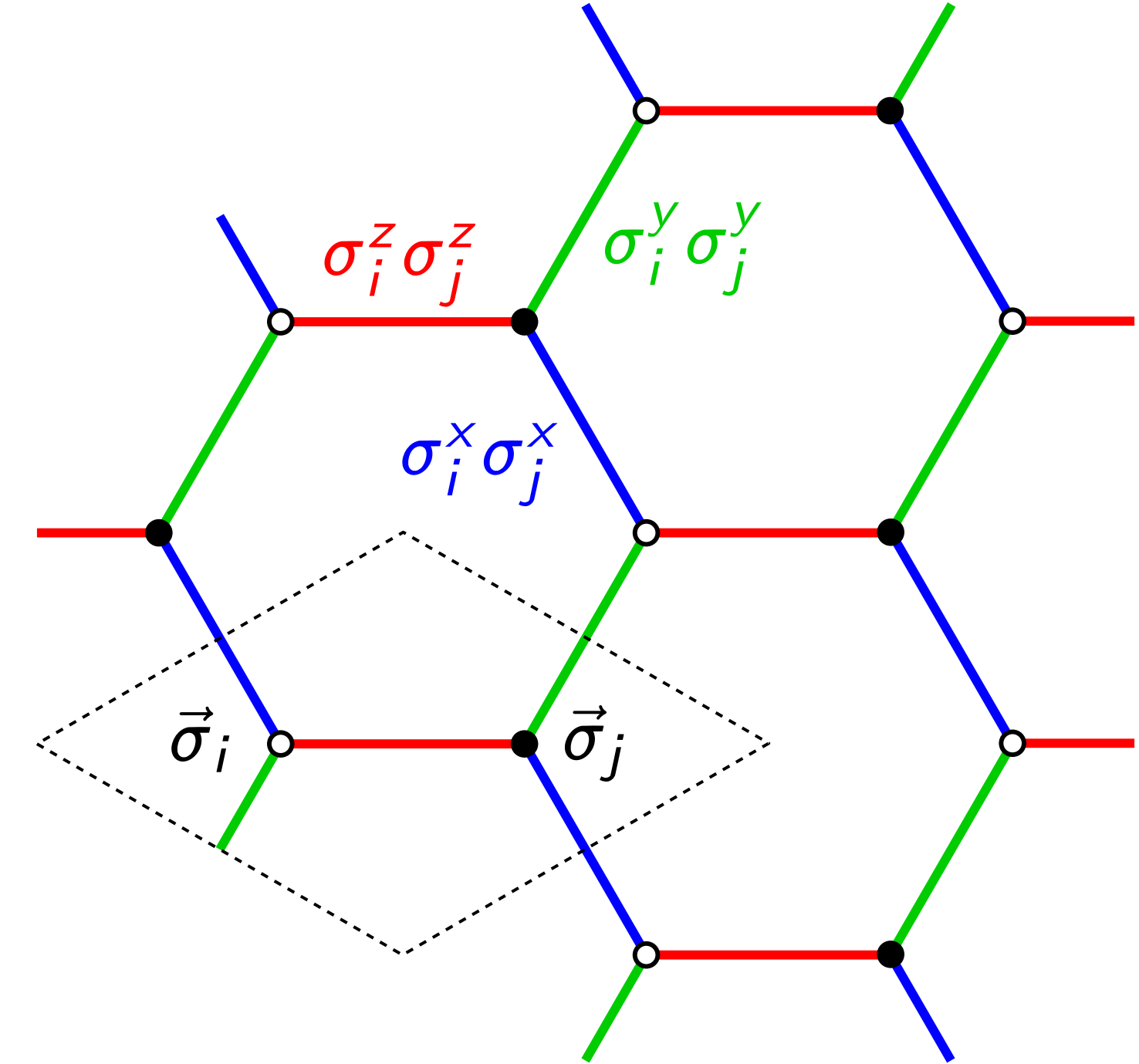
→ talks by N. Perkins, Y. Matsuda,
R. Valentí, ...

[Kitaev, Ann. Phys. '06]

Kitaev spin-1/2 model

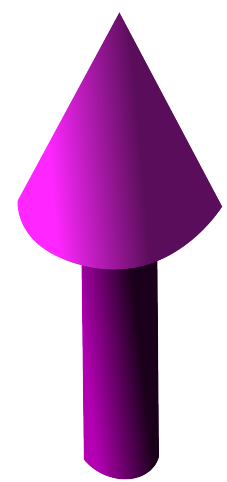
Hamiltonian:

$$\mathcal{H} = K \sum_{\text{blue links}} \sigma_i^x \sigma_j^x + K \sum_{\text{green links}} \sigma_i^y \sigma_j^y + K \sum_{\text{red links}} \sigma_i^z \sigma_j^z$$

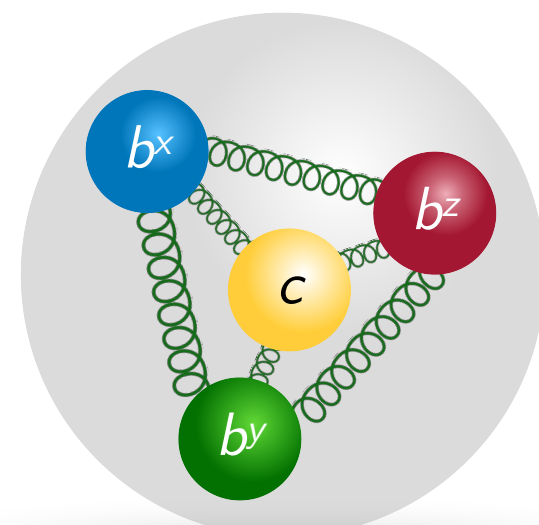
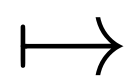


Majorana representation:

$$\begin{aligned} \sigma^x &\mapsto \tilde{\sigma}^x = i b^x c \\ \sigma^y &\mapsto \tilde{\sigma}^y = i b^y c \\ \sigma^z &\mapsto \tilde{\sigma}^z = i b^z c \end{aligned}$$



1 spin

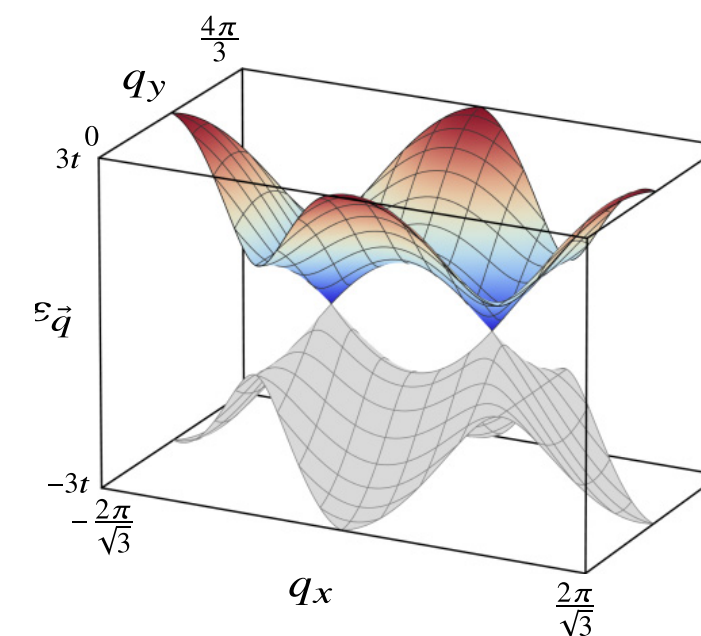


4 Majoranas
with gauge constraint

Fractionalization:

$$\mathcal{H} \mapsto \tilde{\mathcal{H}} = iK \sum_{\langle ij \rangle_\gamma} \underbrace{(i b_i^\gamma b_j^\gamma)}_{\equiv \hat{u}_{ij} = \hat{u}_{ij}^\dagger} c_i c_j$$

with $[\hat{u}_{ij}, \tilde{\mathcal{H}}] = 0 \Rightarrow$ static \mathbb{Z}_2 gauge field!



Ground-state flux pattern: $u \equiv 1$
[Lieb, PRL '94]

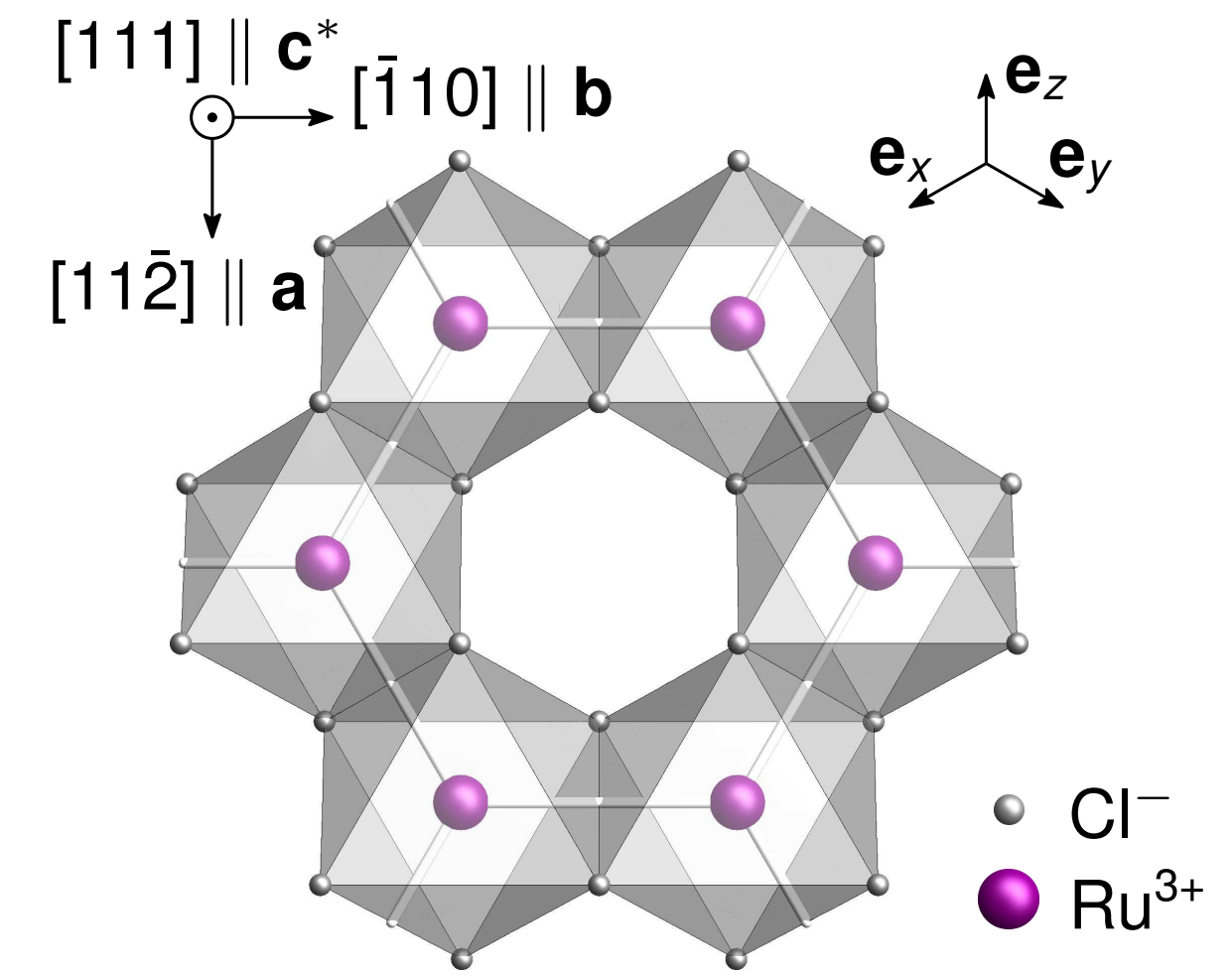
\rightarrow talks by N. Perkins, Y. Matsuda,
R. Valentí, ...

[Kitaev, Ann. Phys. '06]

Kitaev-Heisenberg spin-1/2 model

Hamiltonian:

$$\mathcal{H} = K \sum_{\langle ij \rangle_\gamma} \sigma_i^\gamma \sigma_j^\gamma + J \sum_{\langle ij \rangle} \vec{\sigma}_i \cdot \vec{\sigma}_j$$



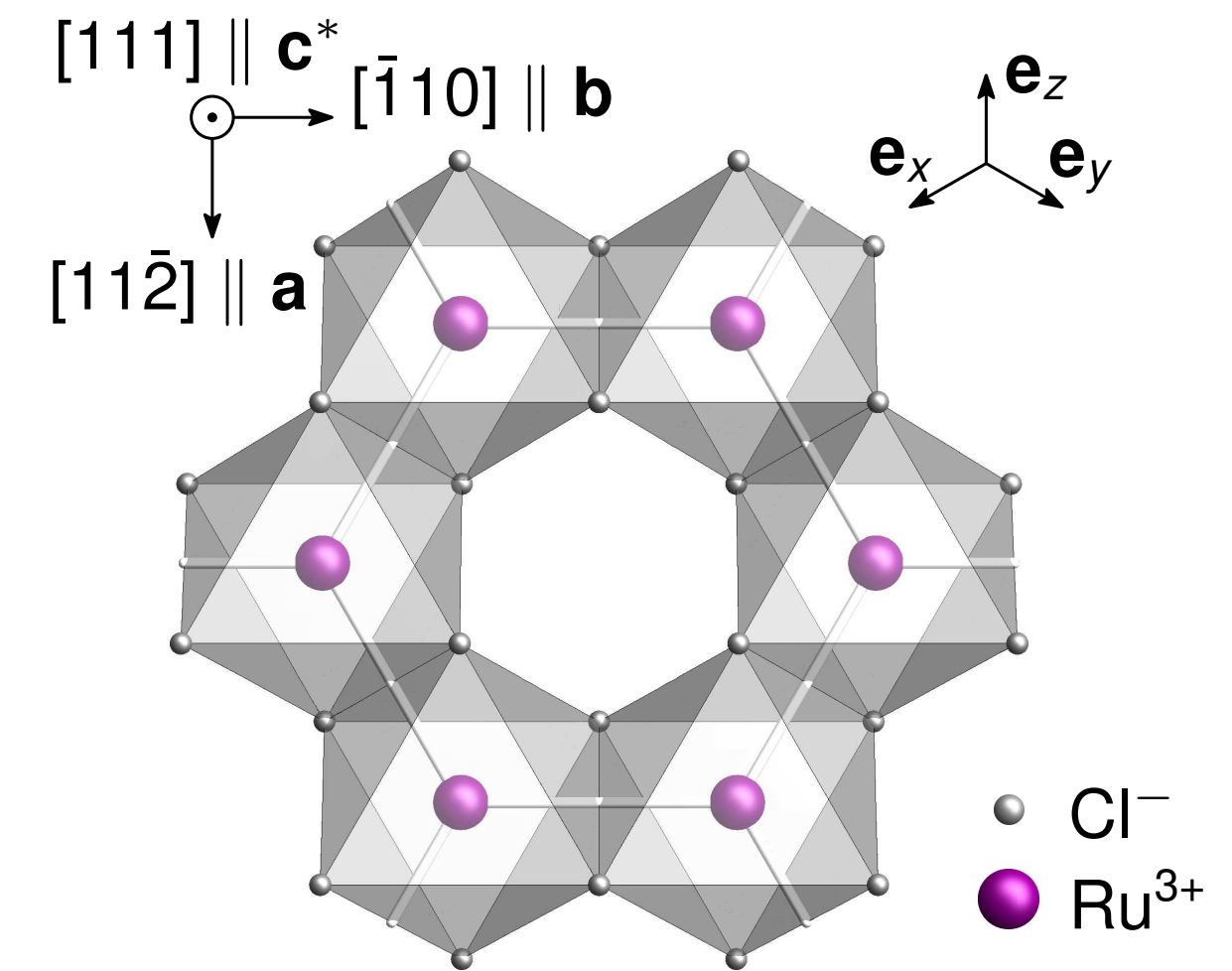
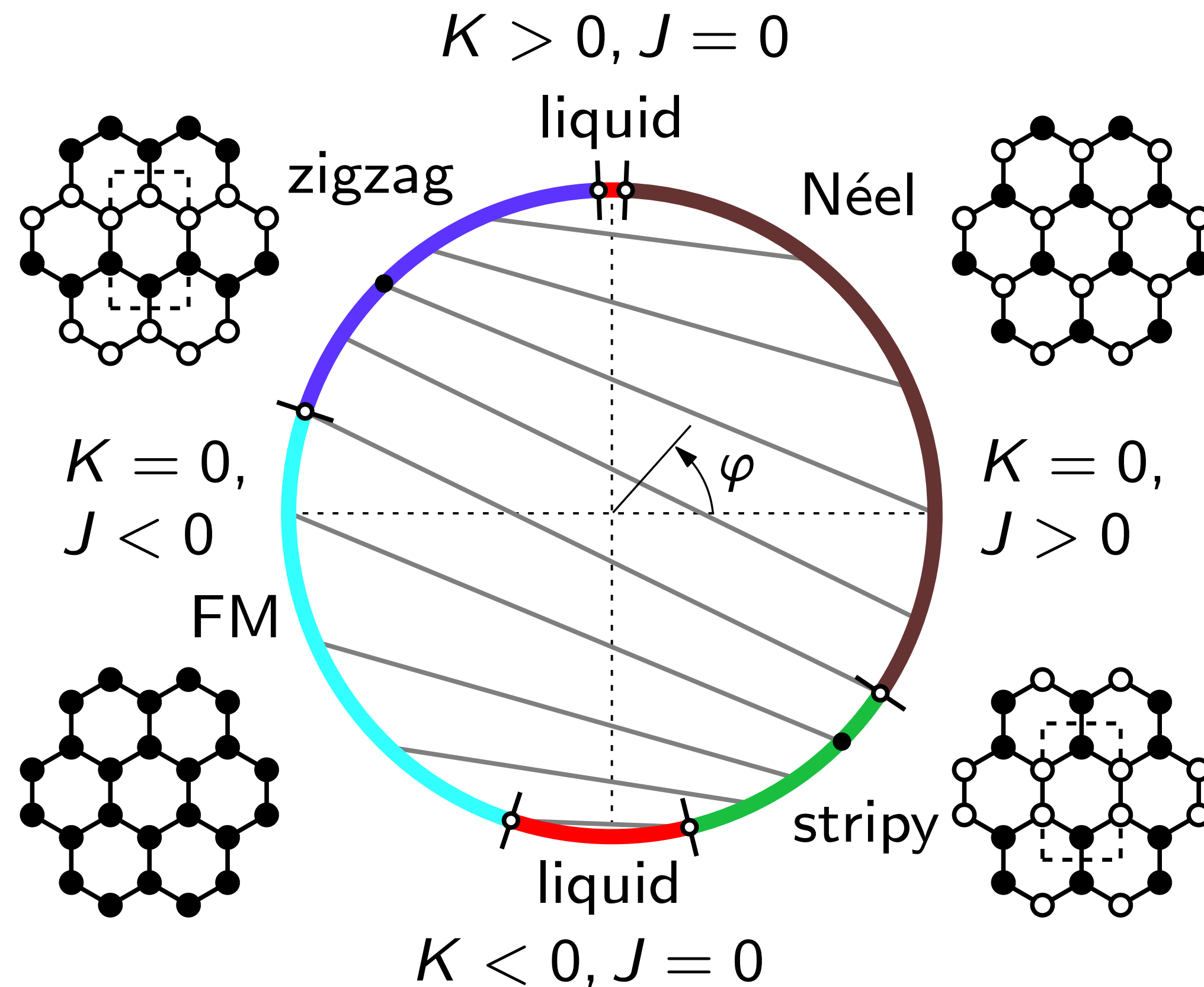
... possible relevance to α -RuCl₃, Na₂IrO₃, Na₂Co₂TeO₆, ...

Kitaev-Heisenberg spin-1/2 model

Hamiltonian:

$$\mathcal{H} = K \sum_{\langle ij \rangle_\gamma} \sigma_i^\gamma \sigma_j^\gamma + J \sum_{\langle ij \rangle} \vec{\sigma}_i \cdot \vec{\sigma}_j$$

Phase diagram:



... possible relevance to α -RuCl₃, Na₂IrO₃, Na₂Co₂TeO₆, ...

$$J = A \cos \varphi$$

$$K = 2A \sin \varphi$$

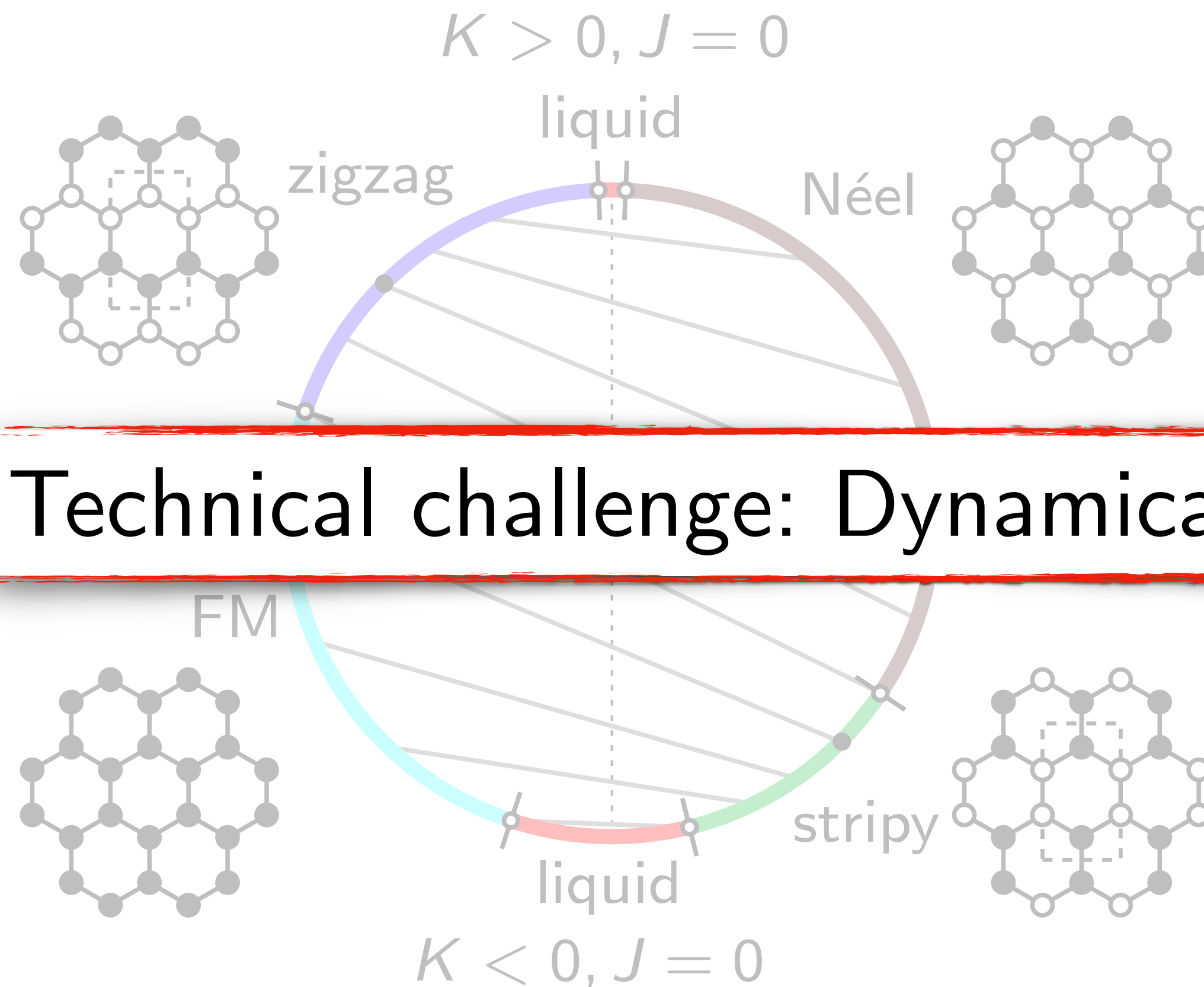
... from 24-site ED: [Chaloupka, Jackeli, Khaliullin, PRL '13]

Kitaev-Heisenberg spin-1/2 model

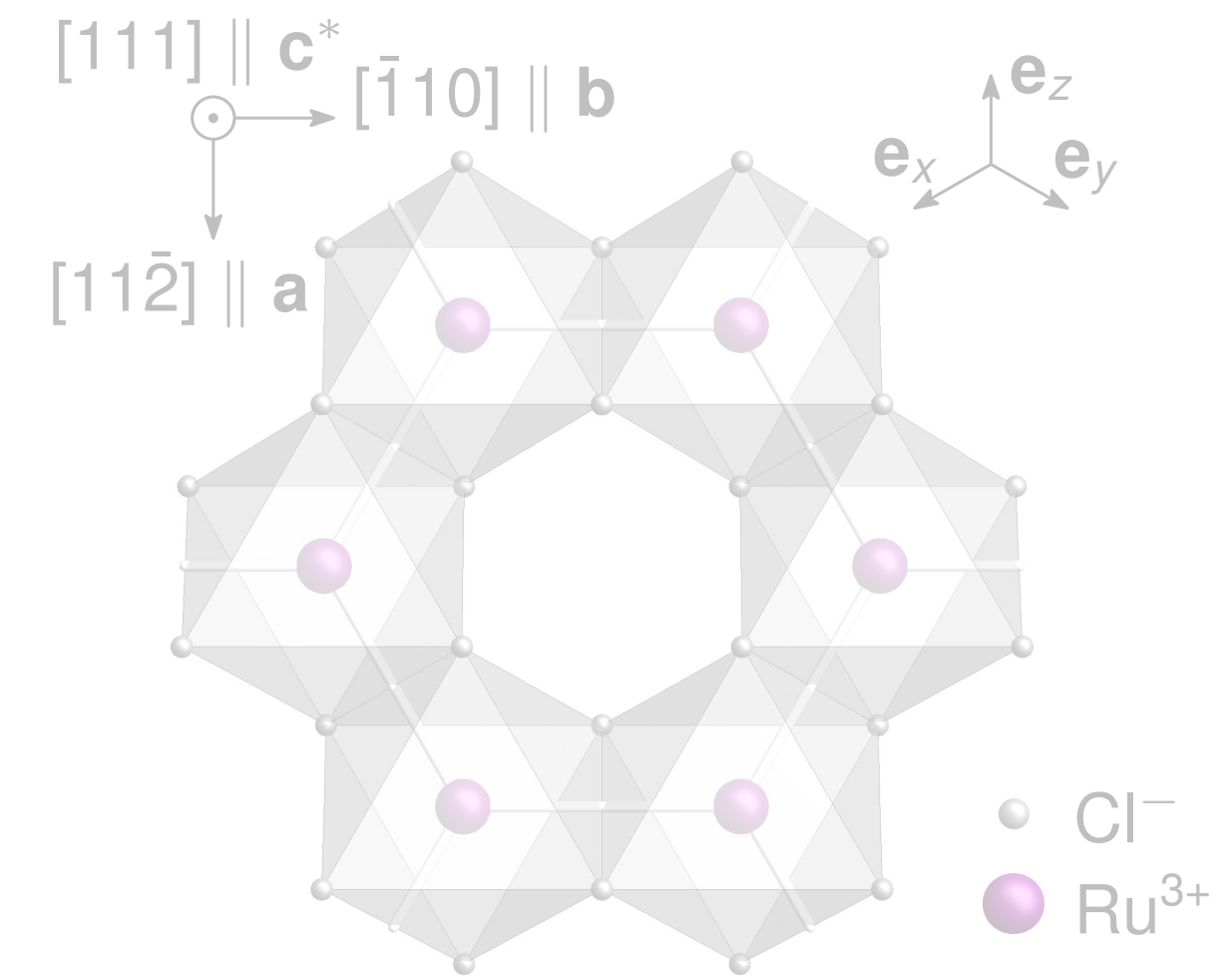
Hamiltonian:

$$\mathcal{H} = K \sum_{\langle ij \rangle_\gamma} \sigma_i^\gamma \sigma_j^\gamma + J \sum_{\langle ij \rangle} \vec{\sigma}_i \cdot \vec{\sigma}_j$$

Phase diagram:



Technical challenge: Dynamical \mathbb{Z}_2 gauge field!



... possible relevance to $\alpha\text{-RuCl}_3$, Na_2IrO_3 , $\text{Na}_2\text{Co}_2\text{TeO}_6$, ...

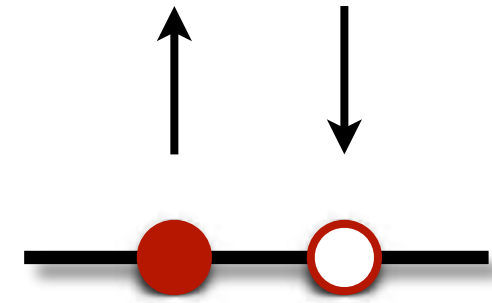
... no sign-problem-free QMC available: [Sato & Assaad, PRB '21]

→ talk by **F. Assaad**

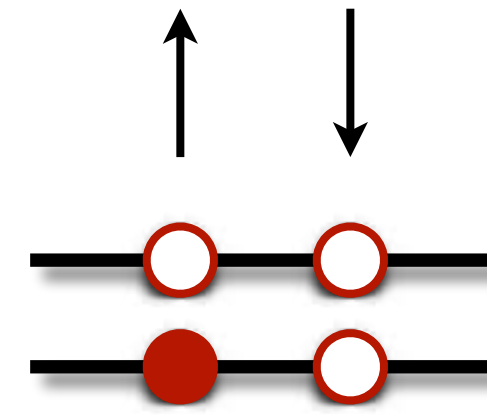
... from 24-site ED: [Chaloupka, Jackeli, Khaliullin, PRL '13]

Kitaev spin-orbital models

Spin-orbital generalization:

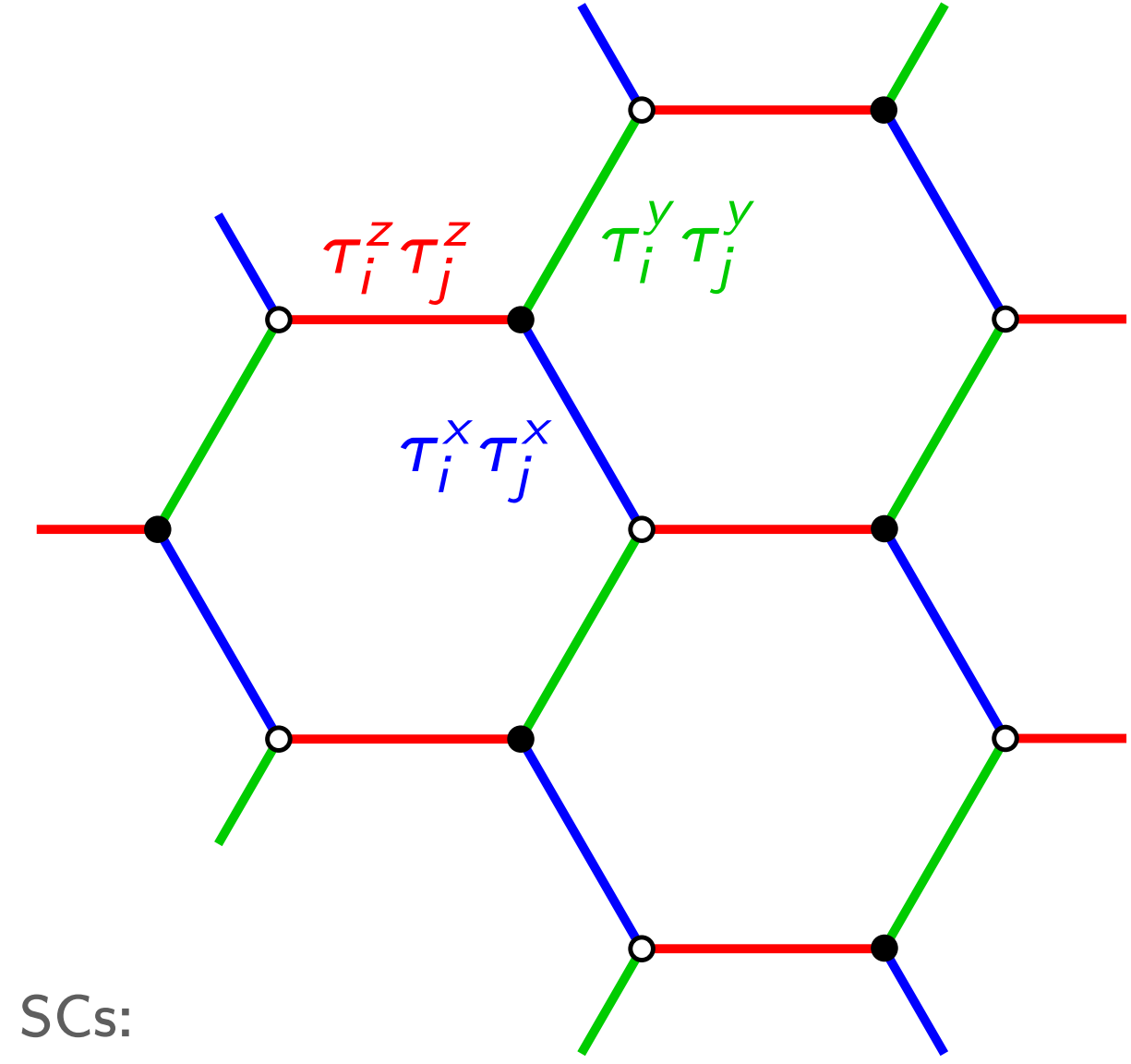


$$\sigma^\alpha \quad 2 \times 2$$



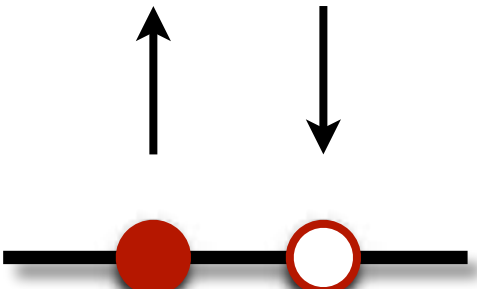
$$\sigma^\alpha \otimes \tau^\beta = \gamma^i \quad 4 \times 4$$

... can realize all 16 topological SCs:
[Chulliparambil, *et int.*, LJ, Tu, PRB '20]

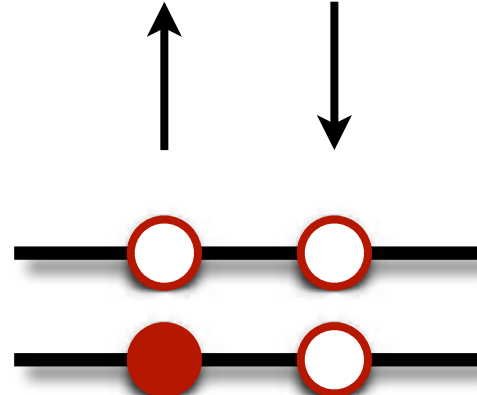


Kitaev spin-orbital models

Spin-orbital generalization:

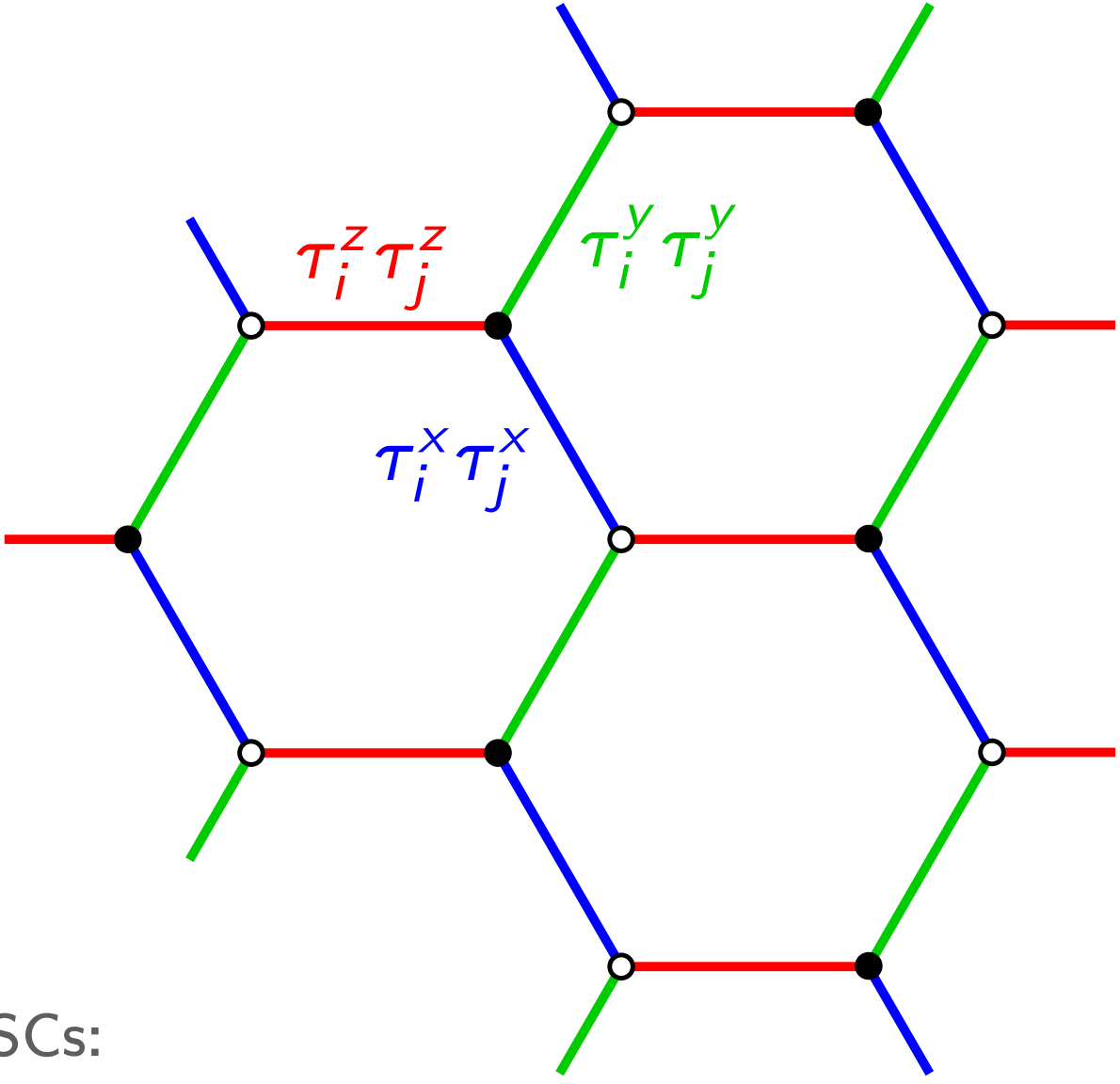


$$\sigma^\alpha \quad 2 \times 2$$



$$\sigma^\alpha \otimes \tau^\beta = \gamma^i \quad 4 \times 4$$

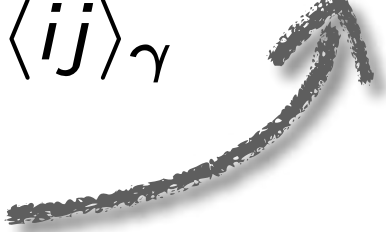
... can realize all 16 topological SCs:
 [Chulliparambil, *et int.*, LJ, Tu, PRB '20]



Hamiltonian:

$$\mathcal{H} = K \sum_{\langle ij \rangle_\gamma} \vec{\sigma}_i \cdot \vec{\sigma}_j \otimes \tau_i^\gamma \tau_j^\gamma$$

Heisenberg spin

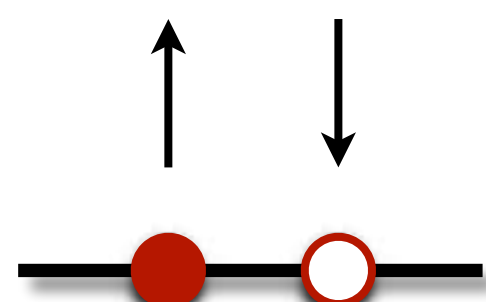


Kitaev orbital

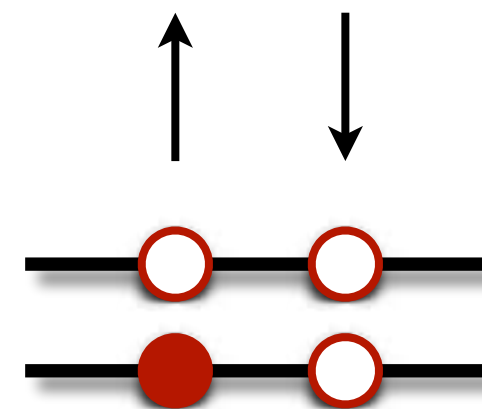


Kitaev spin-orbital models

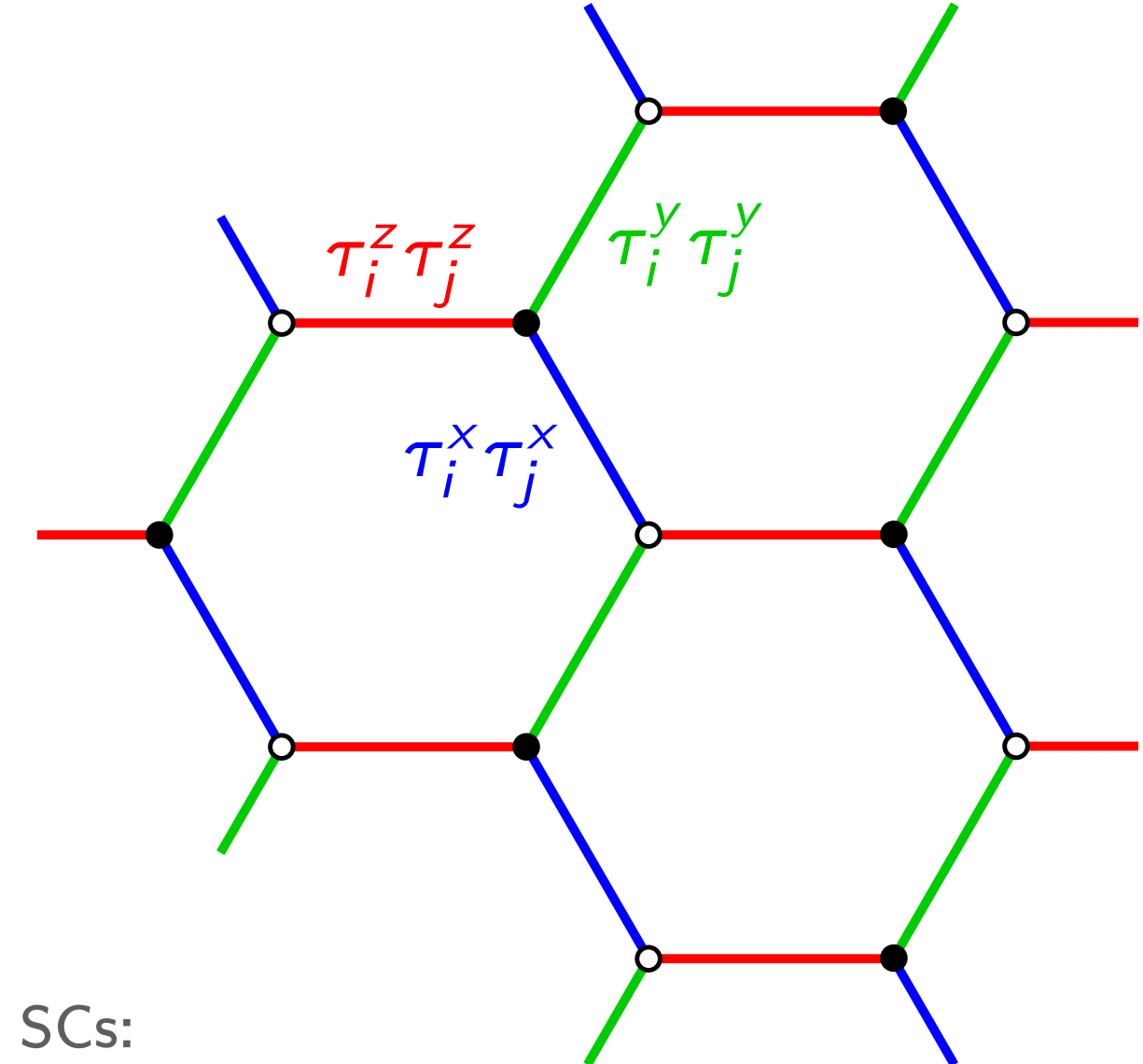
Spin-orbital generalization:



$$\sigma^\alpha \quad 2 \times 2$$



$$\sigma^\alpha \otimes \tau^\beta = \gamma^i \quad 4 \times 4$$



... can realize all 16 topological SCs:
[Chulliparambil, *et int.*, LJ, Tu, PRB '20]

Hamiltonian:

$$\mathcal{H} = K \sum_{\langle ij \rangle_\gamma} \vec{\sigma}_i \cdot \vec{\sigma}_j \otimes \tau_i^\gamma \tau_j^\gamma$$

Heisenberg spin

Kitaev orbital

Majorana representation:

$$\sigma^y \otimes \tau^x = ib^1 c^x$$

$$\sigma^y \otimes \tau^y = ib^2 c^x$$

$$\sigma^y \otimes \tau^z = ib^3 c^x$$

$$\sigma^x \otimes \mathbb{1} = ic^y c^x$$

$$\sigma^z \otimes \mathbb{1} = ic^z c^x$$

Fractionalization:

$$\mathcal{H} \mapsto \tilde{\mathcal{H}} = iK \sum_{\langle ij \rangle_\gamma} \hat{u}_{ij} c_i^\top c_j$$

$$\text{with } [\hat{u}_{ij}, \tilde{\mathcal{H}}] = 0$$

$$c \equiv \begin{pmatrix} c^x \\ c^y \\ c^z \end{pmatrix}$$

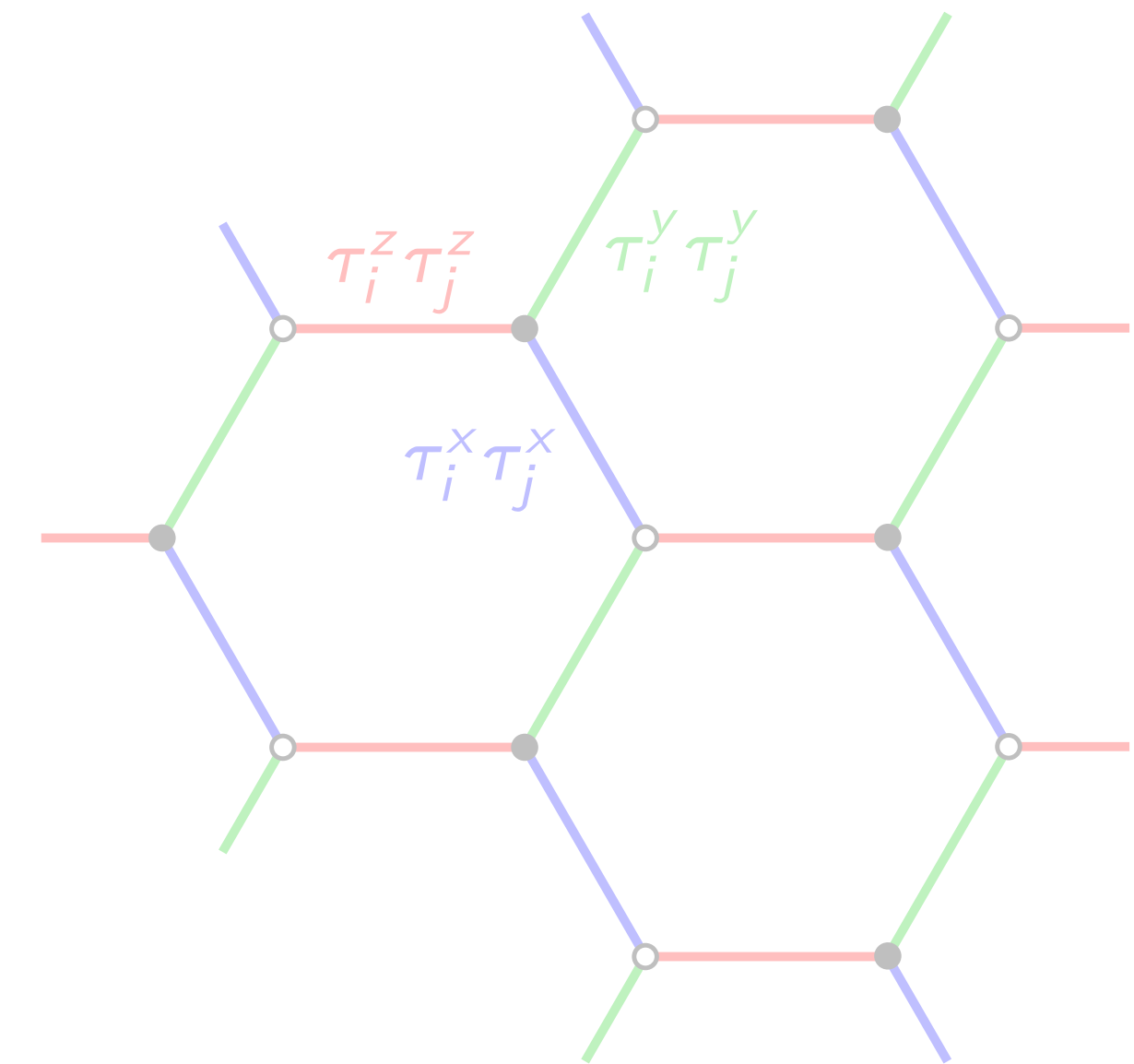
... cf. also [Yao & Lee, PRL '11]

Outline

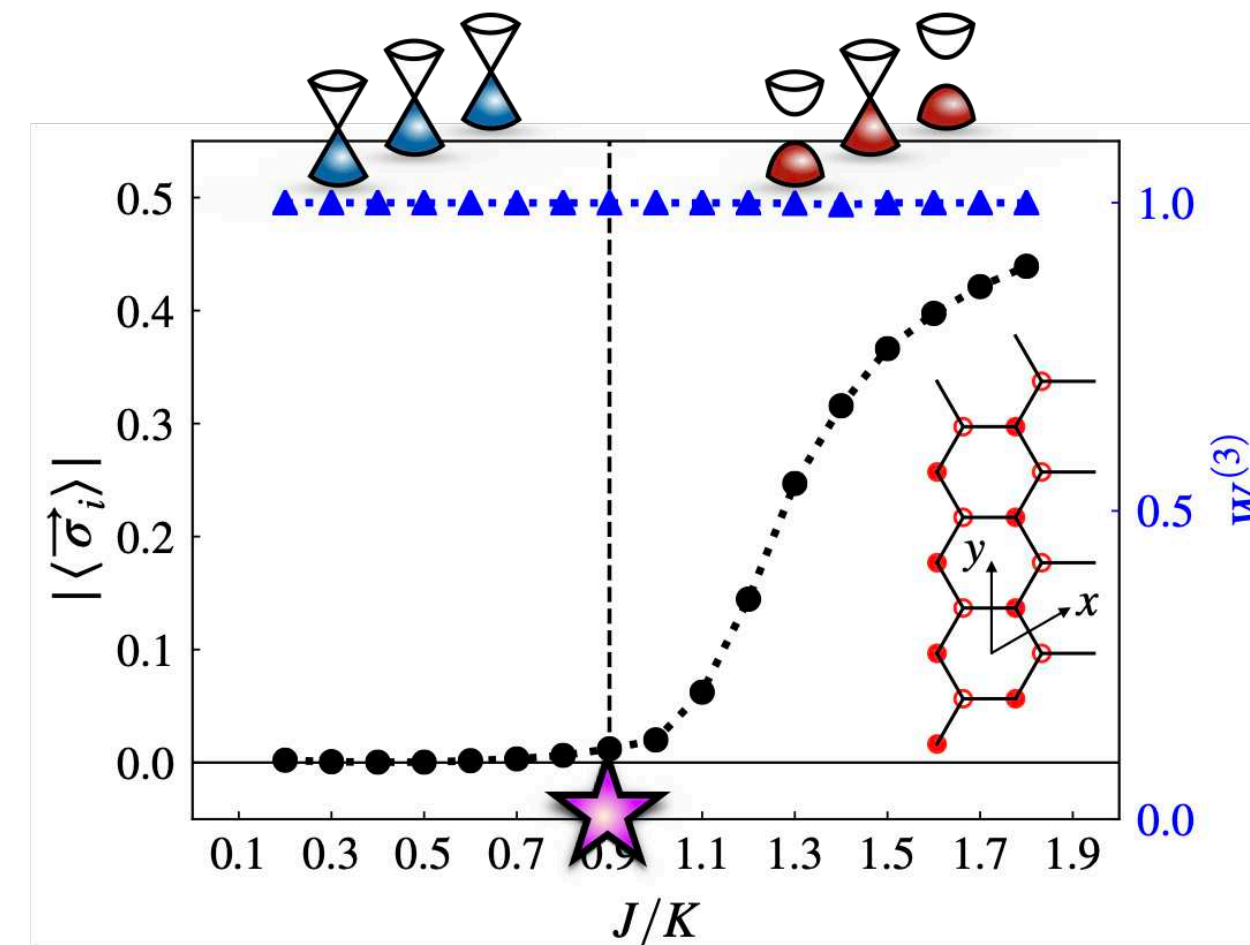
(1) Fractionalized quantum criticality



(2) From Kitaev to Kitaev-Kugel-Khomskii



(3) Kitaev-Heisenberg spin-orbital models



(4) Conclusions

Kitaev-Heisenberg spin-orbital model

Hamiltonian:

$$\mathcal{H} = K \sum_{\langle ij \rangle_\gamma} \vec{\sigma}_i \cdot \vec{\sigma}_j \otimes \tau_i^\gamma \tau_j^\gamma + J \sum_{\langle ij \rangle} \underbrace{\vec{\sigma}_i \cdot \vec{\sigma}_j \otimes \mathbb{1}_i \mathbb{1}_j}_{\substack{\text{spin-1 matrices} \\ \downarrow}} \mapsto \frac{1}{4} (c_i^\top \vec{L} c_i) \cdot (c_j \vec{L} c_j)$$

with $[\hat{u}_{ij}, \mathcal{H}] = 0$ still static!

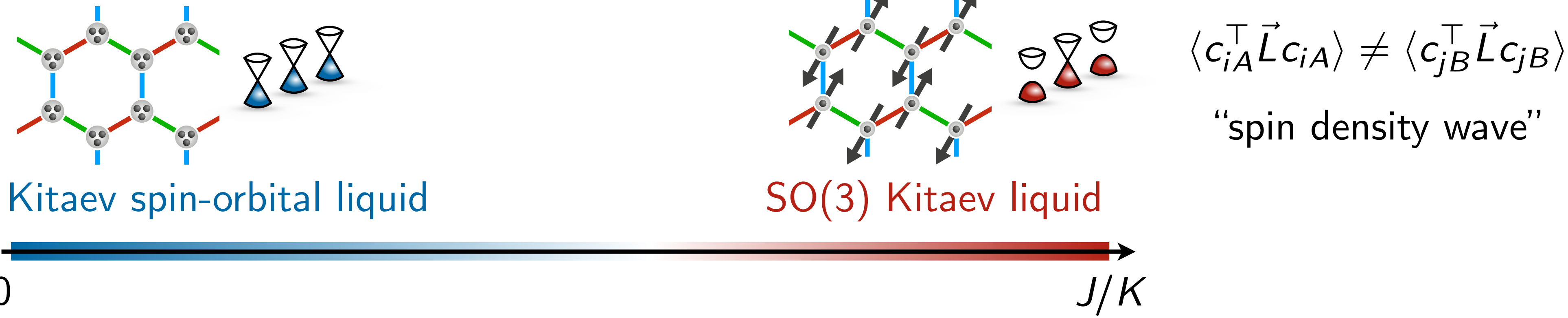
Kitaev-Heisenberg spin-orbital model

Hamiltonian:

$$\mathcal{H} = K \sum_{\langle ij \rangle_\gamma} \vec{\sigma}_i \cdot \vec{\sigma}_j \otimes \tau_i^\gamma \tau_j^\gamma + J \sum_{\langle ij \rangle} \underbrace{\vec{\sigma}_i \cdot \vec{\sigma}_j \otimes \mathbb{1}_i \mathbb{1}_j}_{\substack{\text{spin-1 matrices} \\ \mapsto \frac{1}{4} (c_i^\top \vec{L} c_i) \cdot (c_j \vec{L} c_j)}} + J \sum_{\langle ij \rangle} \vec{\sigma}_i \cdot \vec{\sigma}_j \otimes \mathbb{1}_i \mathbb{1}_j$$

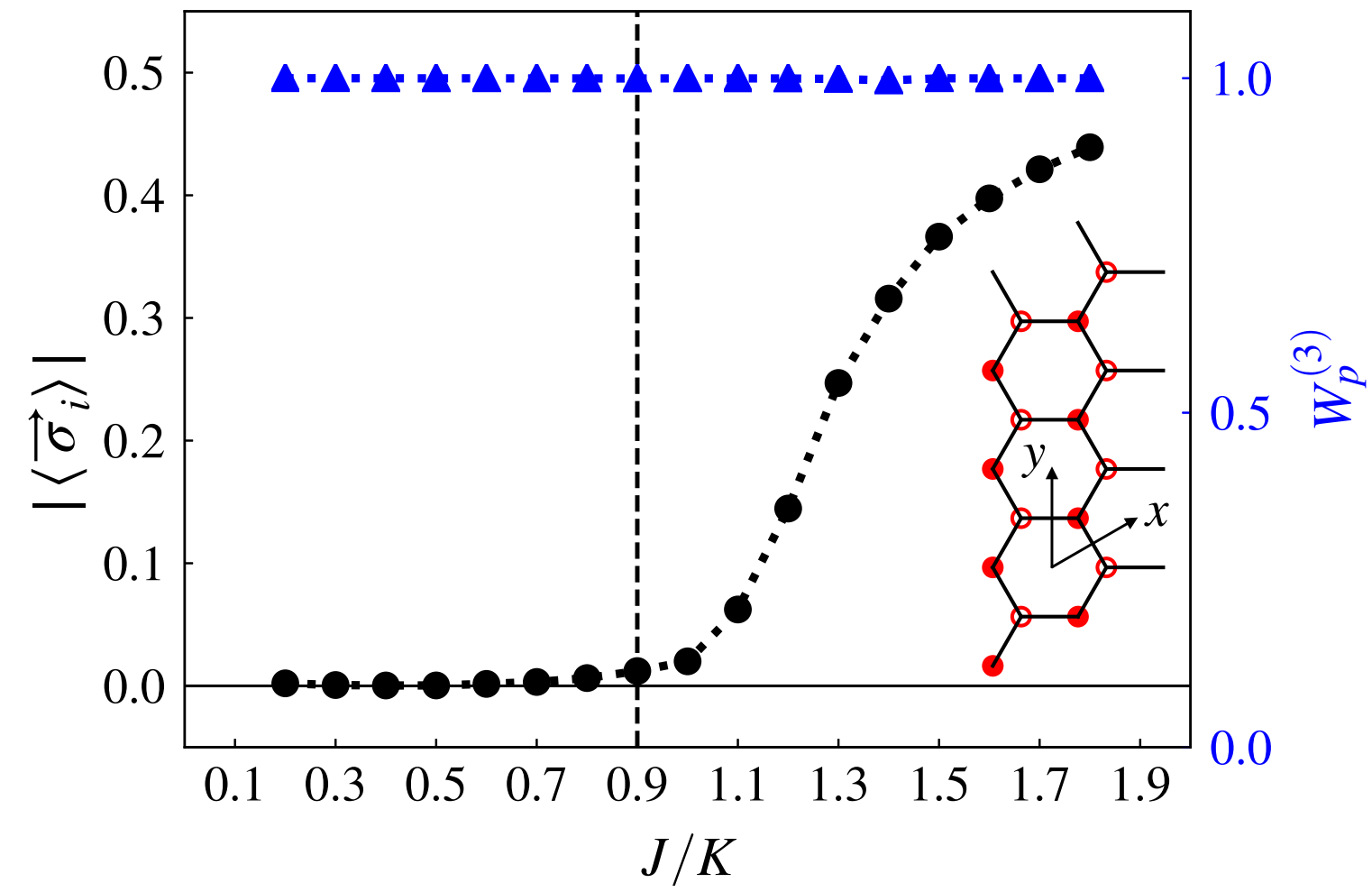
with $[\hat{u}_{ij}, \mathcal{H}] = 0$ still static!

Phase diagram:



Gross-Neveu-SO(3)* transition

iDMRG:

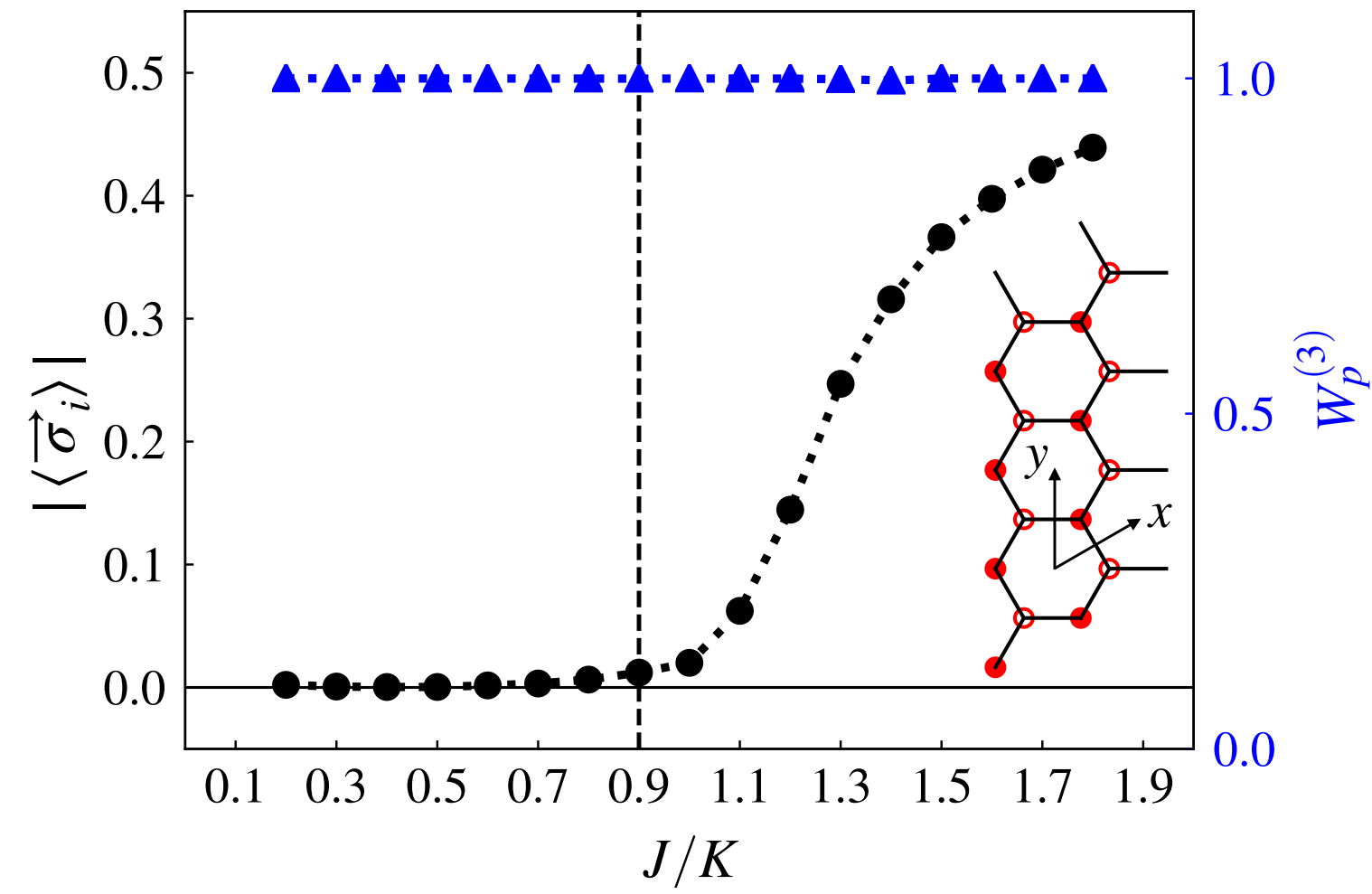


... on cylinder with $L_y = 4$ unit cells

[Seifert, Dong, Chulliparambil, Vojta, Tu, LJ, PRL '20]

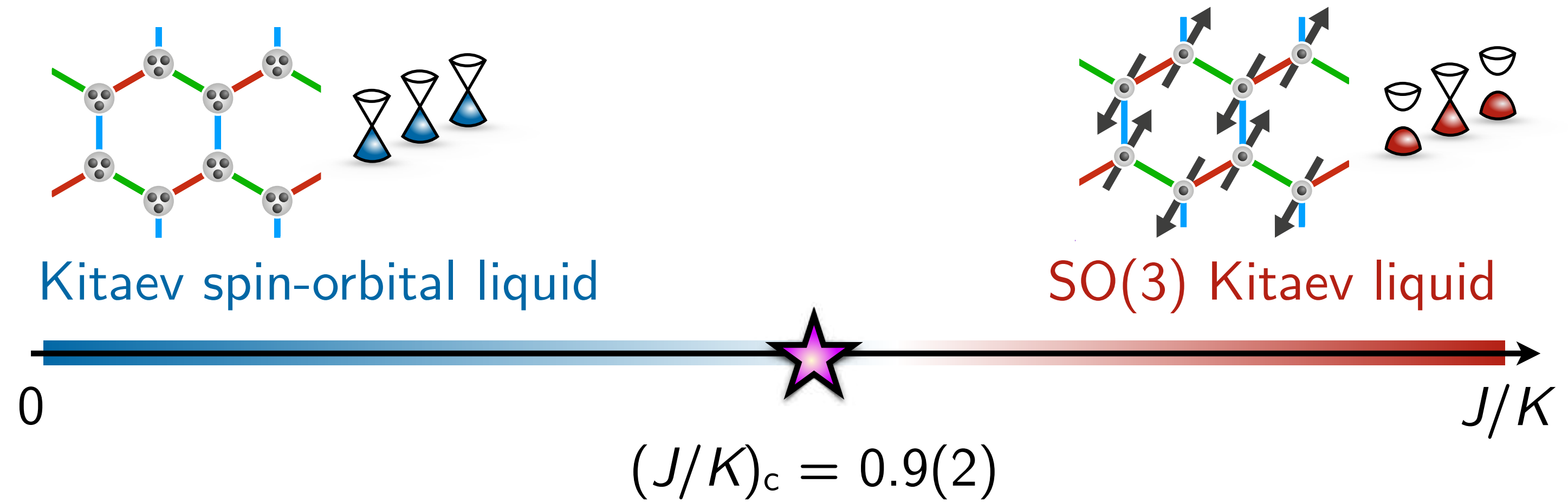
Gross-Neveu-SO(3)* transition

iDMRG:



... on cylinder with $L_y = 4$ unit cells

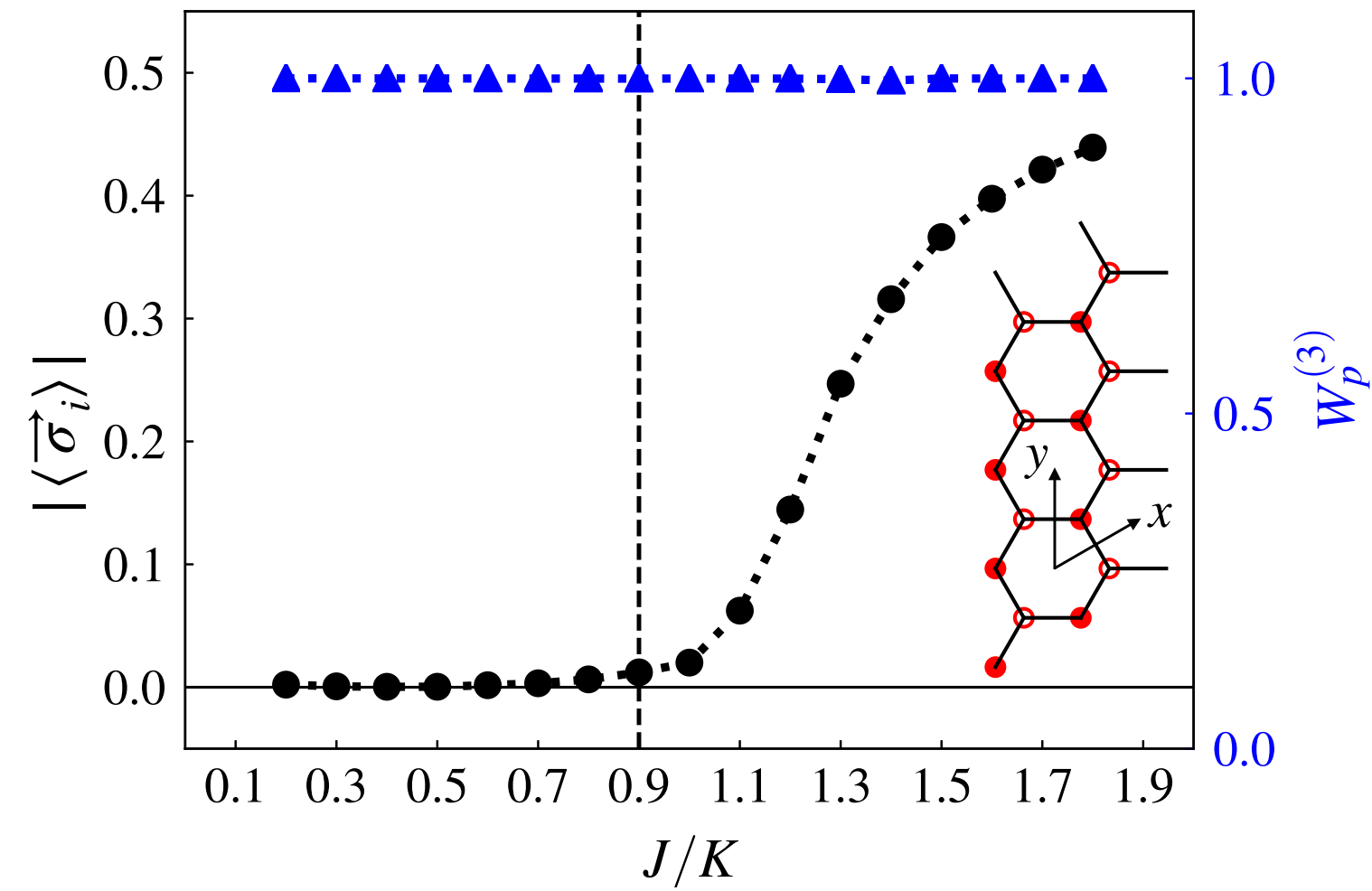
Phase diagram:



[Seifert, Dong, Chulliparambil, Vojta, Tu, LJ, PRL '20]

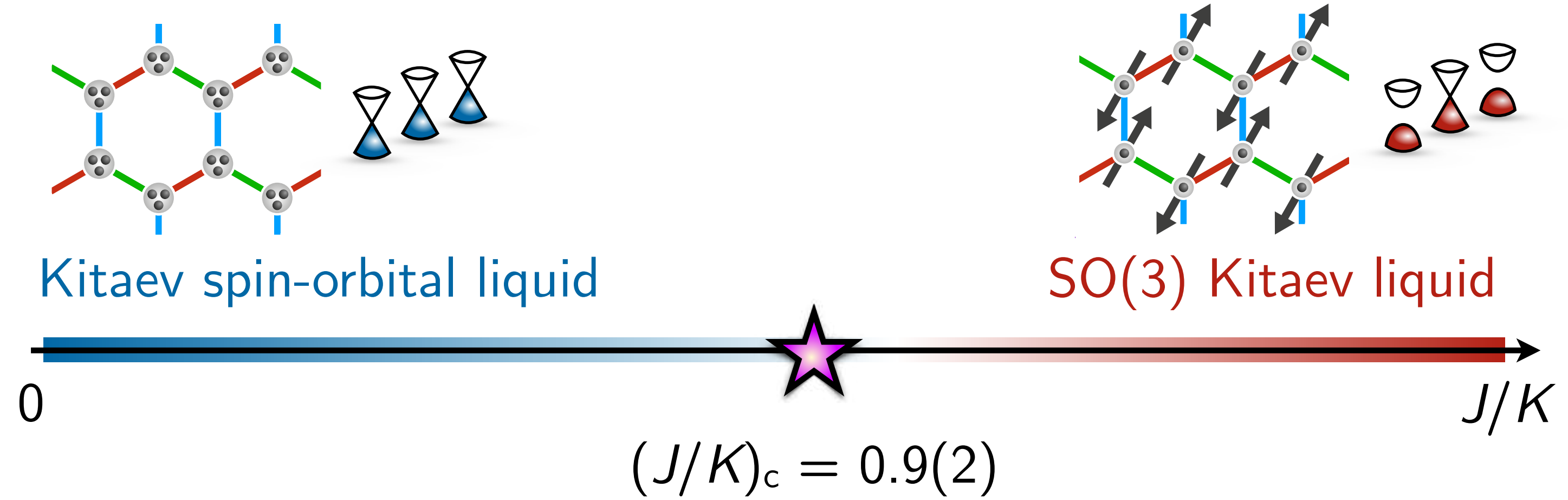
Gross-Neveu-SO(3)* transition

iDMRG:



... on cylinder with $L_y = 4$ unit cells

Phase diagram:



[Seifert, Dong, Chulliparambil, Vojta, Tu, LJ, PRL '20]

Effective field theory:
$$\mathcal{S} = \int d^2\vec{x}d\tau \left[\bar{\psi}\gamma^\mu\partial_\mu\psi + g\vec{\varphi} \cdot \bar{\psi}(\mathbb{1}_2 \otimes \vec{L})\psi \right] \quad \text{“Gross-Neveu-SO(3)”}$$

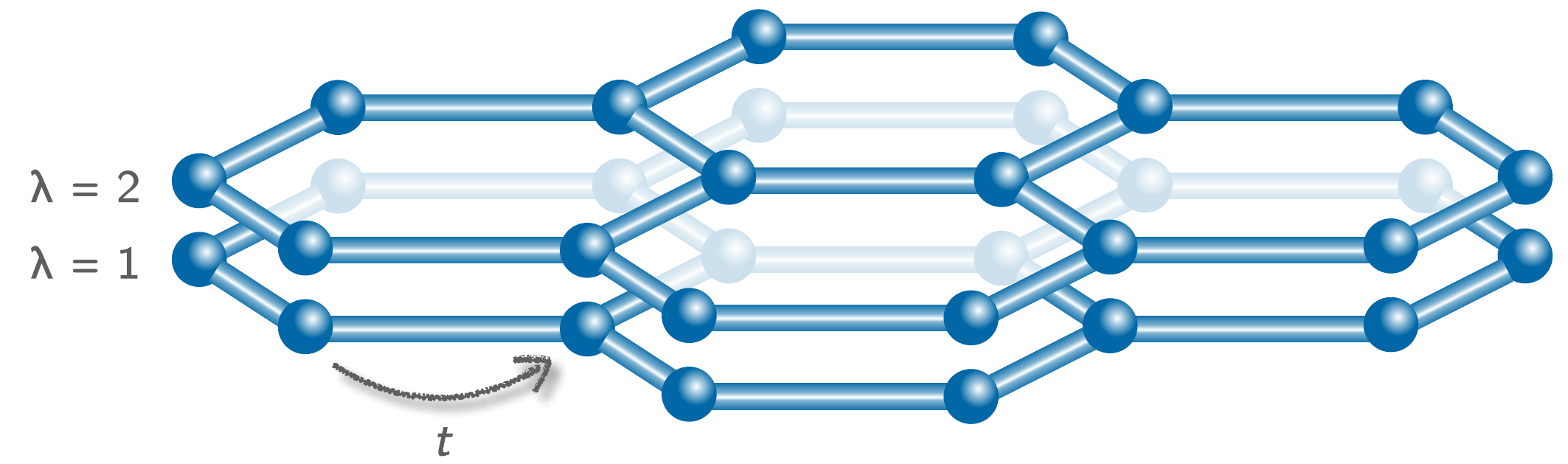
[Ray, Ihrig, Kruti, Gracey, Scherer, LJ, PRB '21]

Sign-problem-free bilayer model

Hamiltonian:

$$\mathcal{H} = -t \sum_{\langle ij \rangle} c_{i\lambda}^\dagger c_{j\lambda} - J \sum_i \left(c_{i\lambda}^\dagger \vec{L} \tau_{\lambda\lambda'}^z c_{i\lambda'} \right)^2$$

spin-1 matrices ↙



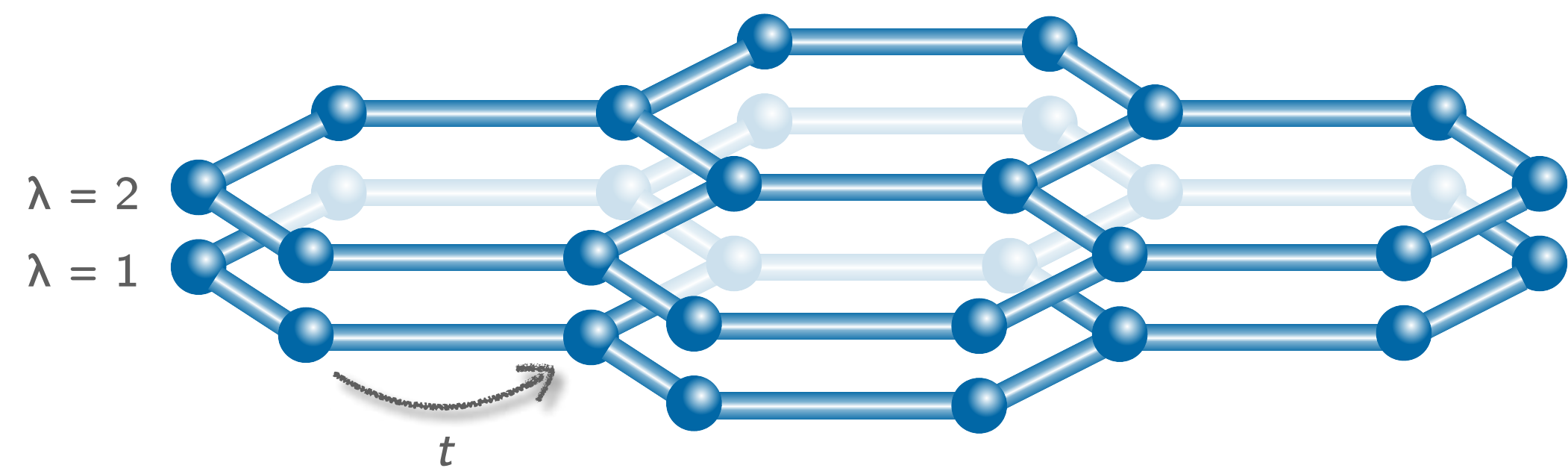
... with $SO(3) \times U_\lambda(1) \times U_c \times \mathbb{Z}_2$ symmetry

Sign-problem-free bilayer model

Hamiltonian:

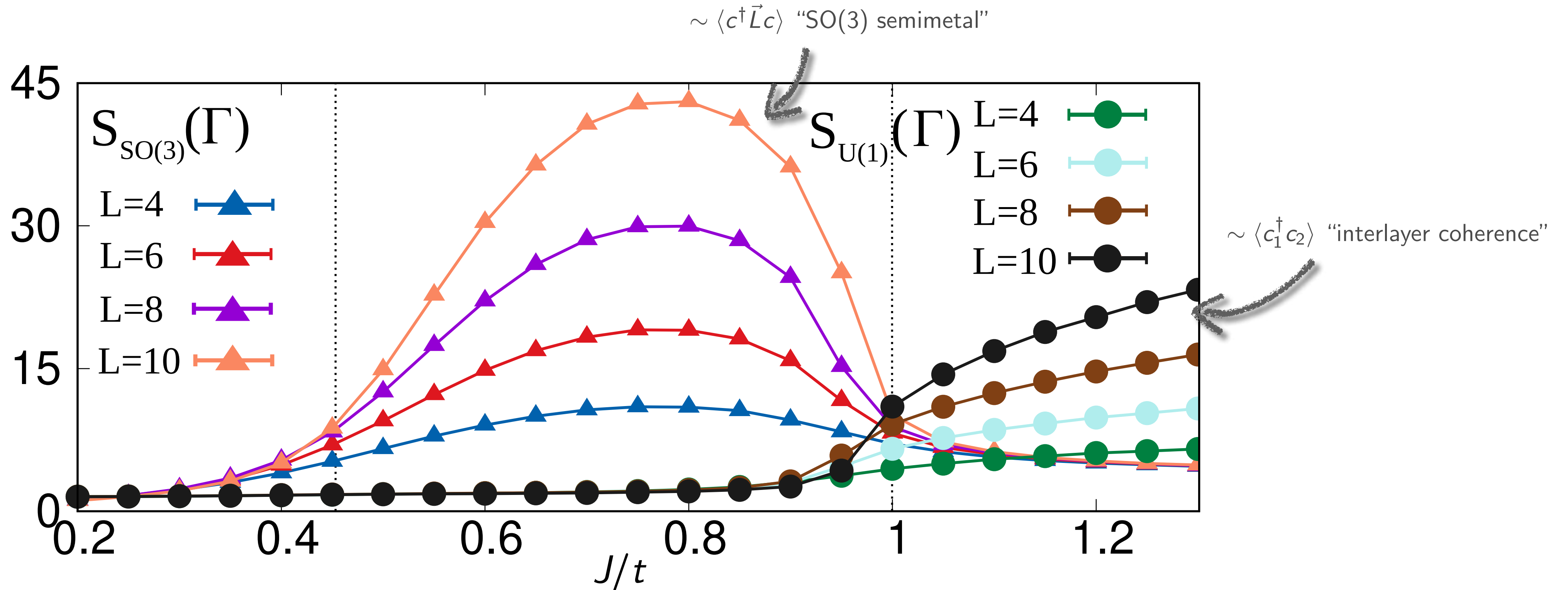
$$\mathcal{H} = -t \sum_{\langle ij \rangle} c_{i\lambda}^\dagger c_{j\lambda} - J \sum_i \left(c_{i\lambda}^\dagger \vec{L} \tau_{\lambda\lambda'}^z c_{i\lambda'} \right)^2$$

spin-1 matrices



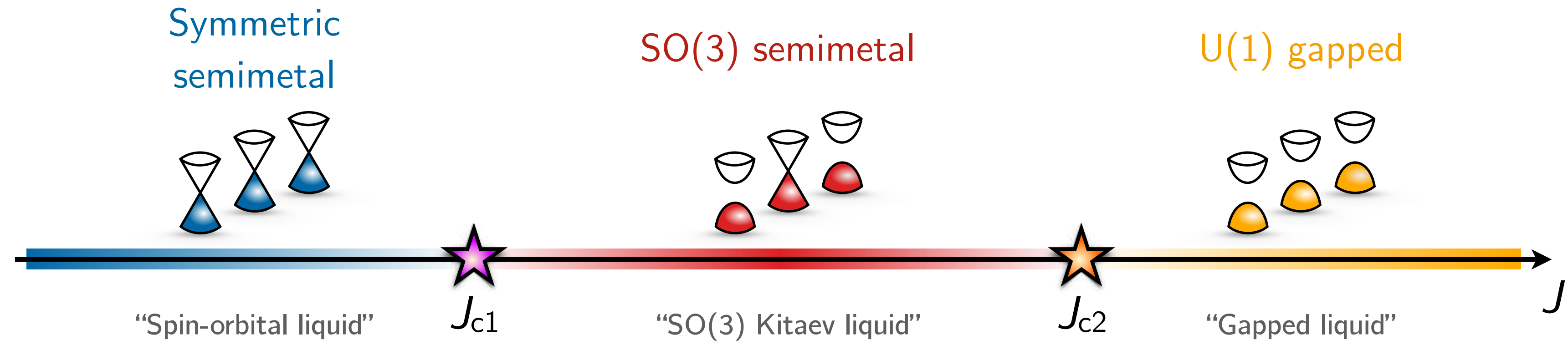
... with $SO(3) \times U_\lambda(1) \times U_c \times \mathbb{Z}_2$ symmetry

QMC structure factors:



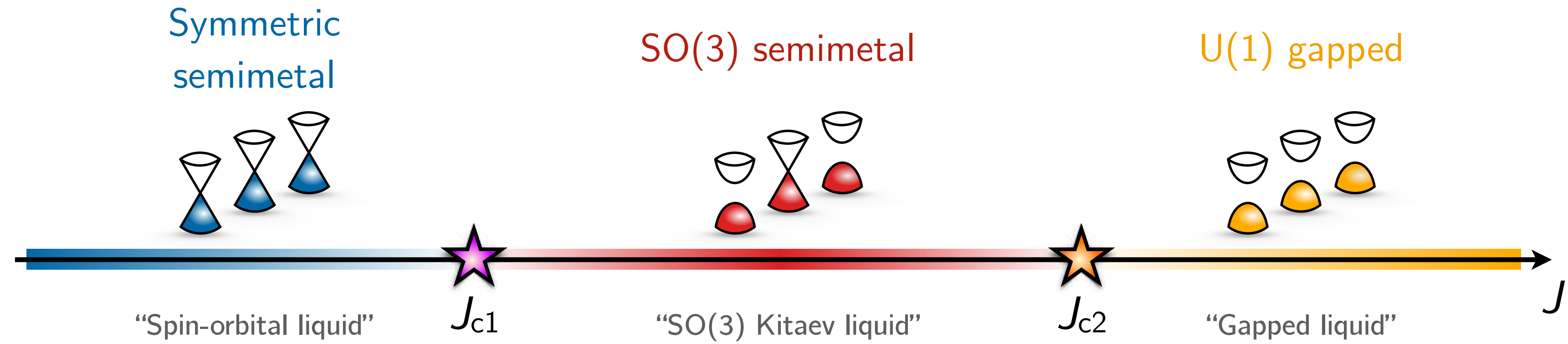
Sign-problem-free bilayer model

Phase diagram:

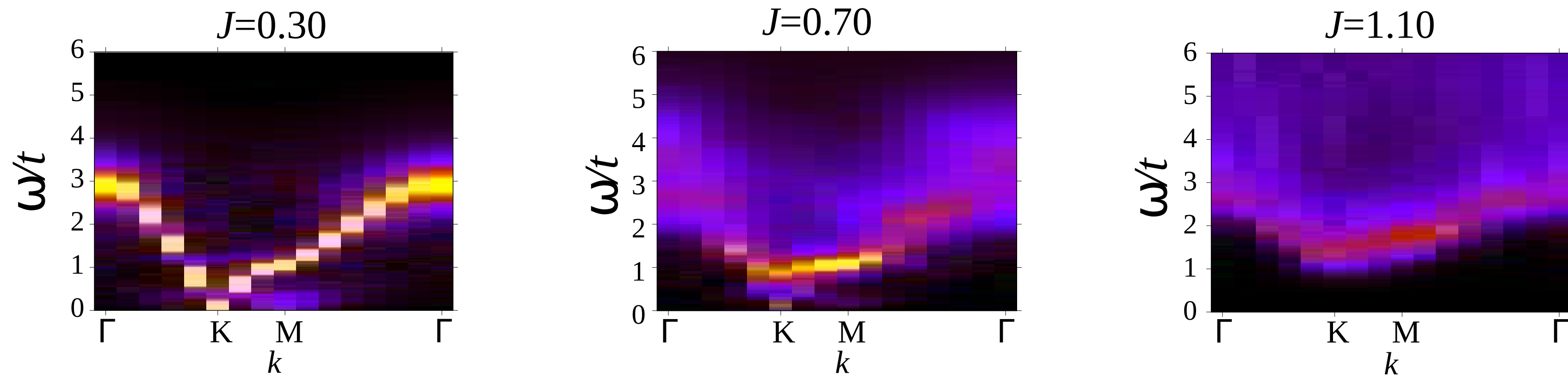


Sign-problem-free bilayer model

Phase diagram:

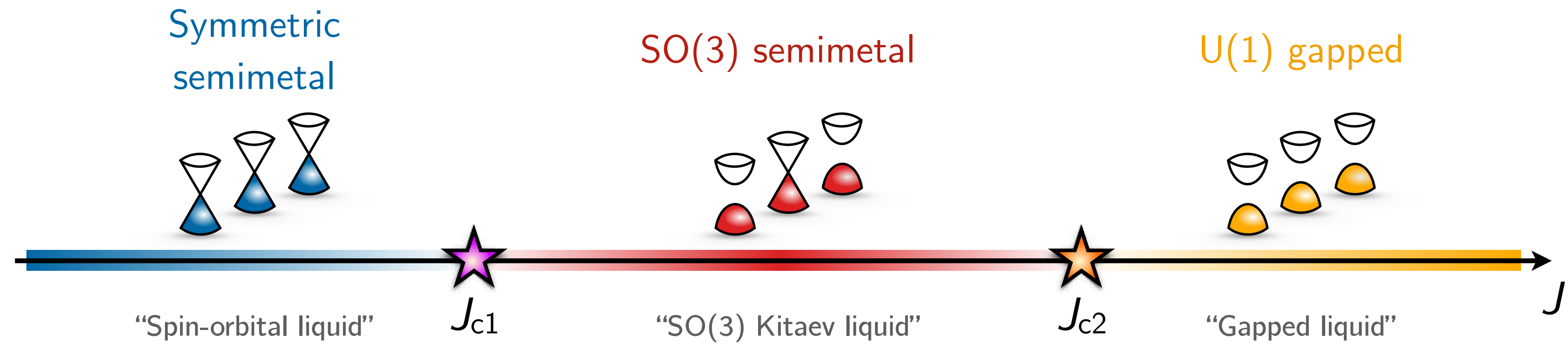


Fermion spectral function:

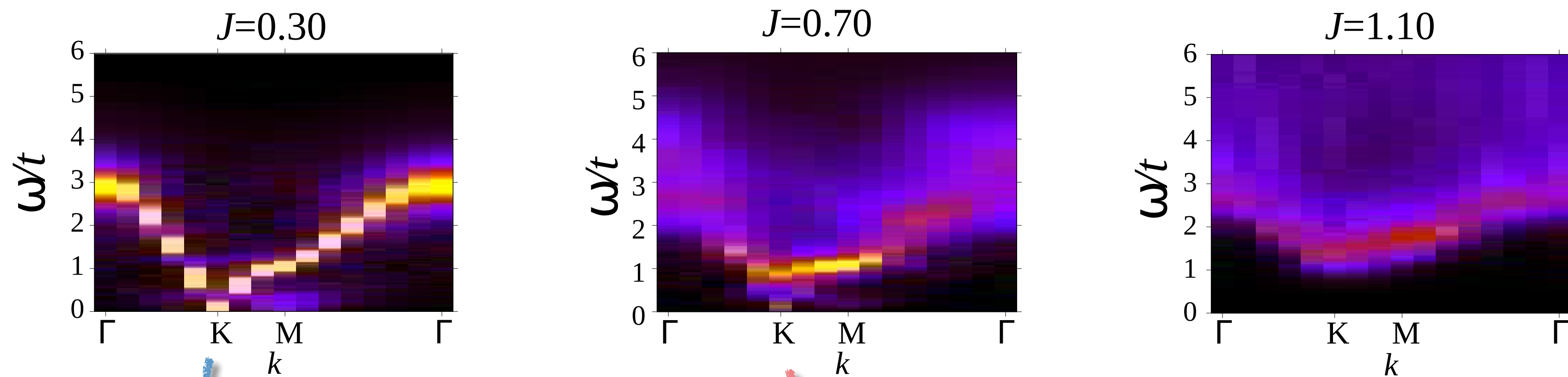


Sign-problem-free bilayer model

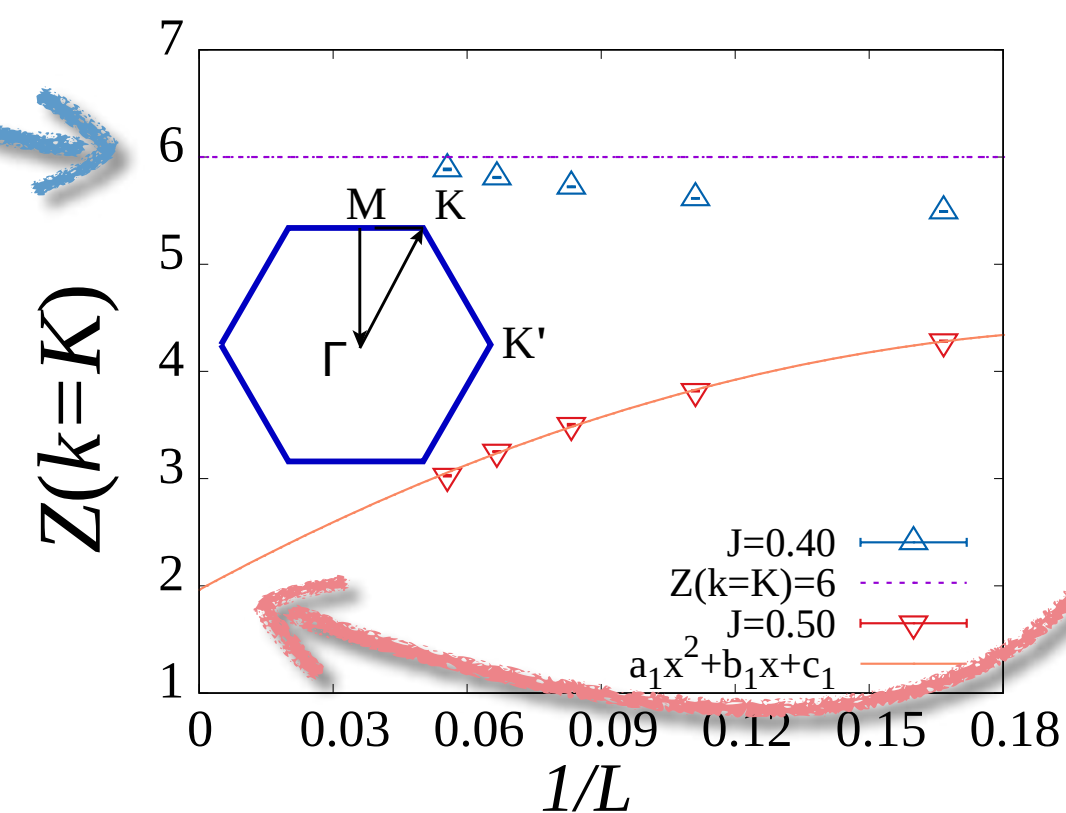
Phase diagram:



Fermion spectral function:

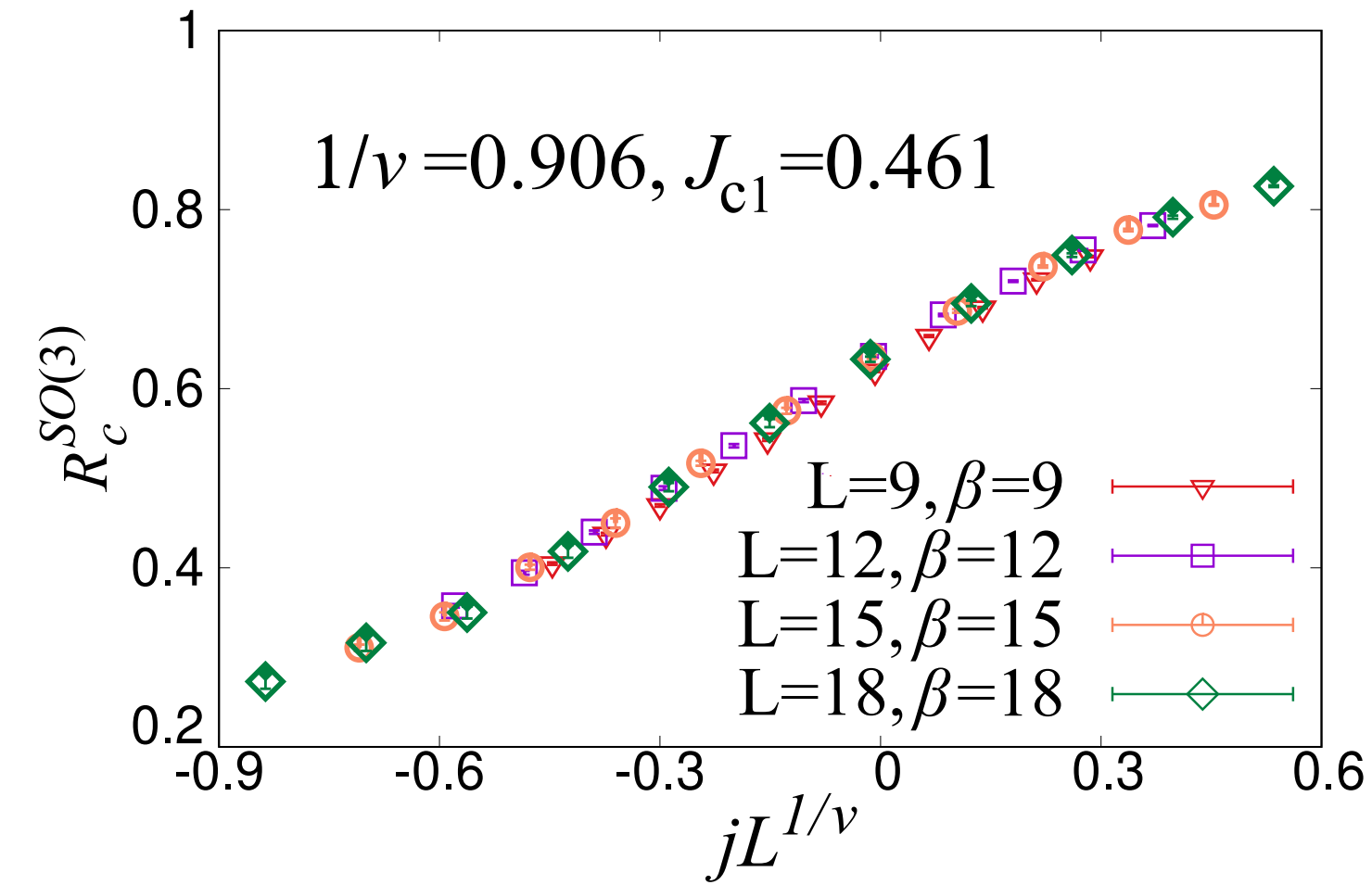
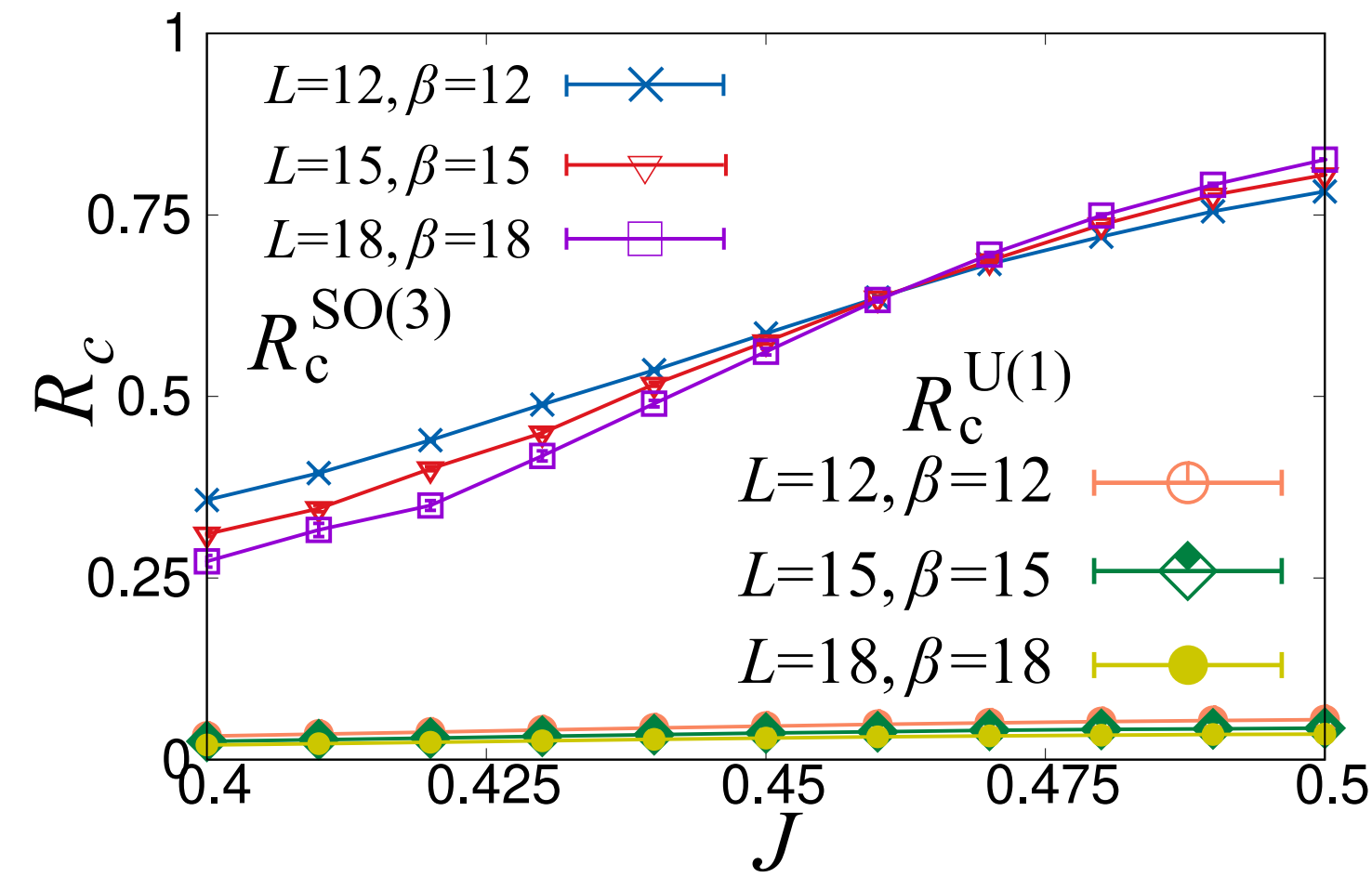
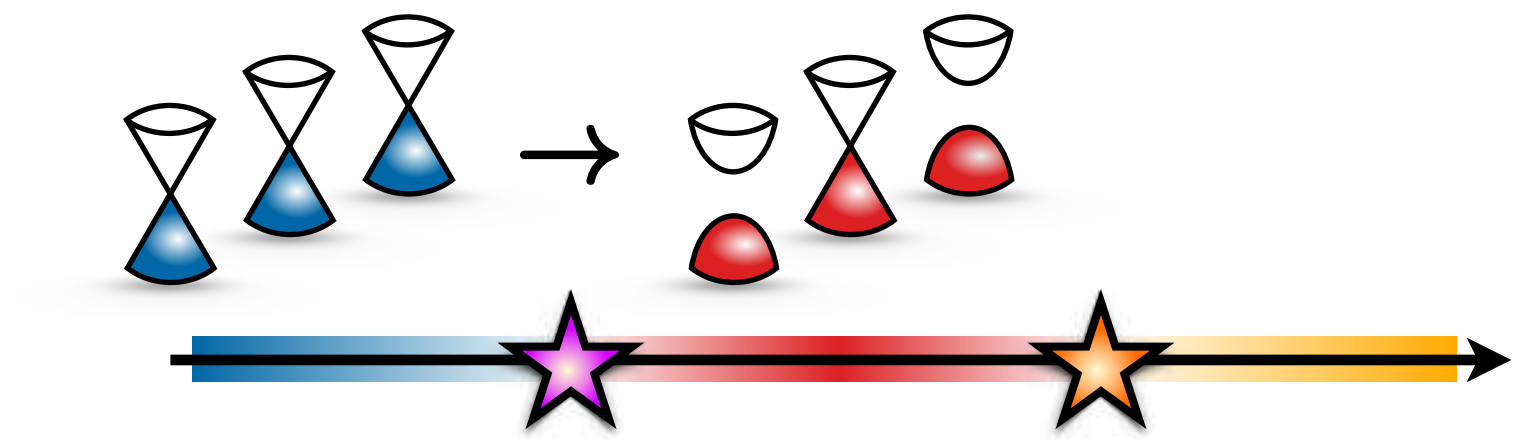


Quasiparticle weight:



Gross-Neveu-SO(3) transition at J_{c1}

Correlation ratio:
$$R_c = 1 - \frac{S(\Gamma + d\vec{k})}{S(\Gamma)}$$

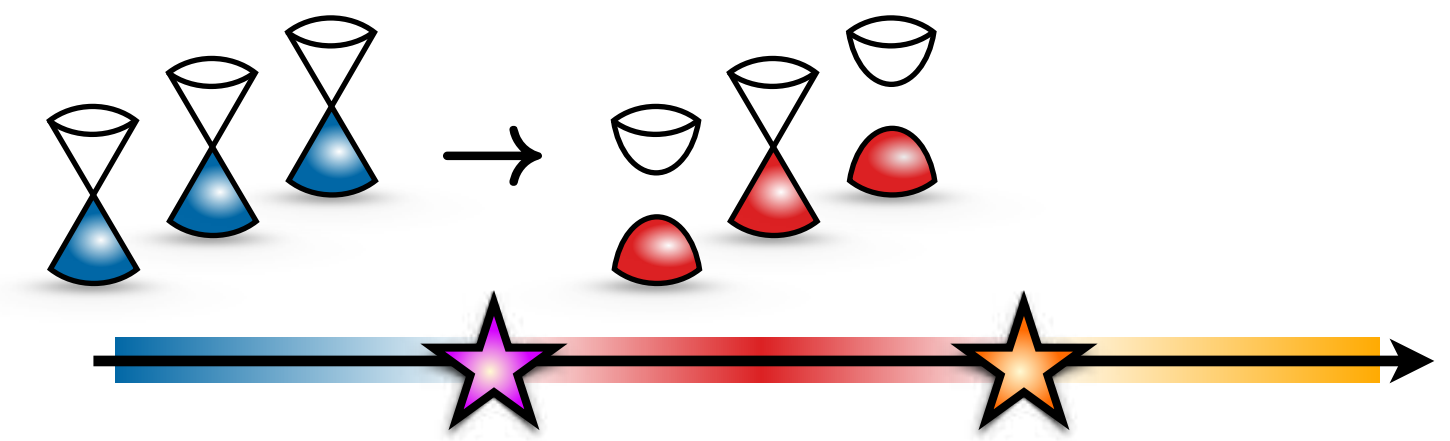
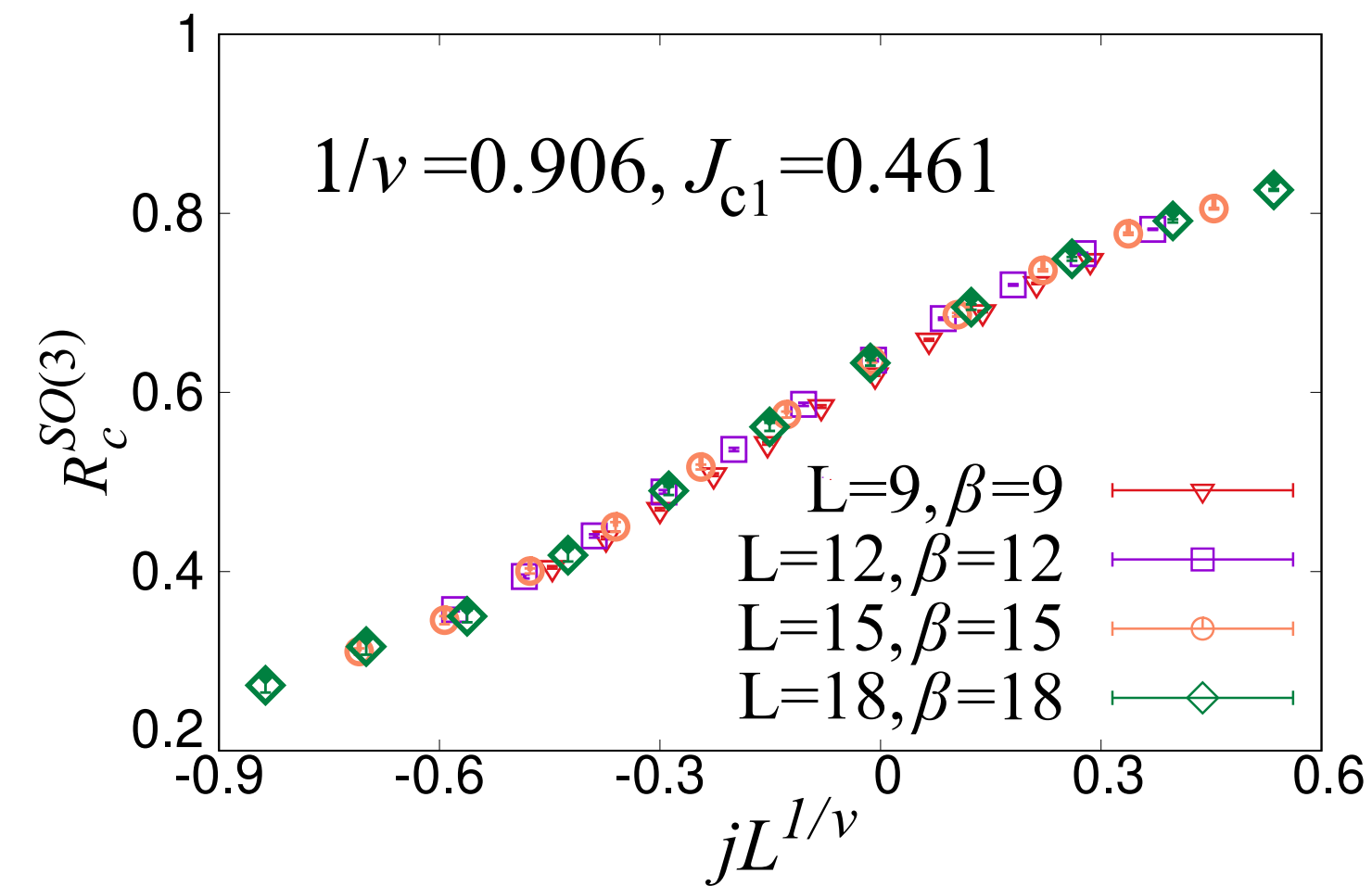
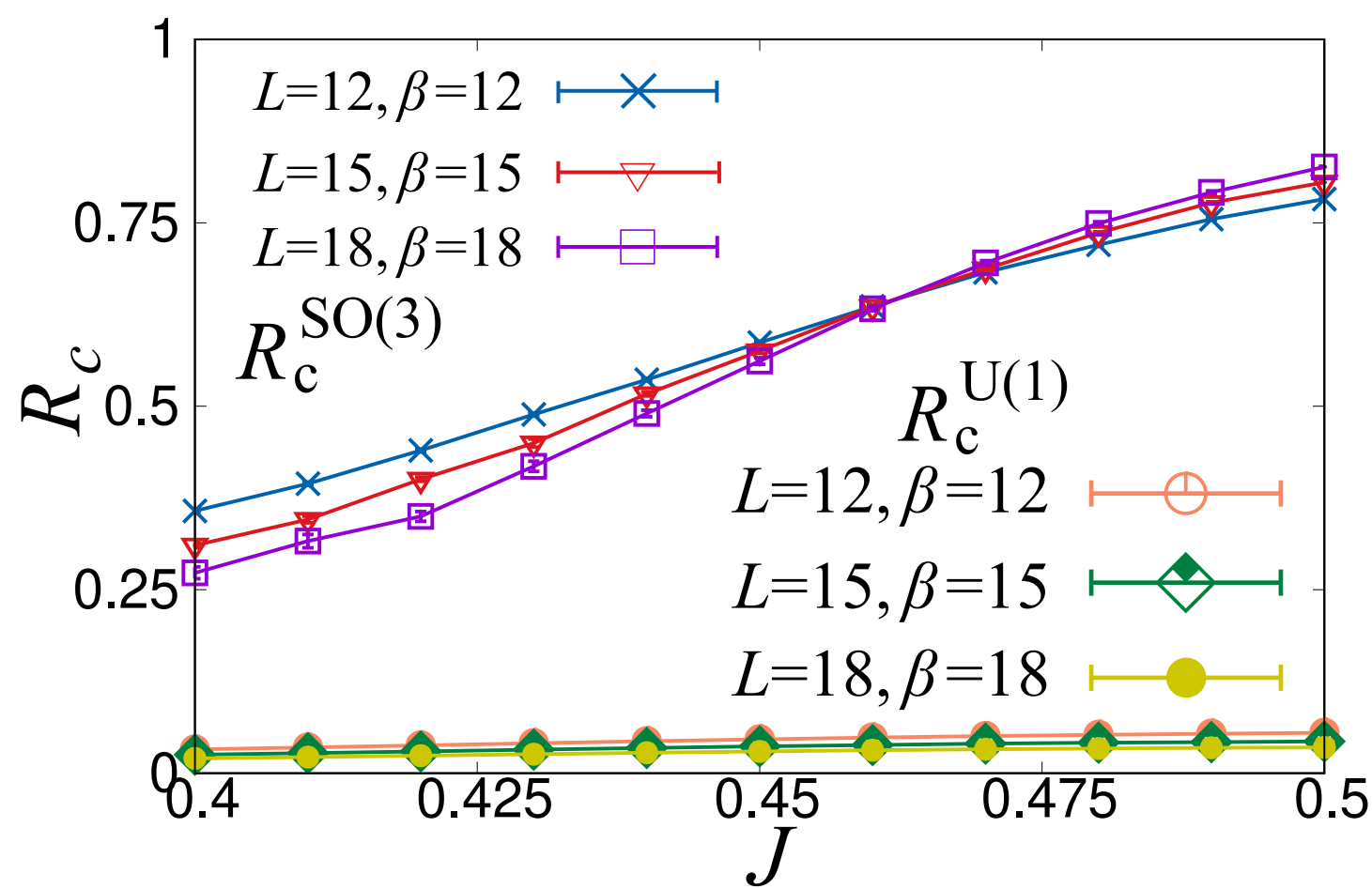


$\Rightarrow 1/\nu = 0.906(35)$

... cf. $1/\nu = 0.93(4)$ and $\eta_\phi = 0.83(4)$ from field theory ($N = 12$)
 [Ray, Ihrig, Kruti, Gracey, Scherer, LJ, PRB '21]

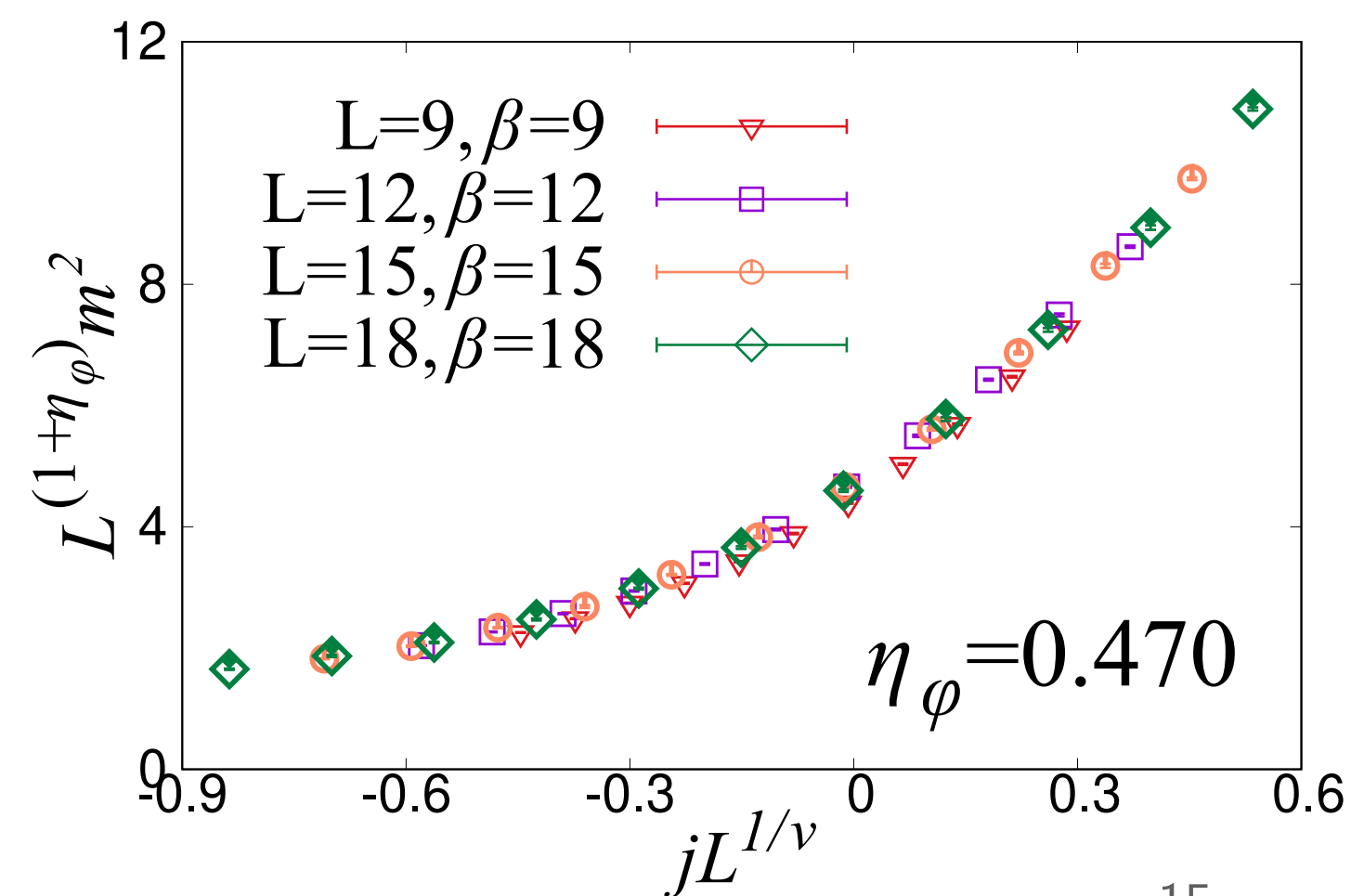
Gross-Neveu-SO(3) transition at J_{c1}

Correlation ratio:
$$R_c = 1 - \frac{S(\Gamma + d\vec{k})}{S(\Gamma)}$$



$\Rightarrow 1/\nu = 0.906(35)$

Order parameter:

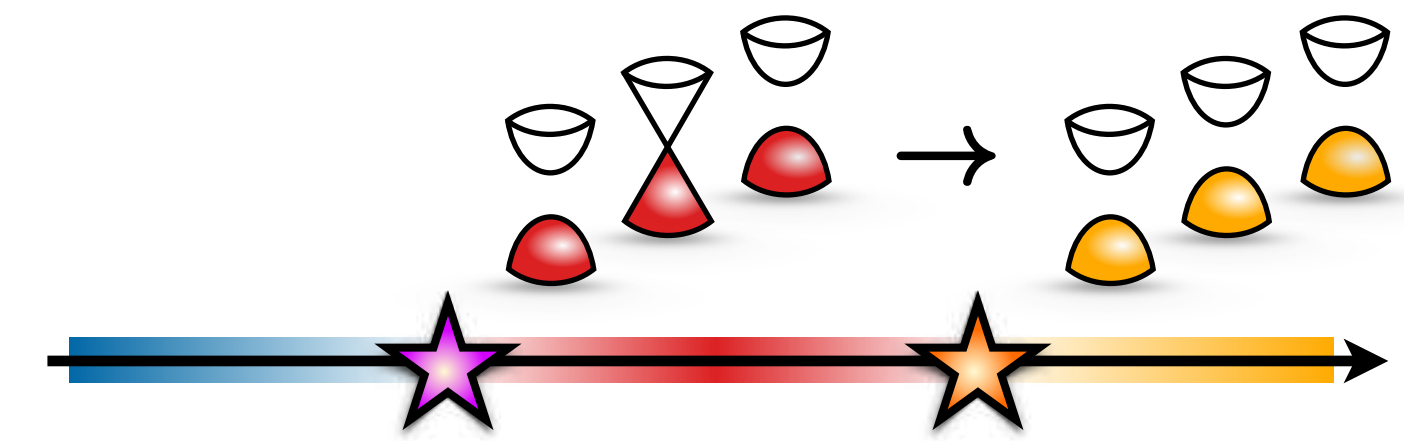
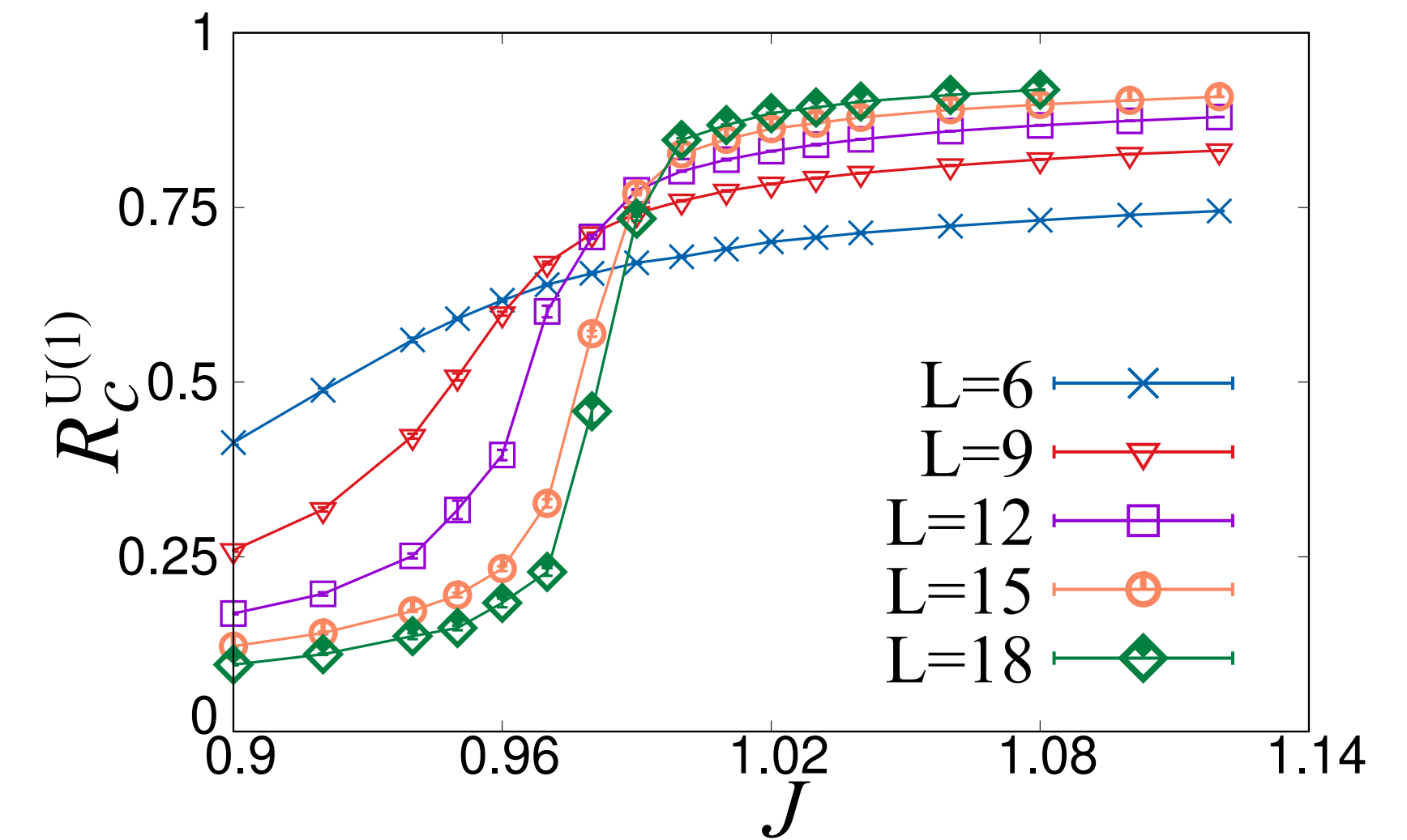
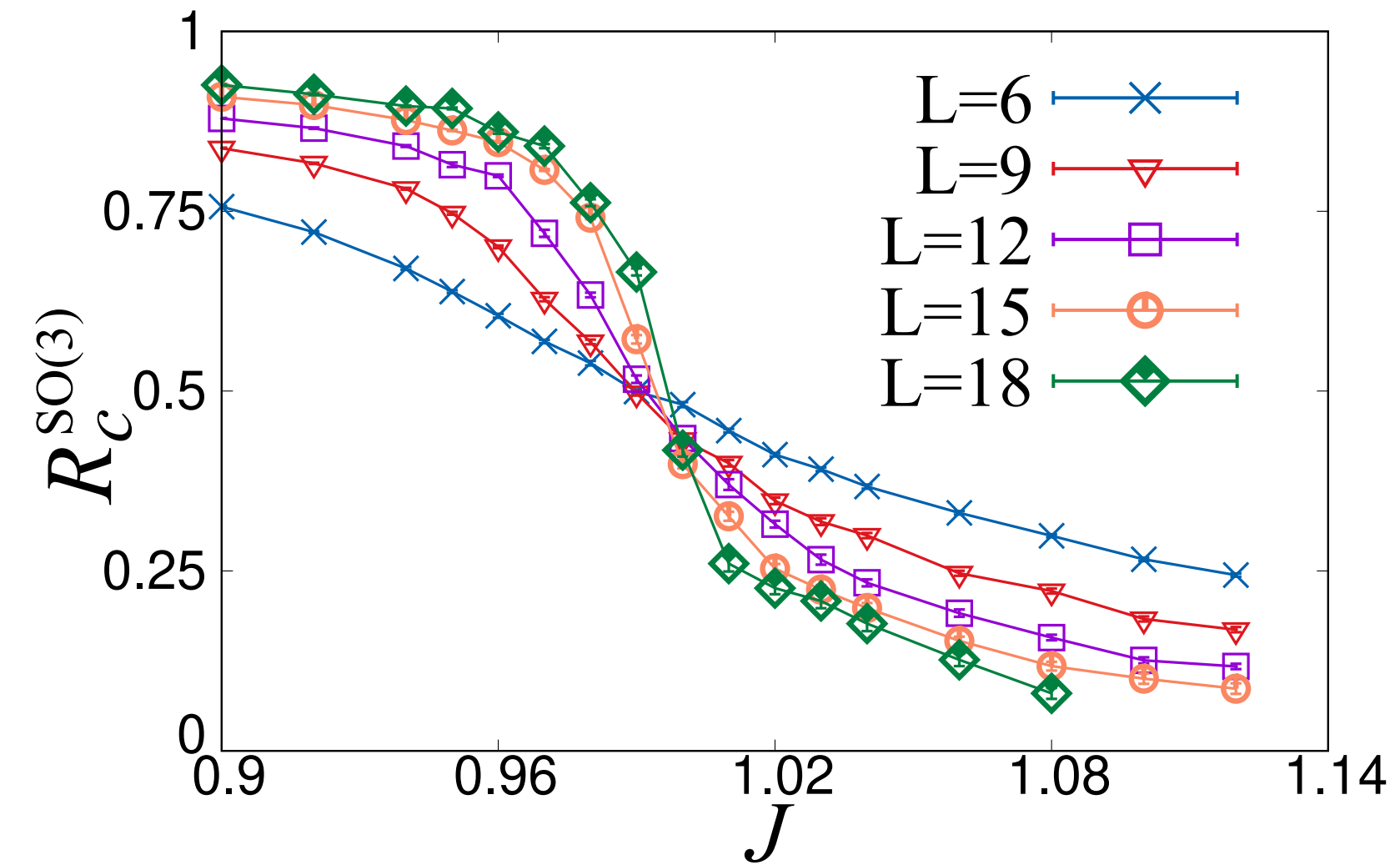


$\Rightarrow \eta_\phi = 0.470(13)$

... cf. $1/\nu = 0.93(4)$ and $\eta_\phi = 0.83(4)$ from field theory ($N = 12$)
 [Ray, Ihrig, Kruti, Gracey, Scherer, LJ, PRB '21]

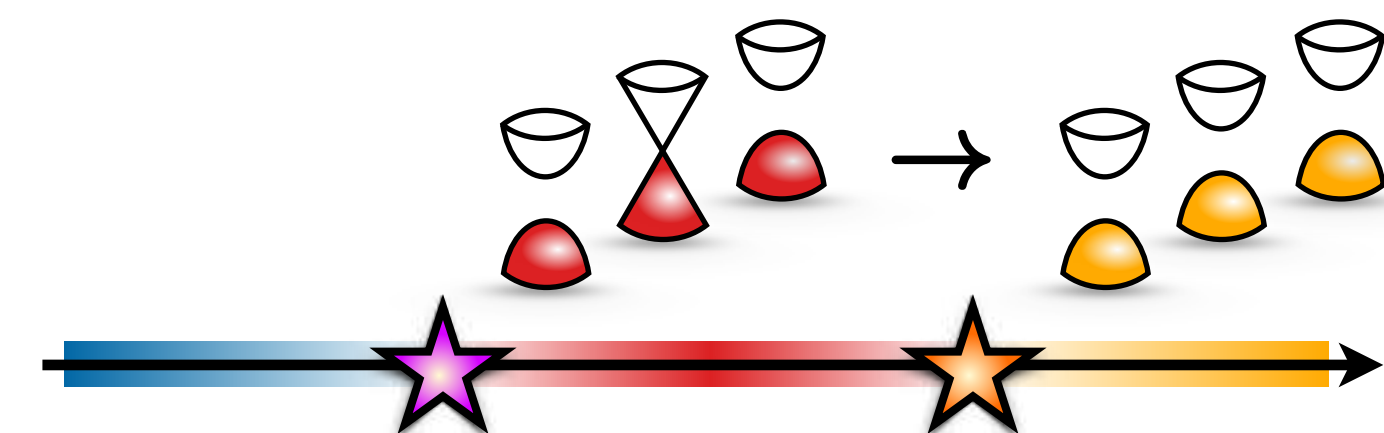
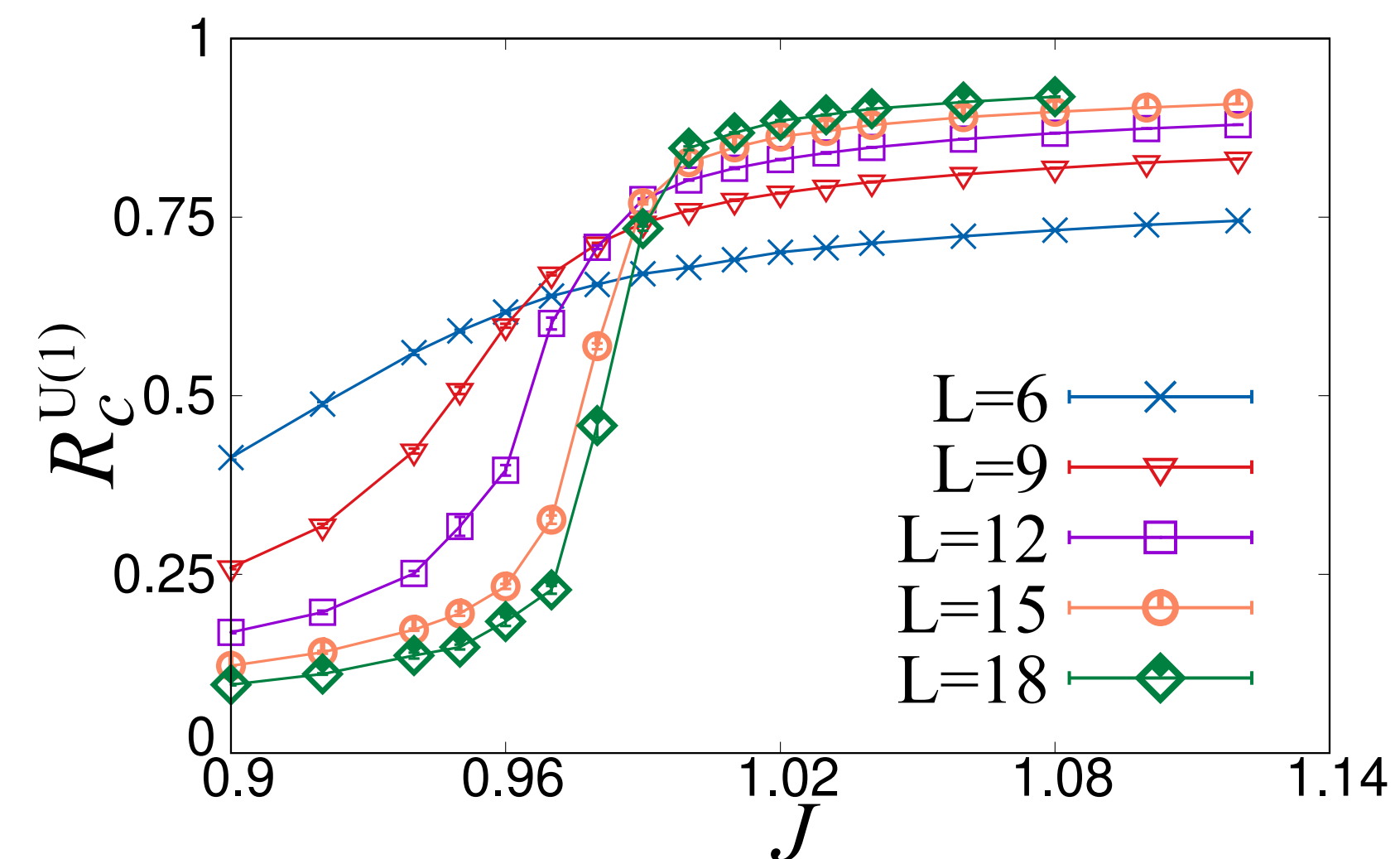
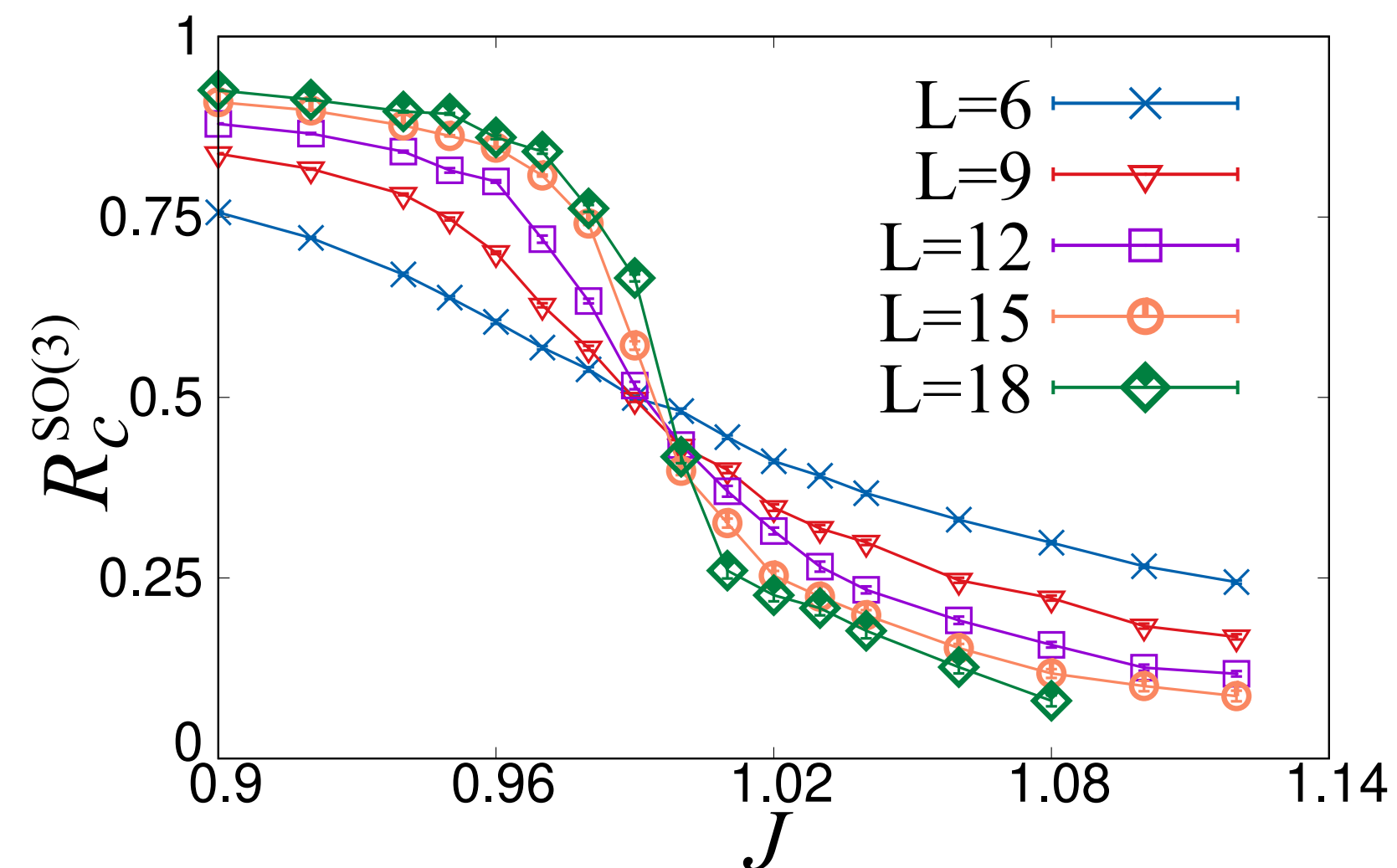
Order-to-order transition at J_{c2}

Correlation ratios:

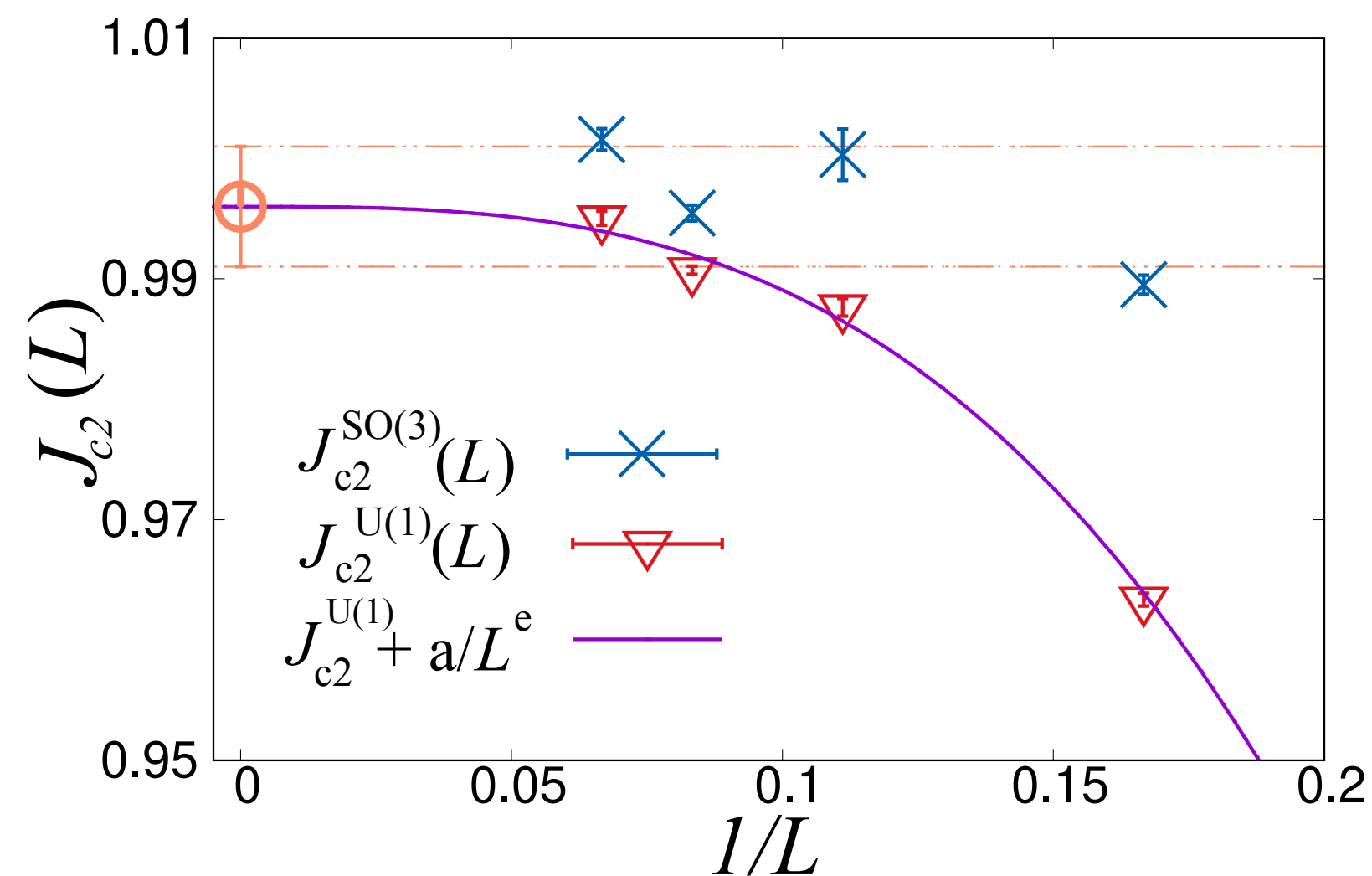


Order-to-order transition at J_{c2}

Correlation ratios:



Critical couplings:



$$\Rightarrow J_{c2}^{SO(3)} = J_{c2}^{U(1)} \text{ unique!}$$

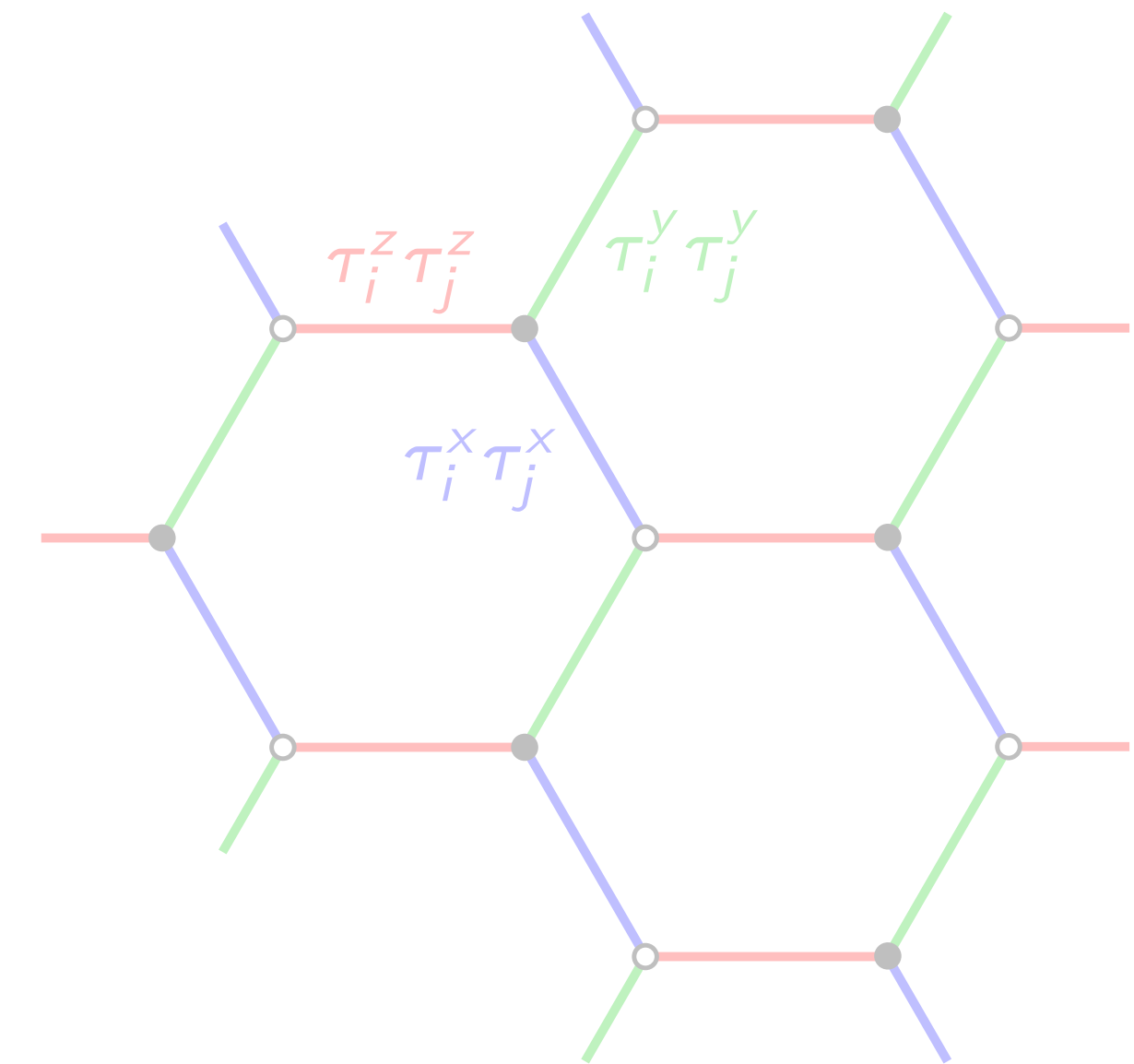
Metallic deconfined QCP?

Outline

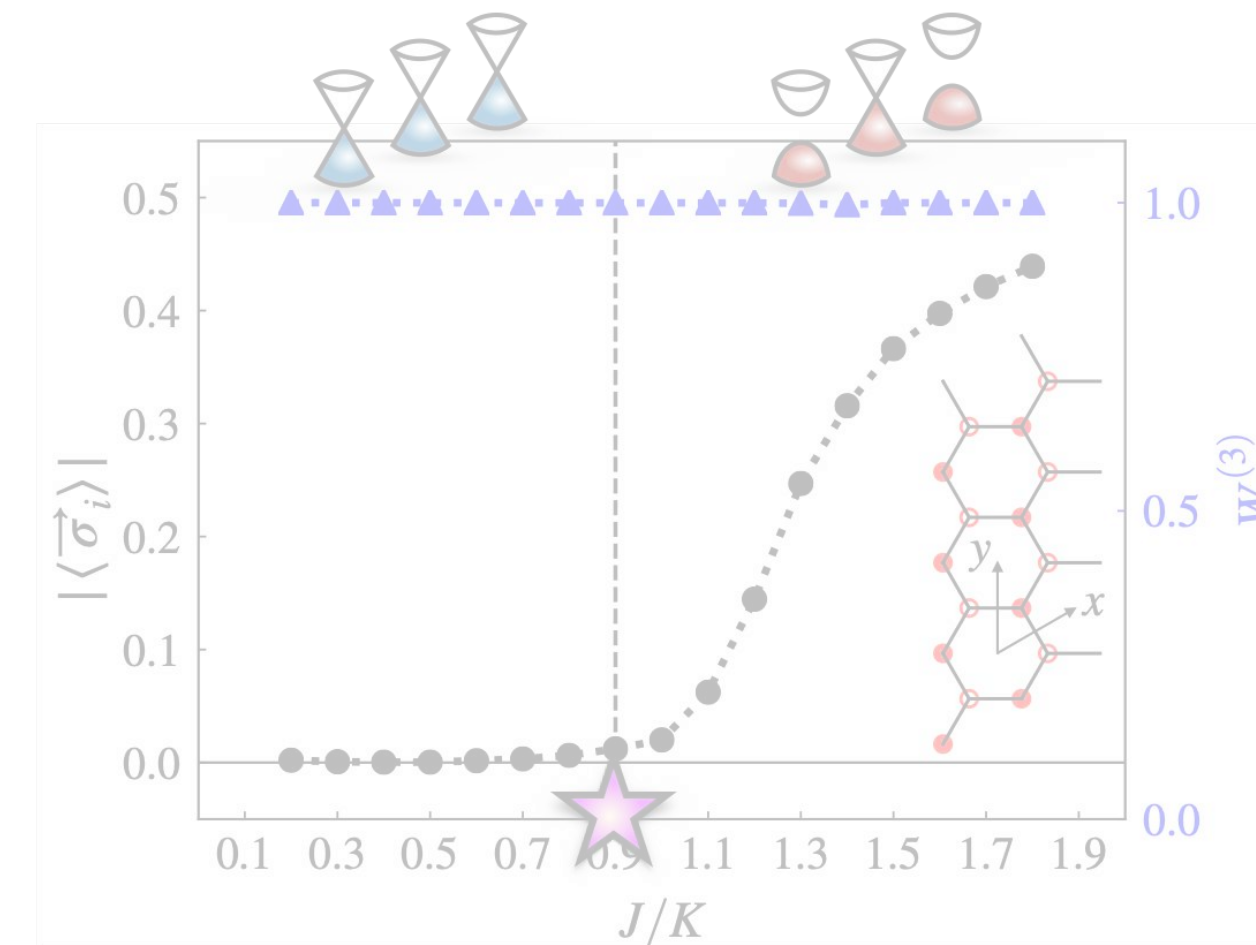
(1) Fractionalized quantum criticality



(2) From Kitaev to Kitaev-Kugel-Khomskii



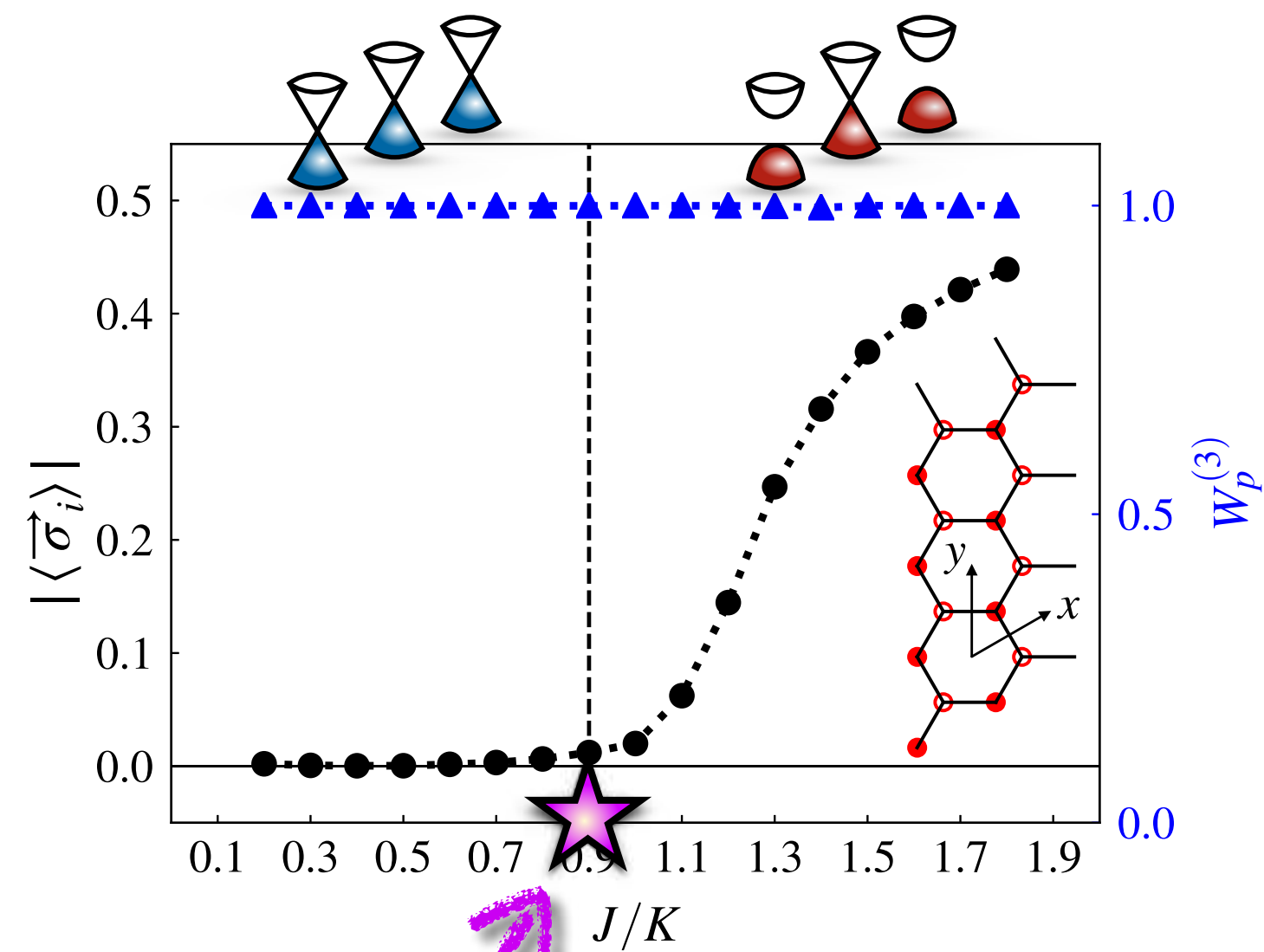
(3) Kitaev-Heisenberg spin-orbital models



(4) Conclusions

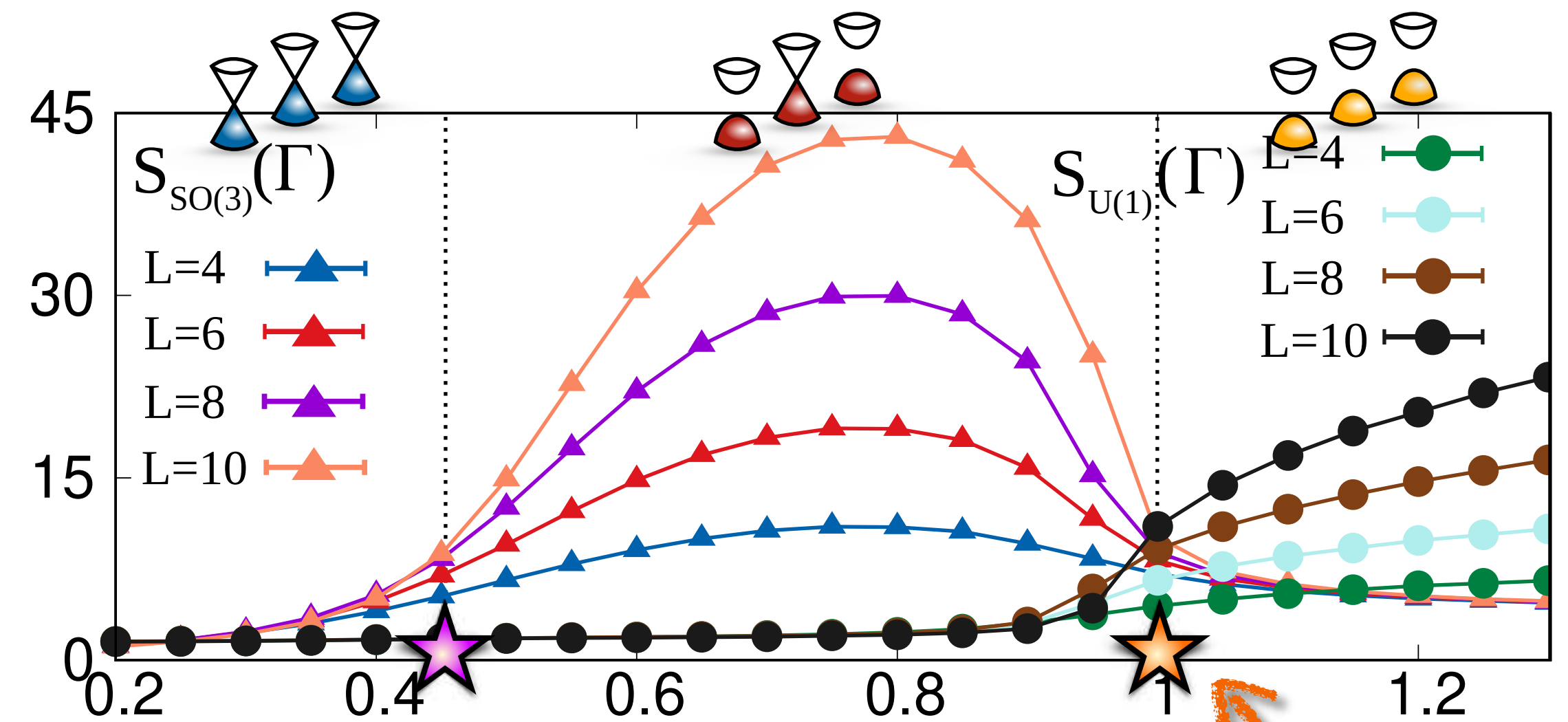
Conclusions

Kitaev-Heisenberg spin-orbital model:



Gross-Neveu-SO(3)*

Effective bilayer honeycomb model:

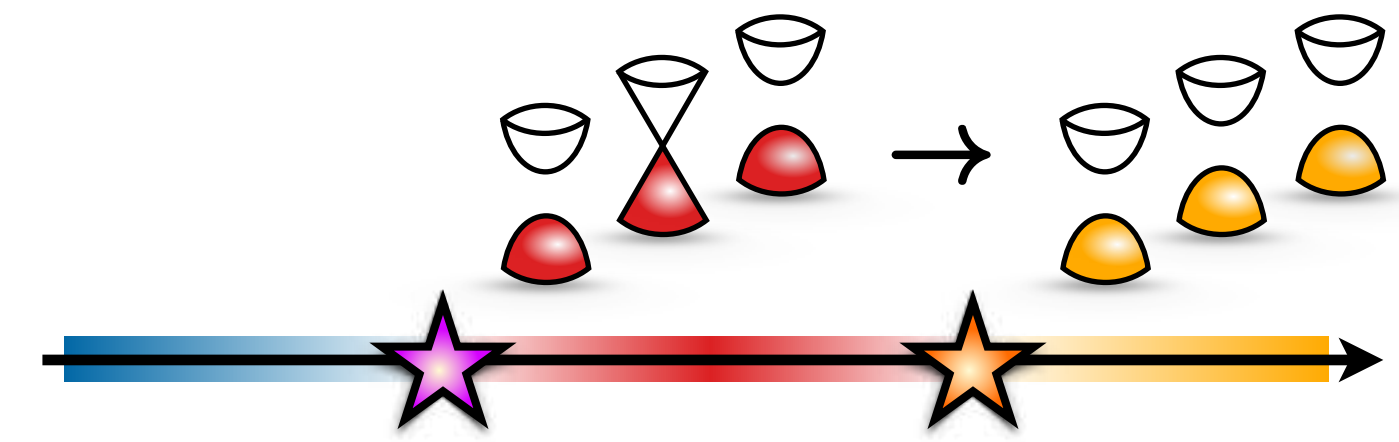
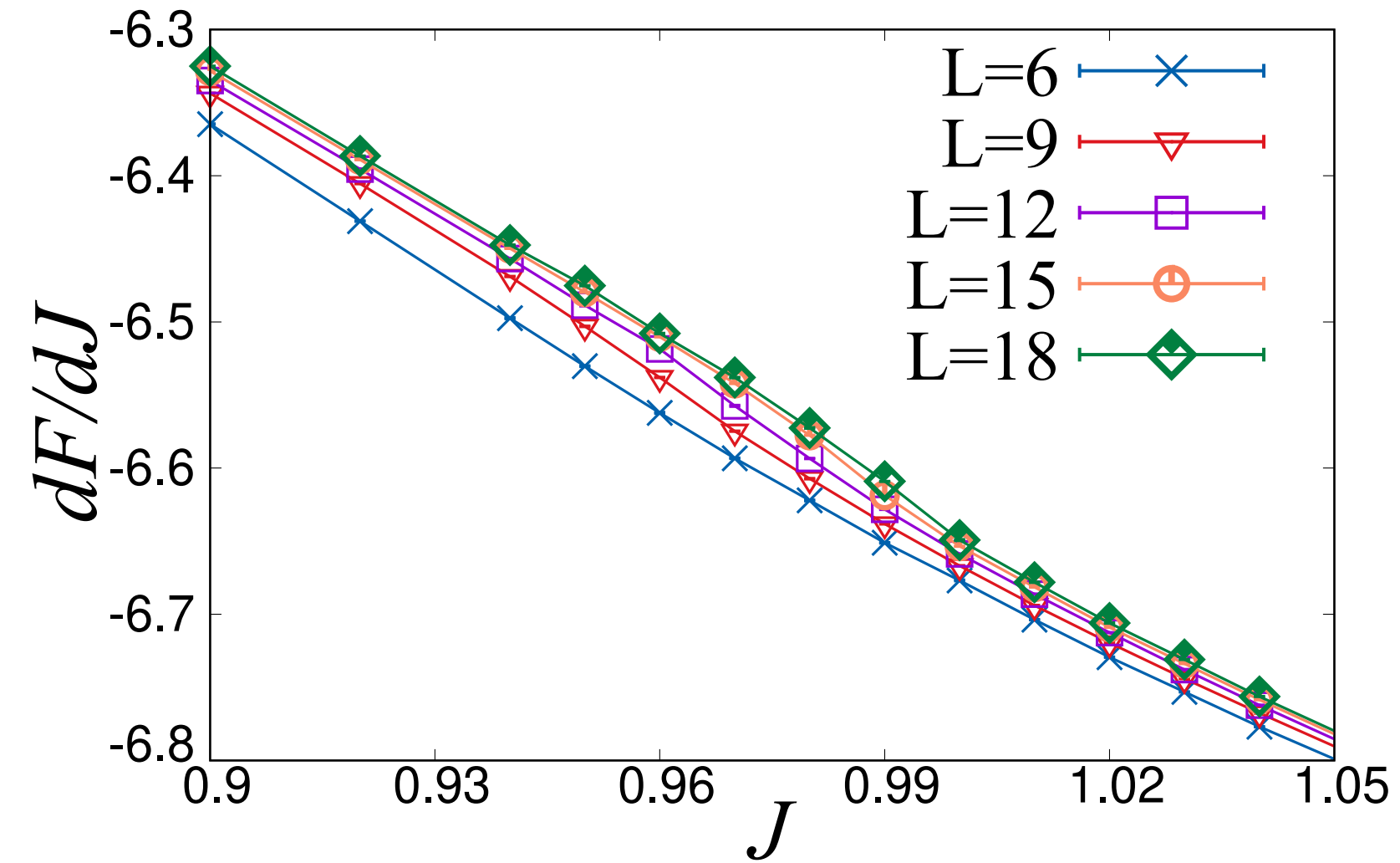


Gross-Neveu-SO(3)

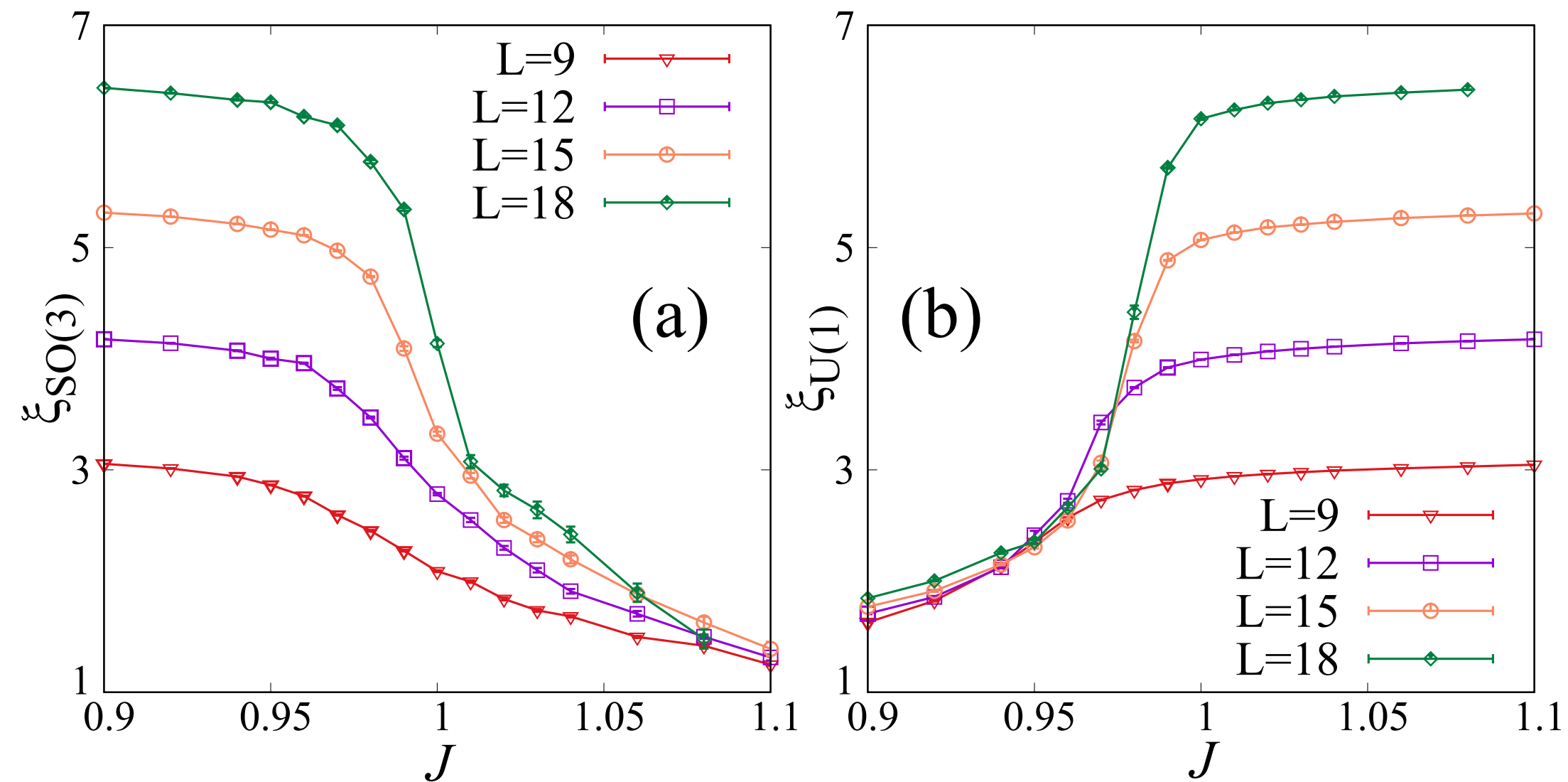
Metallic
deconfined QCP?

Order-to-order transition at J_{c2}

Free energy:



Correlation lengths:

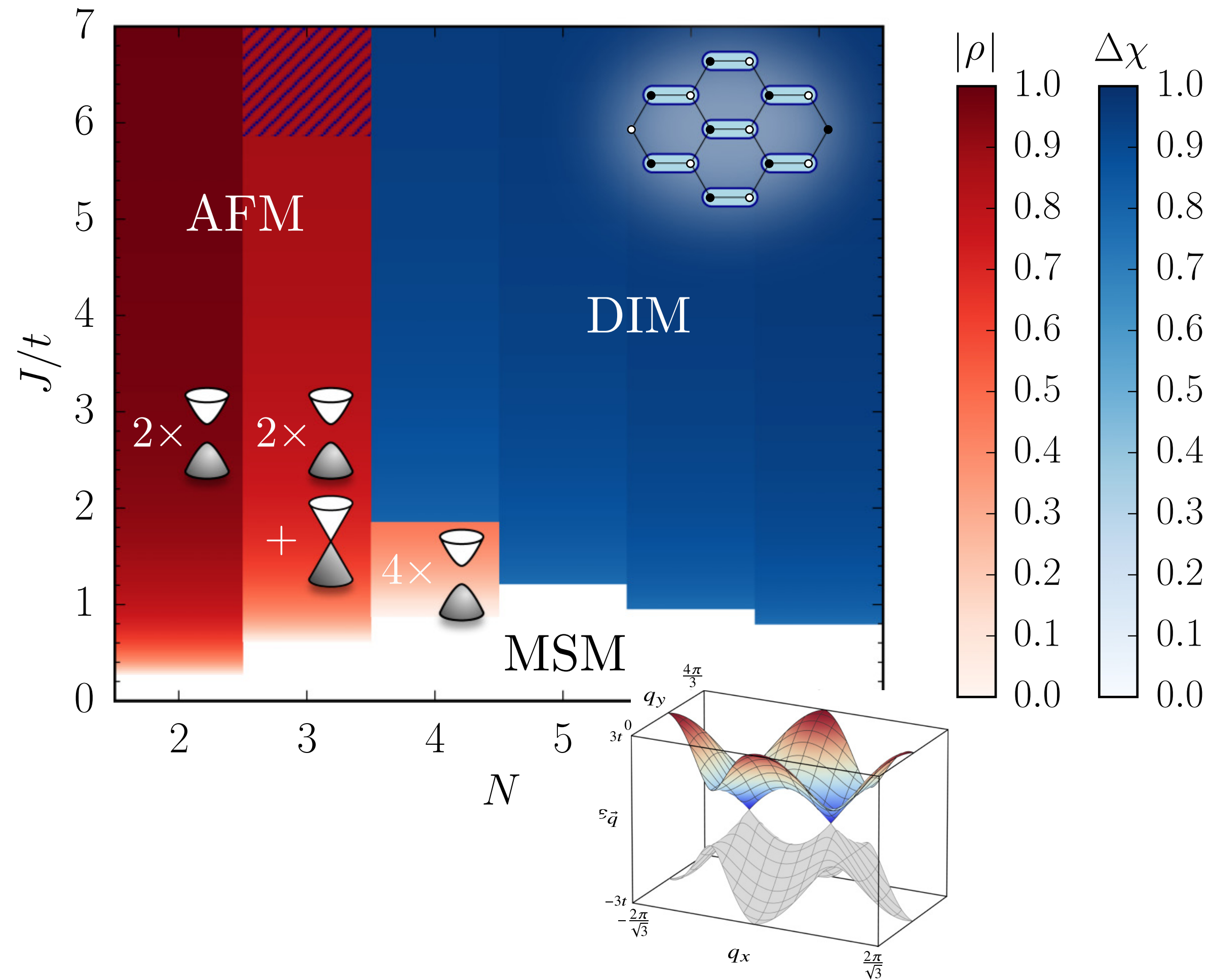


$$\xi^2 = \frac{1}{2d} \frac{\sum_{\vec{r}} |\vec{r}|^2 S(\vec{r})}{\sum_{\vec{r}} S(\vec{r})}$$

[Liu, Vojta, Assaad, LJ, PRL '22 (Editors' Suggestion)]

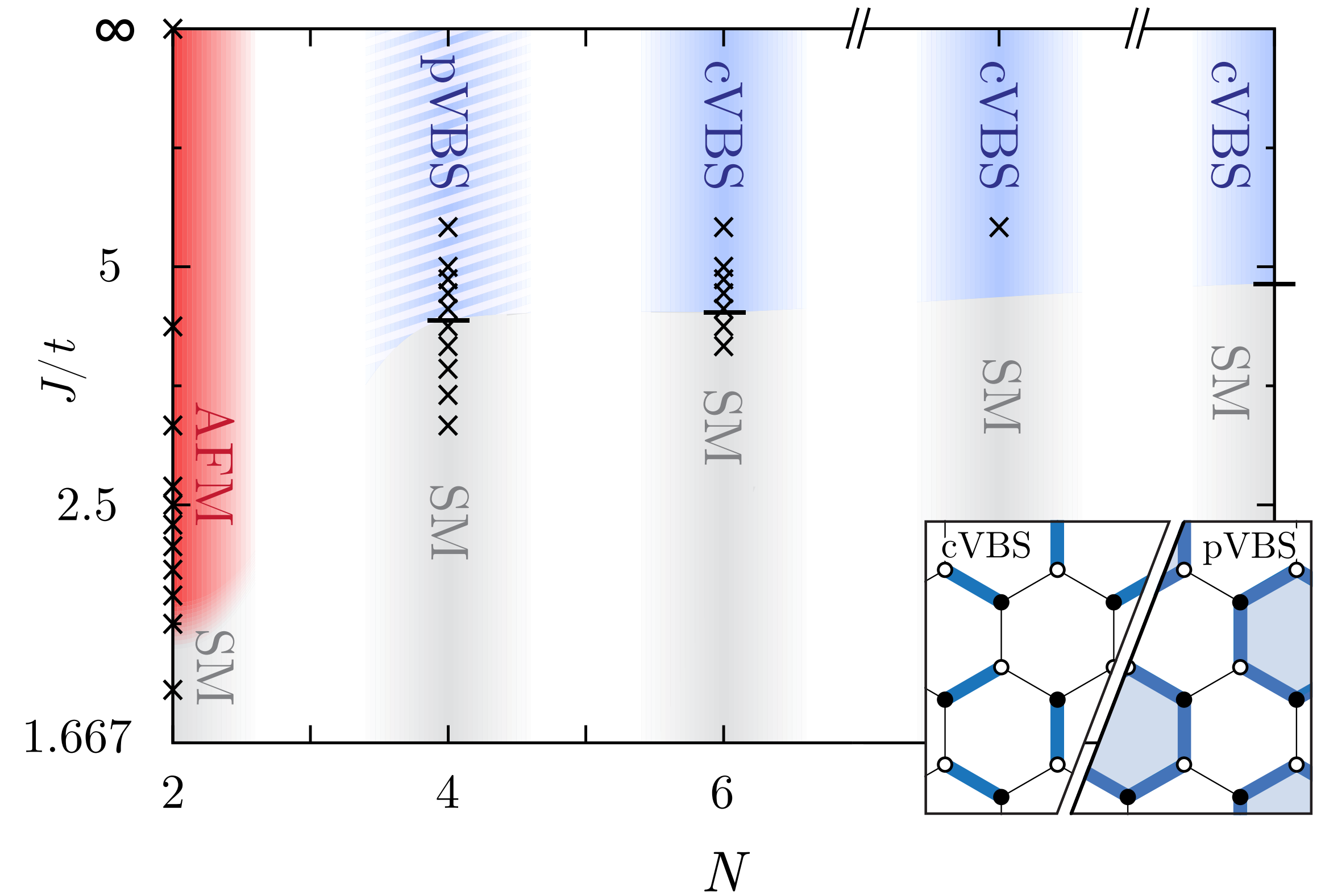
SO(N) generalization

SO(N) Majorana-Hubbard models



[LJ & Seifert, PRB '22]

SU(N) Hubbard-Heisenberg models



[Lang, Meng, Muramatsu, Wessel, Assaad, PRL '13]

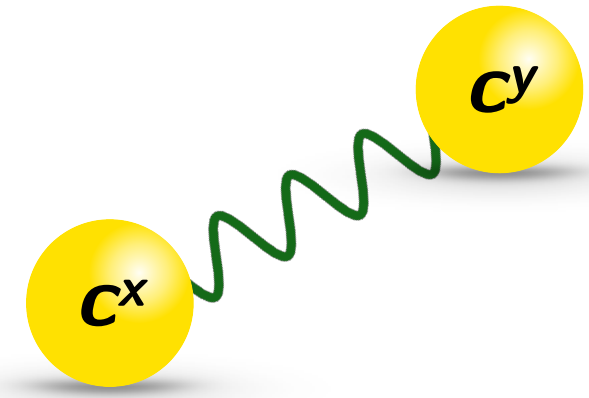
[Affleck & Marston, PRB '88]

[Read & Sachdev, NPB '89]

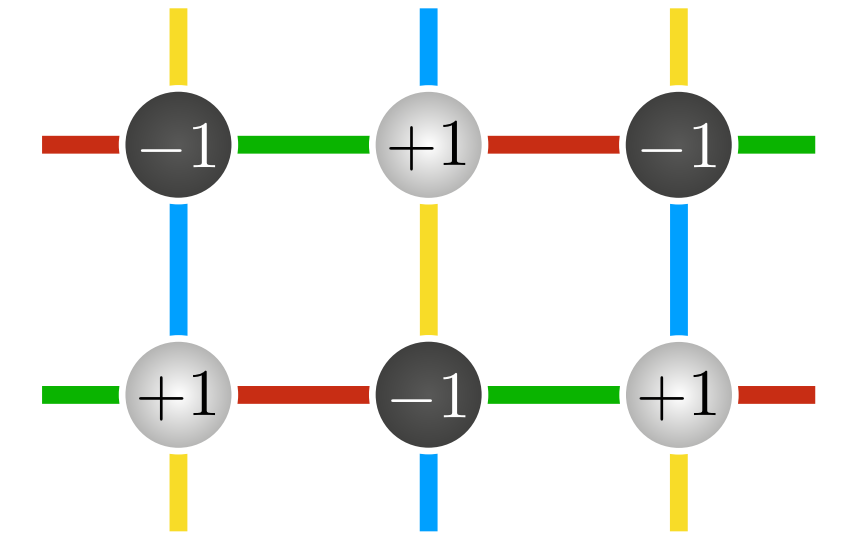
Kitaev-Ising spin-orbital model

Ising perturbation:

$$H = H_K + J^z \sum_{\langle ij \rangle} \sigma_i^z \sigma_j^z \otimes \mathbb{1}_i \mathbb{1}_j$$



“Kitaev” spin-orbital liquid



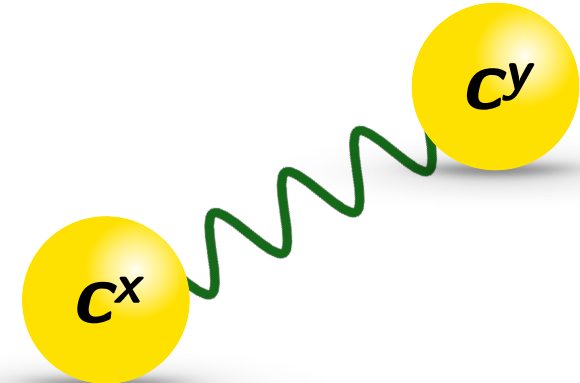
Ising spin order



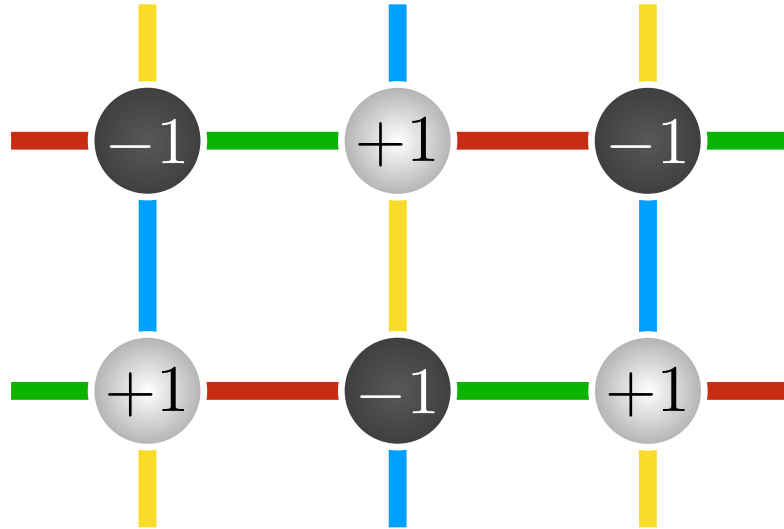
Kitaev-Ising spin-orbital model

Ising perturbation:

$$H = H_K + J^z \sum_{\langle ij \rangle} \sigma_i^z \sigma_j^z \otimes \mathbb{1}_i \mathbb{1}_j$$



“Kitaev” spin-orbital liquid



Ising spin order



Parton representation:

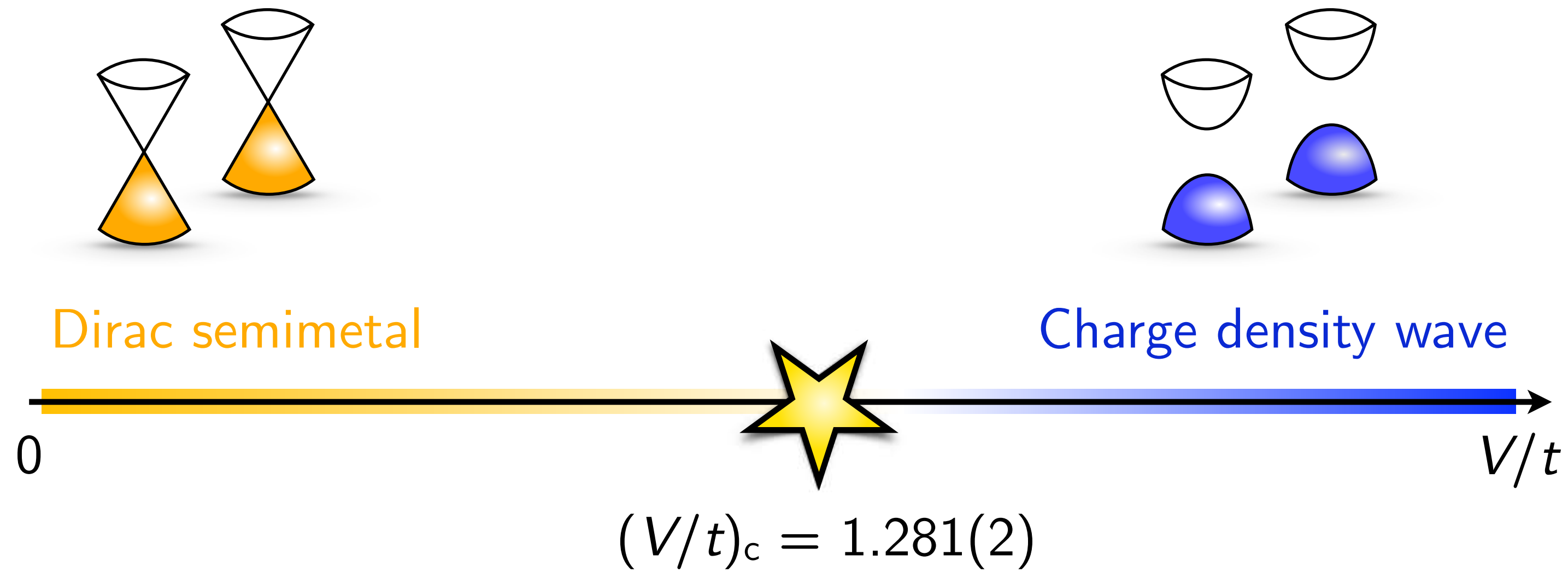
$$H \mapsto \sum_{\langle ij \rangle} \left[2K u_{ij} (f_i^\dagger f_j + f_j^\dagger f_i) + 4J^z (n_i - \frac{1}{2})(n_j - \frac{1}{2}) \right]$$

hopping parameter $t = 2K$
 π flux
nearest-neighbor repulsion $V = 4J^z$
 $f = \frac{1}{2}(c^x + ic^y)$
electron density $f^\dagger f$

Ground-state flux pattern:
[Lieb, PRL '94]

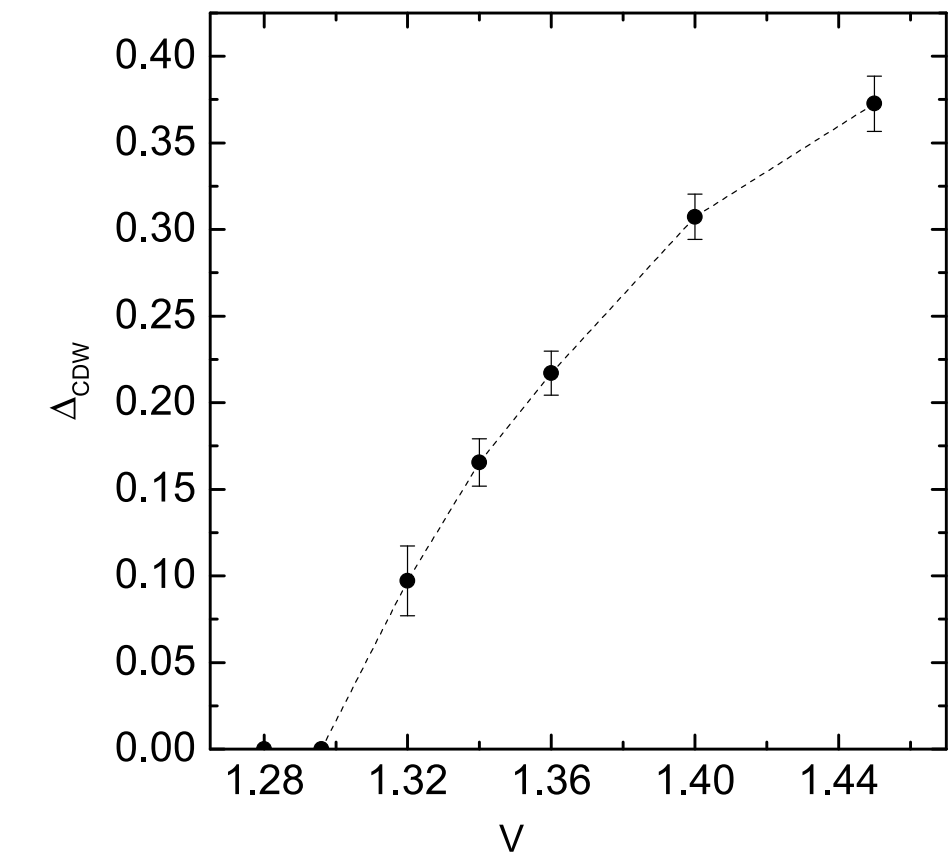
Spin-orbital model \mapsto interacting fermions on π -flux lattice

Spinless fermions on π -flux lattice: QMC



Gross-Neveu- \mathbb{Z}_2 universality:

$$1/\nu = 1.12(1), \quad \eta_\phi = 0.51(3)$$



[Wang, Corboz, Troyer, NJP '14]

[Li, Jiang, Yao, NJP '15]

[Huffman & Chandrasekharan, PRD '17; PRD '20]

[Gracey, IJMP '94]

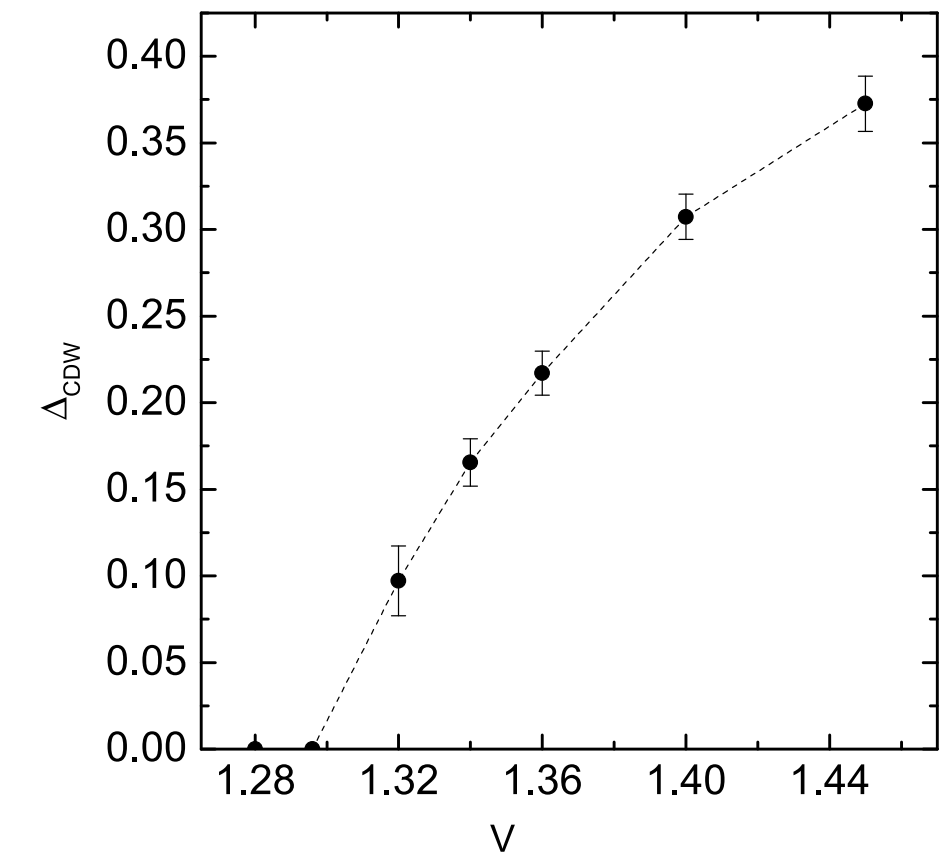
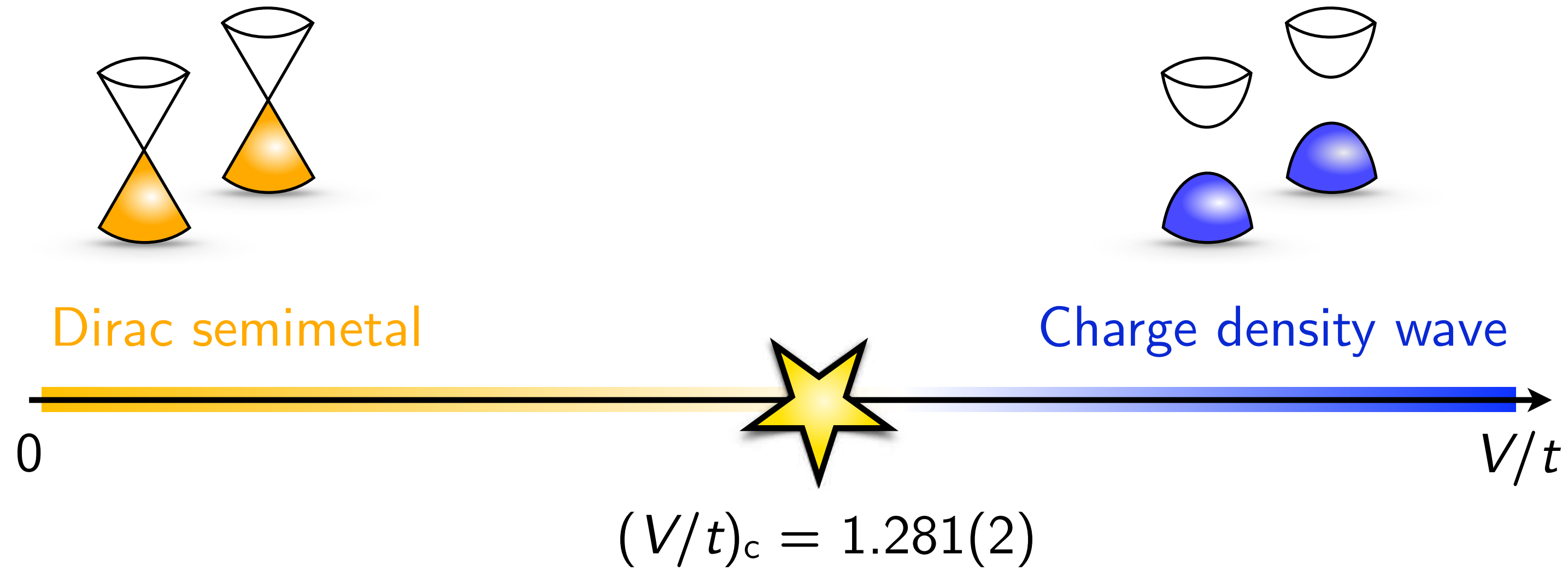
[LJ & Herbut, PRB '14]

[Iliesiu *et al.*, JHEP '18]

[Ihrig, Mihaila, Scherer, PRB '18]

...

Spinless fermions on π -flux lattice: QMC



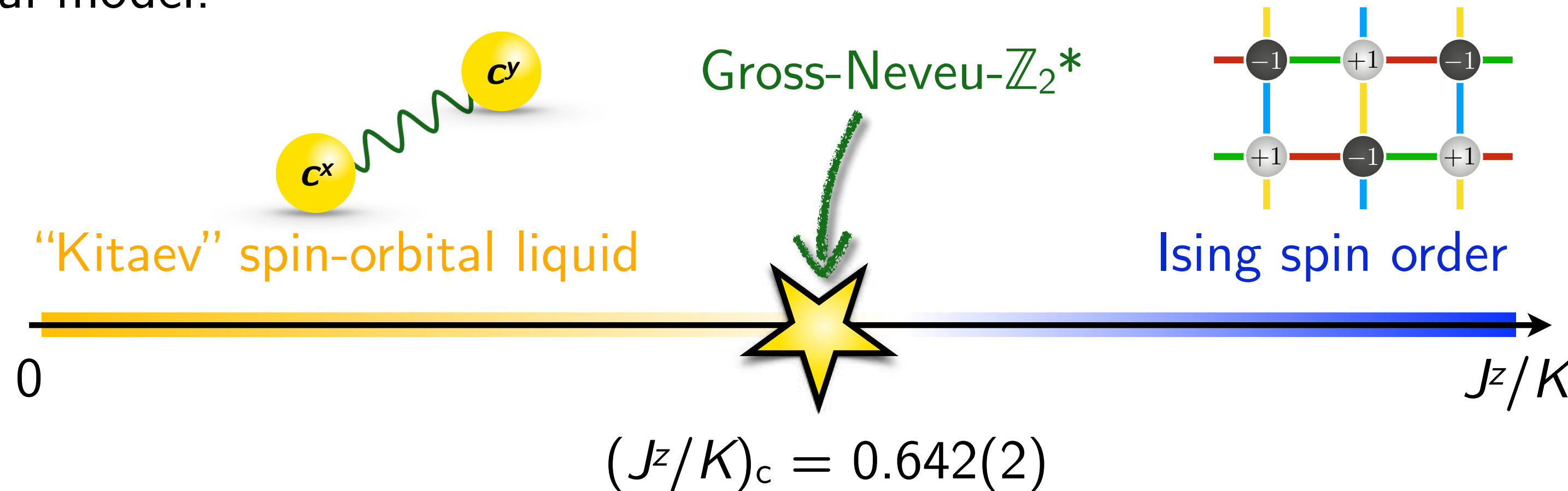
[Wang, Corboz, Troyer, NJP '14]
 [Li, Jiang, Yao, NJP '15]
 [Huffman & Chandrasekharan, PRD '17; PRD '20]

Gross-Neveu- \mathbb{Z}_2 universality:

$$1/\nu = 1.12(1), \quad \eta_\phi = 0.51(3)$$

[Gracey, IJMP '94]
 [LJ & Herbut, PRB '14]
 [Iliesiu *et al.*, JHEP '18]
 [Ihrig, Mihaila, Scherer, PRB '18]

Spin-orbital model:



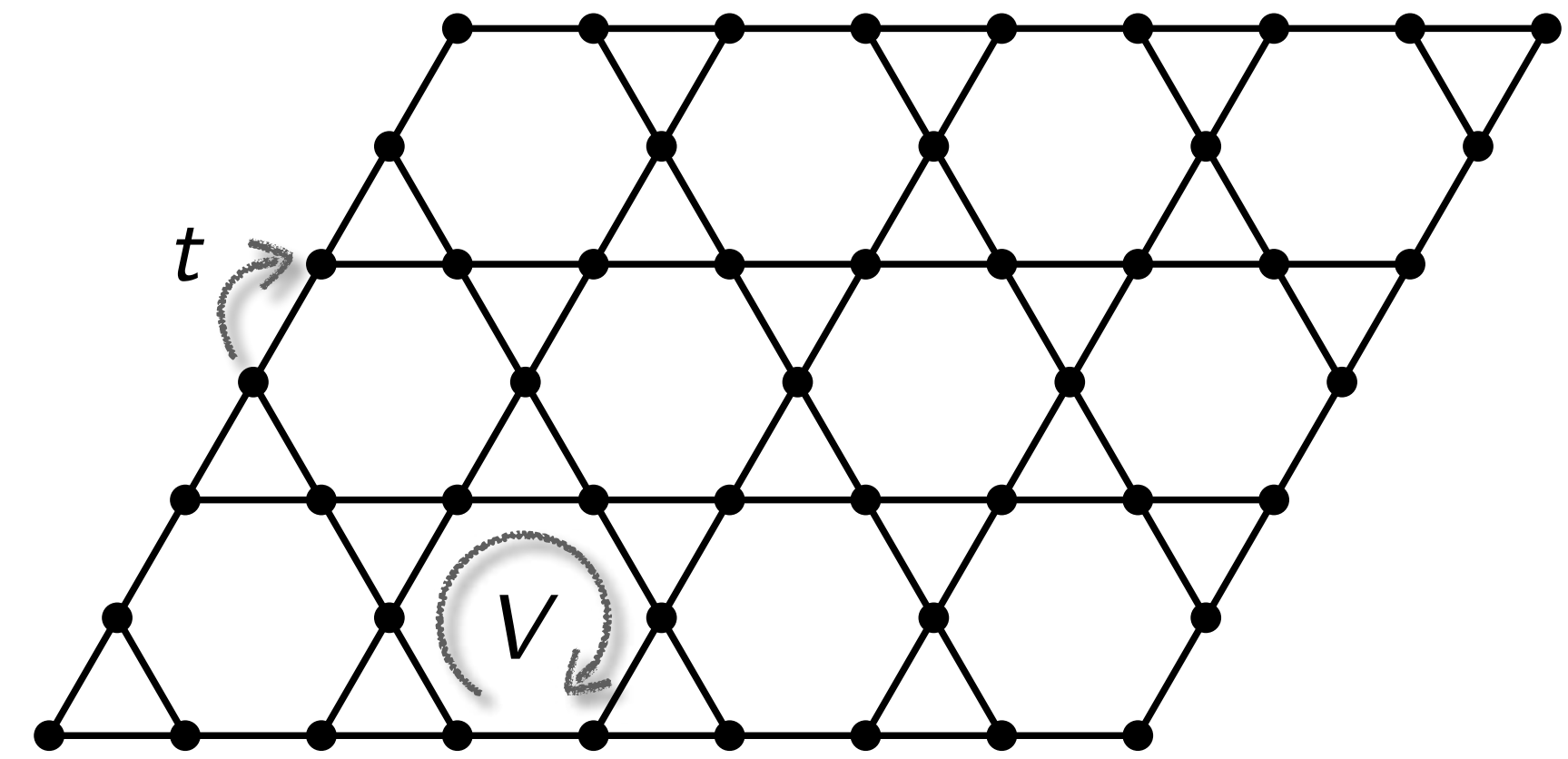
[Seifert, Dong, Chulliparambil, Vojta, Tu, LJ, PRL '20]

Example: Kagome-lattice Bose-Hubbard model

Hamiltonian:

$$\mathcal{H} = -t \sum_{\langle ij \rangle} \left[b_i^\dagger b_j + b_i b_j^\dagger \right] + V \sum_{\hexagon} (n_{\hexagon})^2$$

... b_i hard-core bosons

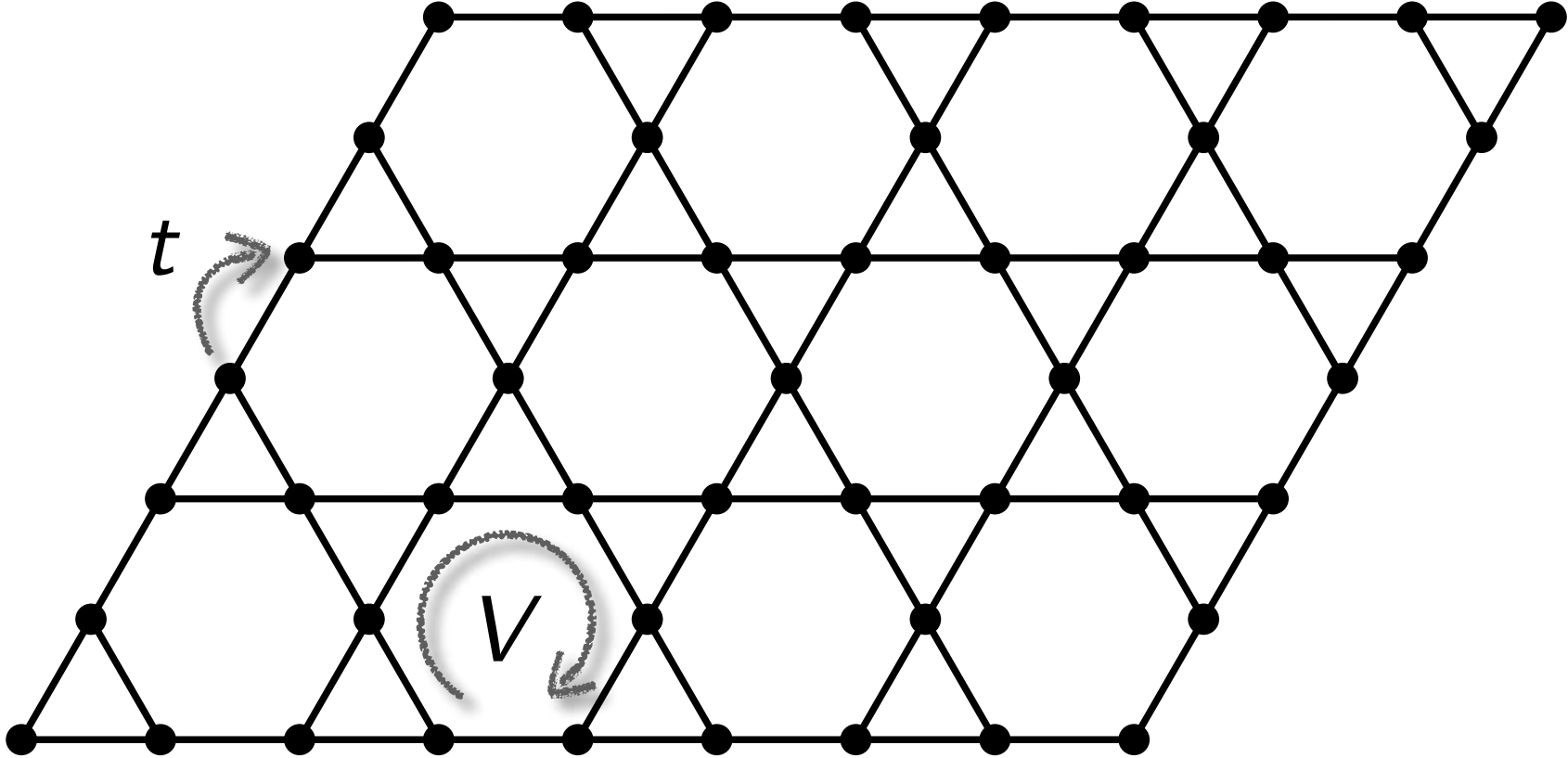


Example: Kagome-lattice Bose-Hubbard model

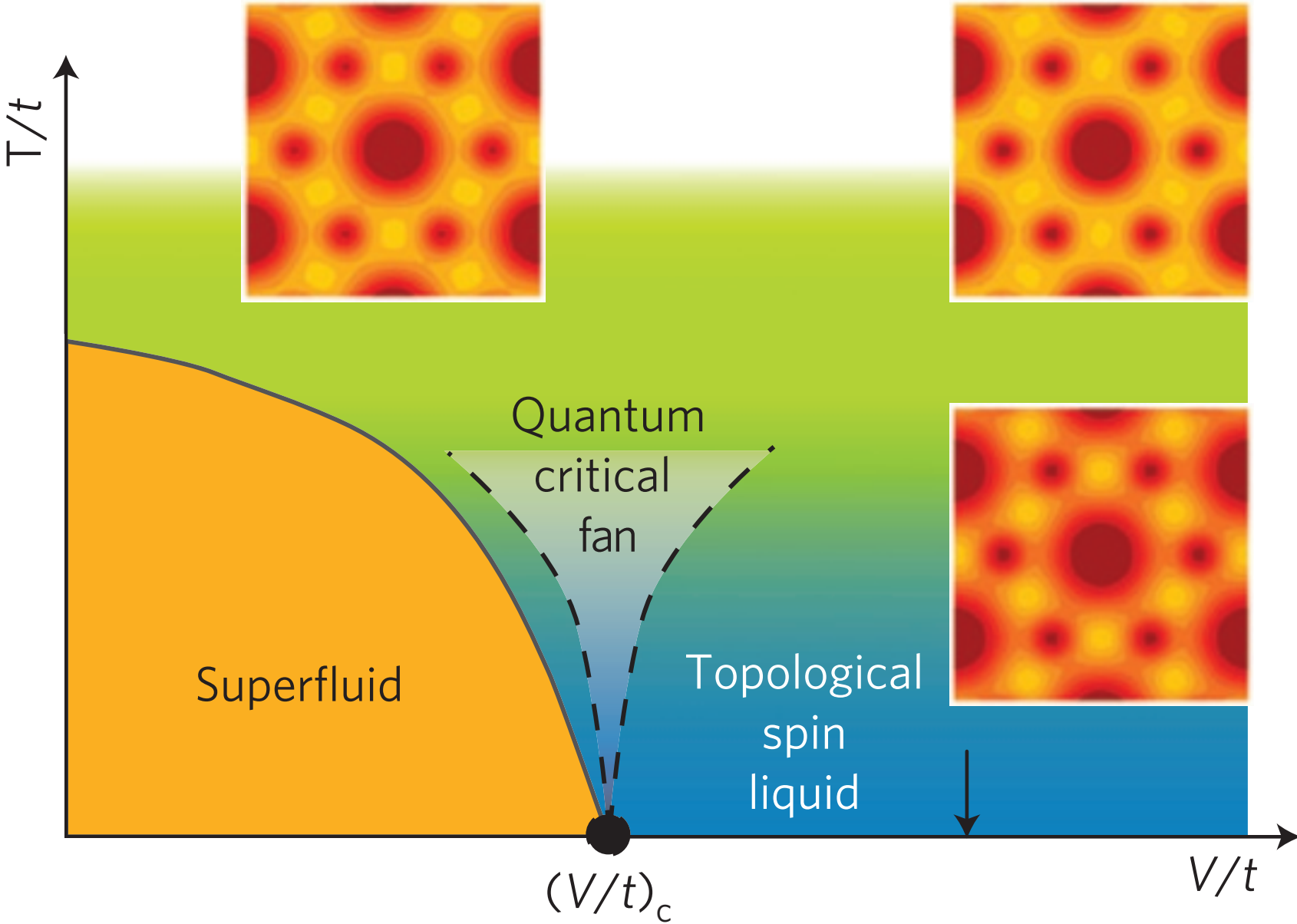
Hamiltonian:

$$\mathcal{H} = -t \sum_{\langle ij \rangle} \left[b_i^\dagger b_j + b_i b_j^\dagger \right] + V \sum_{\hexagon} (n_{\hexagon})^2$$

... b_i hard-core bosons



Phase diagram:



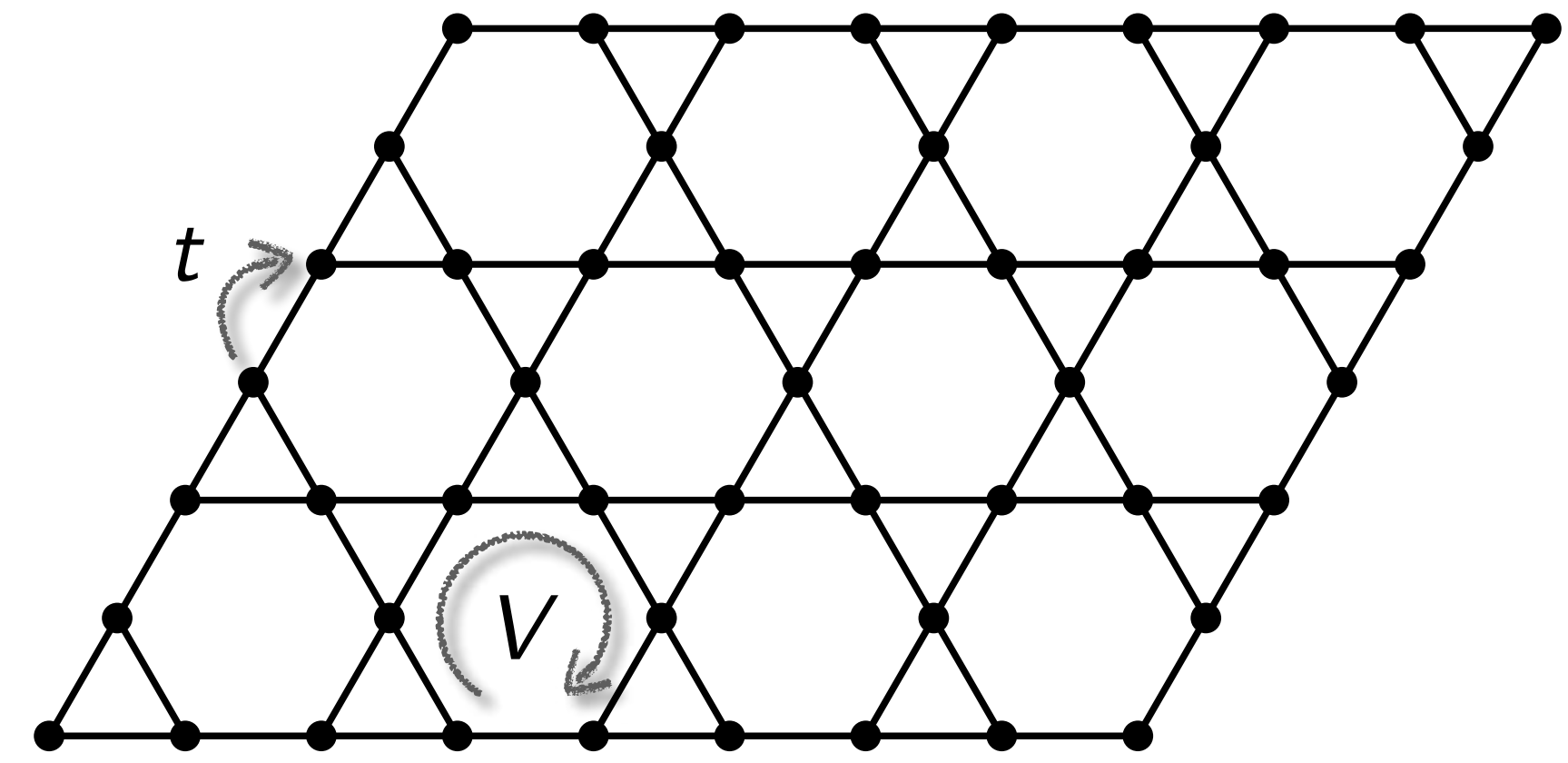
[Isakov, Hastings, Melko, Nat. Phys. '11]

Example: Kagome-lattice Bose-Hubbard model

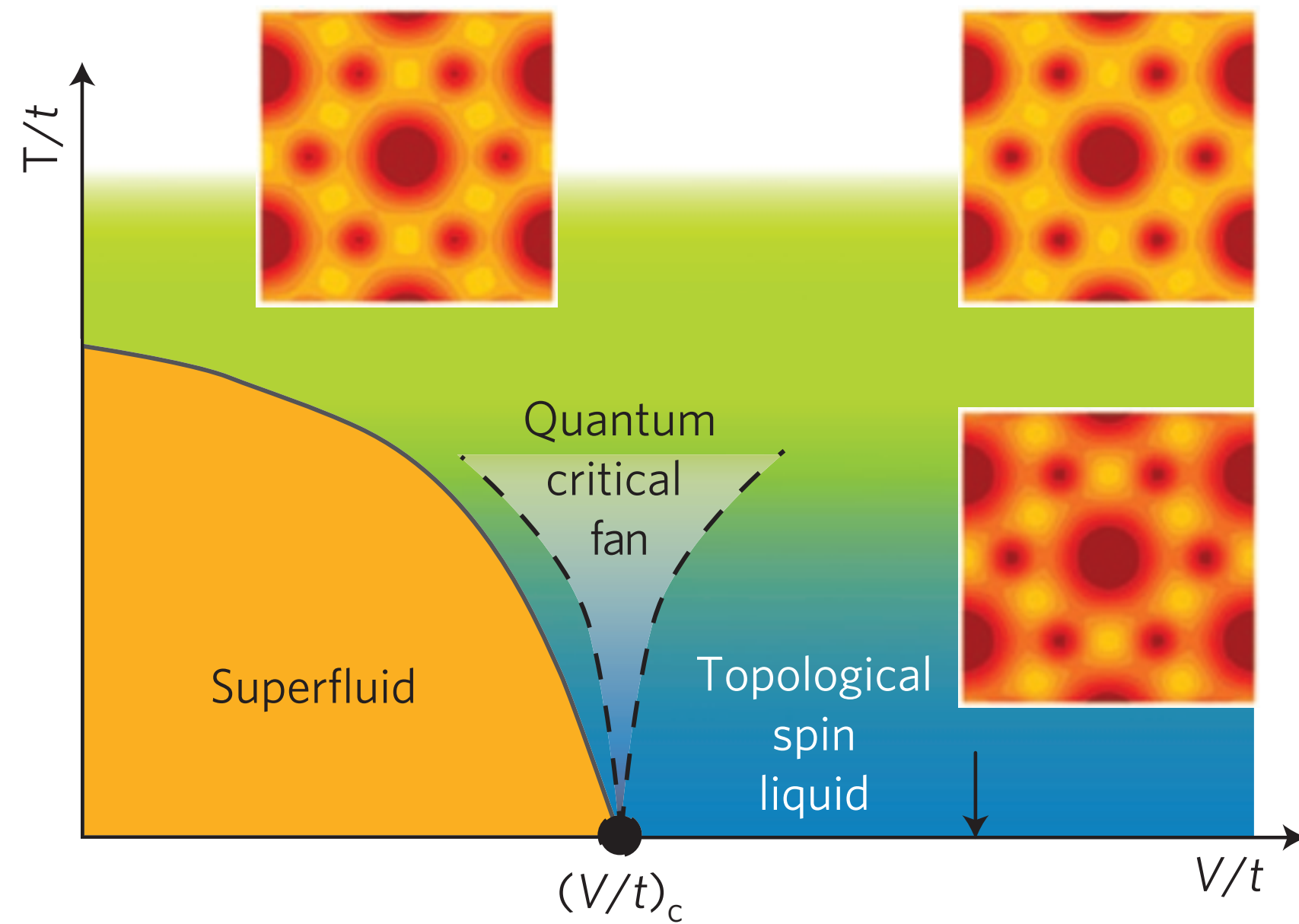
Hamiltonian:

$$\mathcal{H} = -t \sum_{\langle ij \rangle} \left[b_i^\dagger b_j + b_i b_j^\dagger \right] + V \sum_{\hexagon} (n_{\hexagon})^2$$

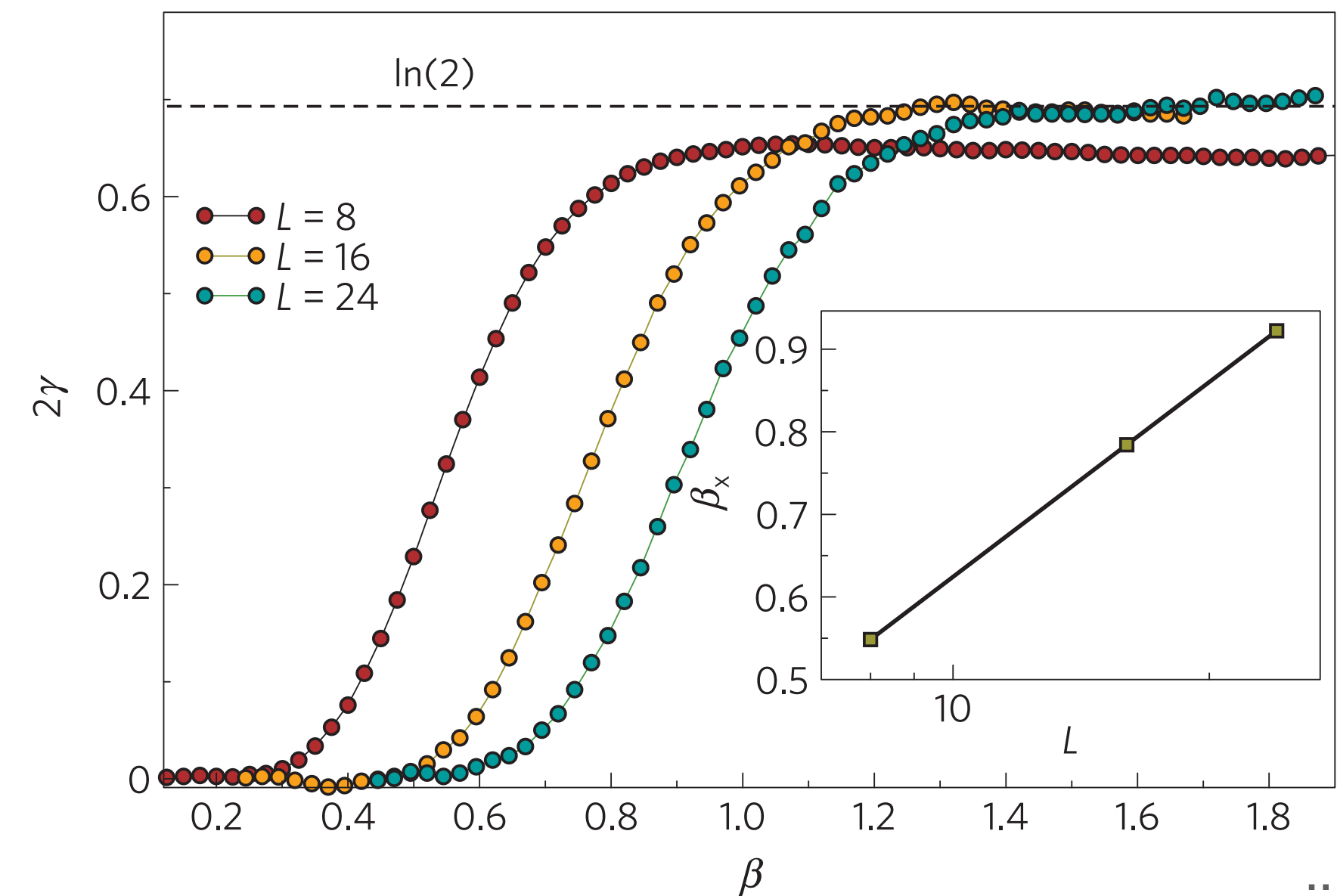
... b_i hard-core bosons



Phase diagram:



Entanglement entropy: $S_n(A) = a\ell - \gamma + \dots$

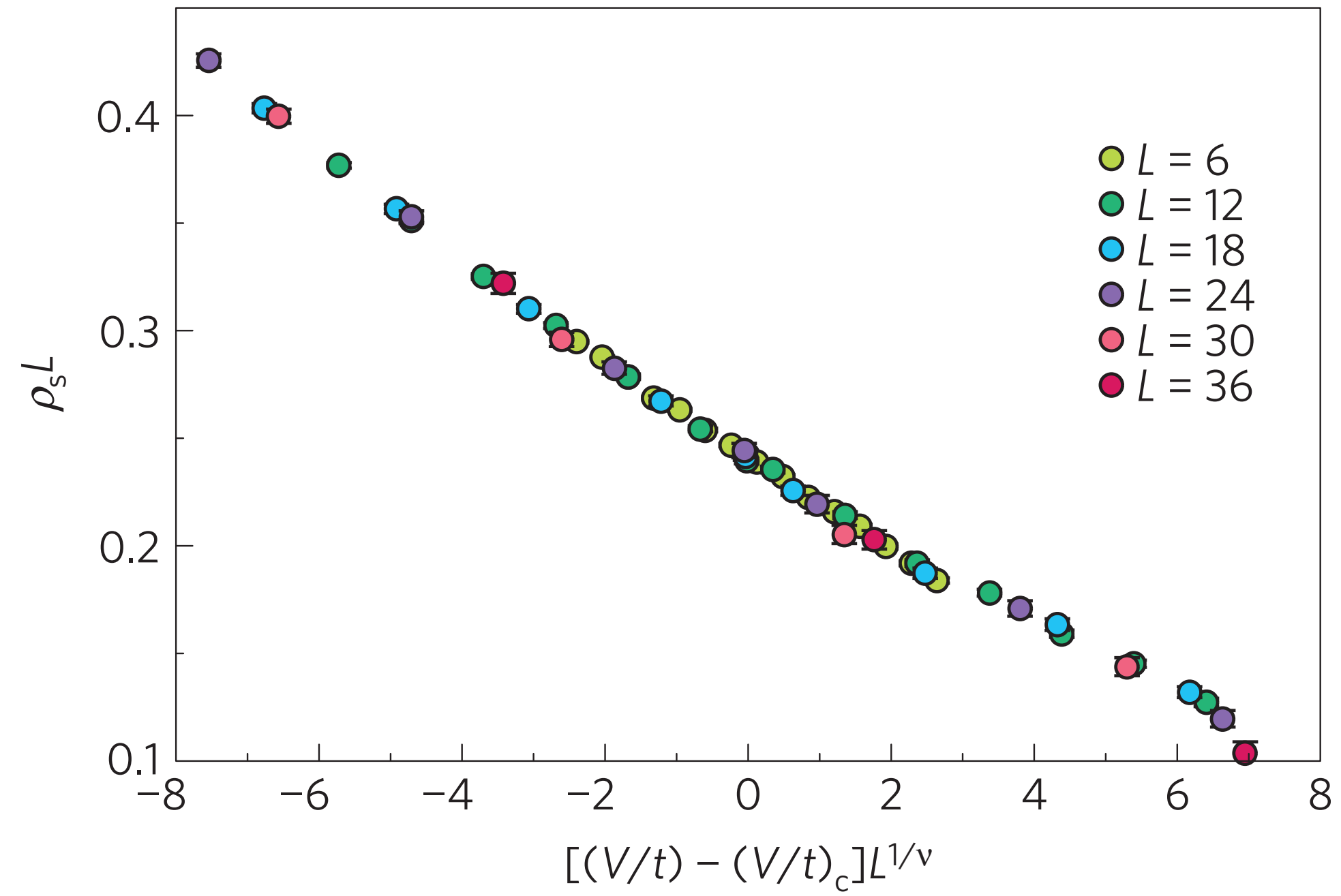


... in spin liquid phase

[Isakov, Hastings, Melko, Nat. Phys. '11]

Quantum critical scaling: XY*

Superfluid density:

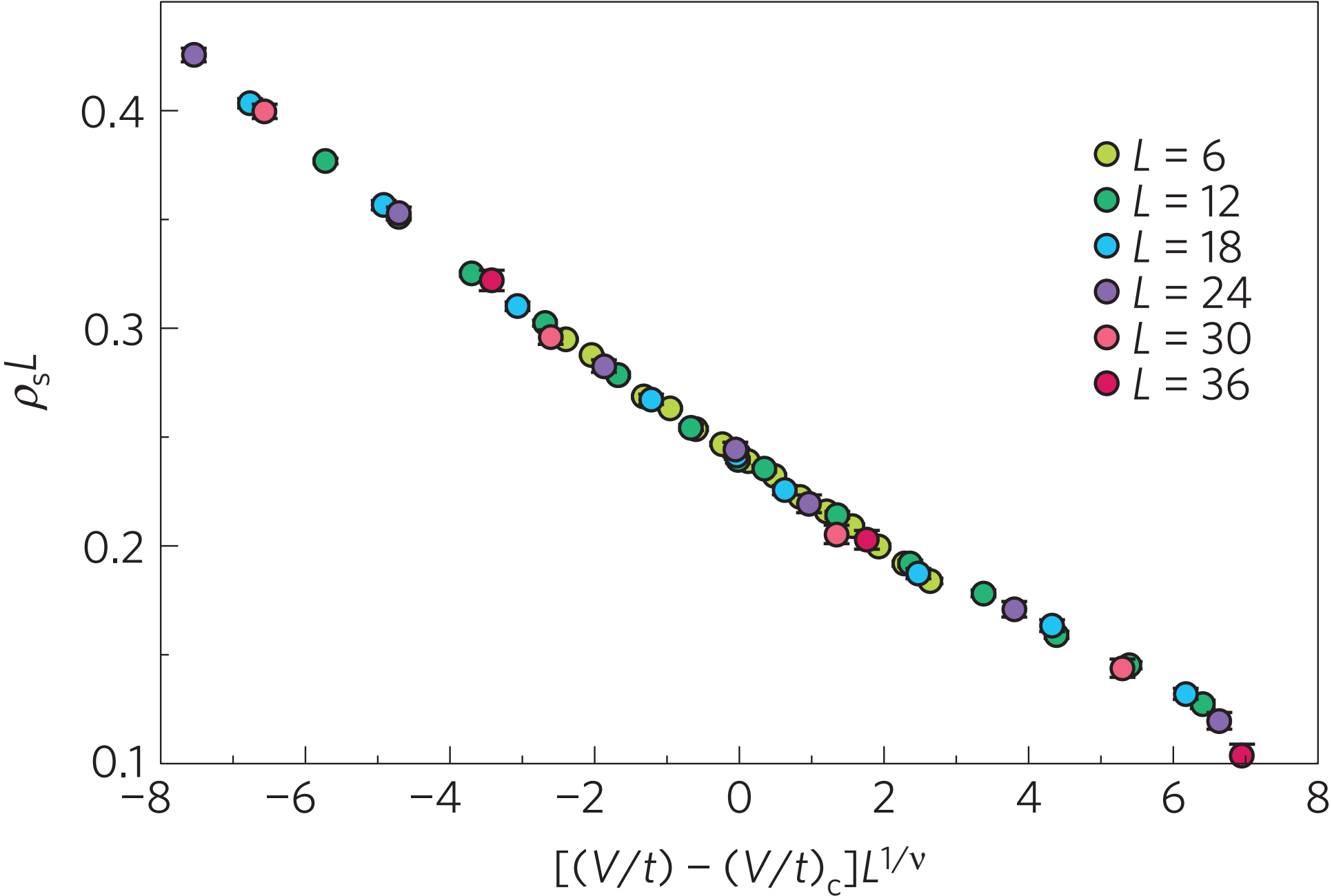


[Isakov, Hastings, Melko, Nat. Phys. '11]

$$\nu \approx 0.67 = \nu_{XY}$$

Quantum critical scaling: XY*

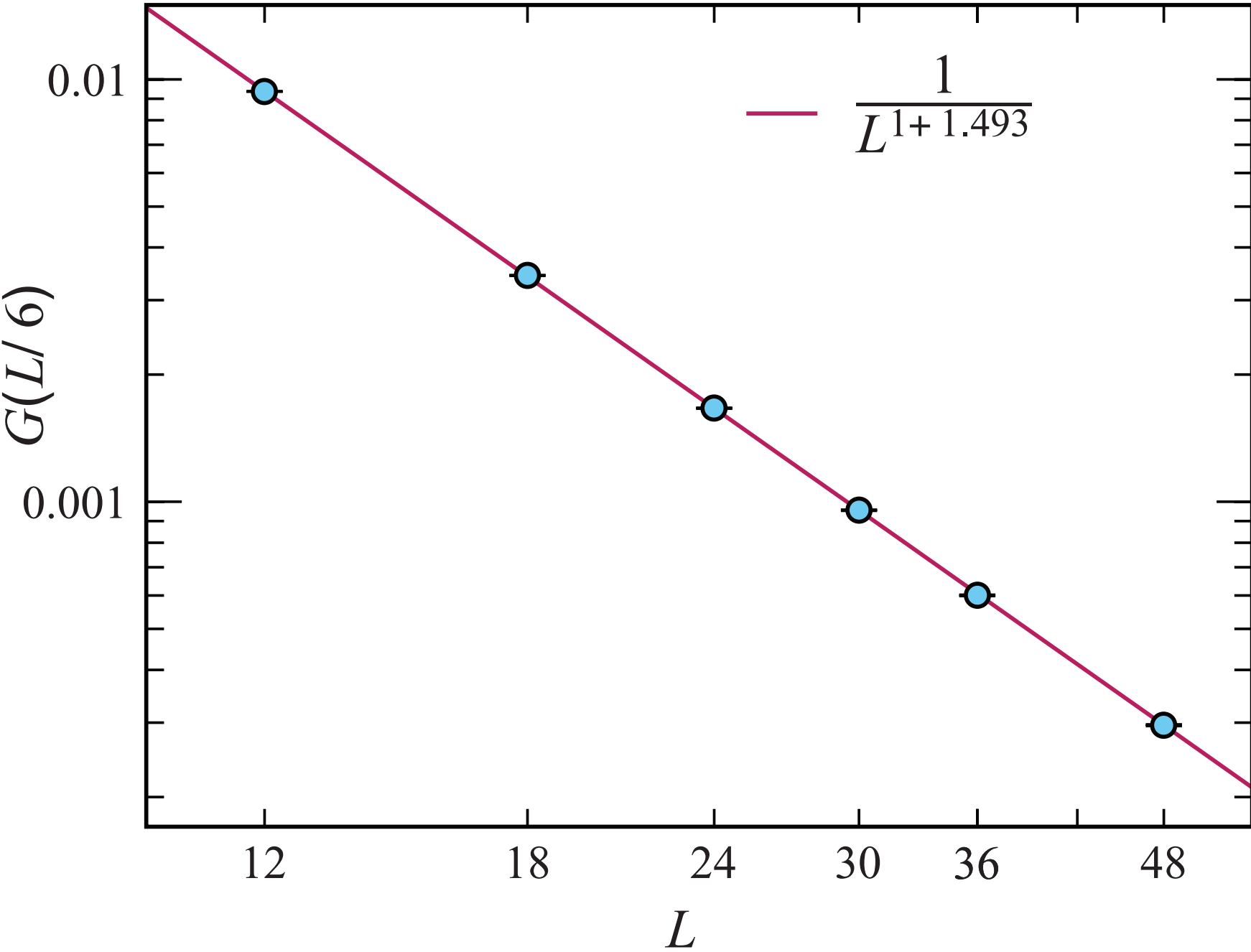
Superfluid density:



[Isakov, Hastings, Melko, Nat. Phys. '11]

$$\nu \approx 0.67 = \nu_{XY}$$

Two-point superfluid correlator:



[Isakov, Melko, Hastings, Science '12]

$$\eta \approx 1.49 \neq \eta_{XY} \approx 0.038$$

Order parameter *composite* of fractionalized particles!

... cf. $\eta_T \approx 1.47$ from XY field theory
[Calabrese, Pelissetto, Vicari, PRE '02]

Finite-size spectroscopy: Ising vs Ising*

Transverse-field Ising:

$$\mathcal{H} = -J \sum_{\langle ij \rangle} \sigma_i^z \sigma_j^z - h \sum_i \sigma_i^x$$

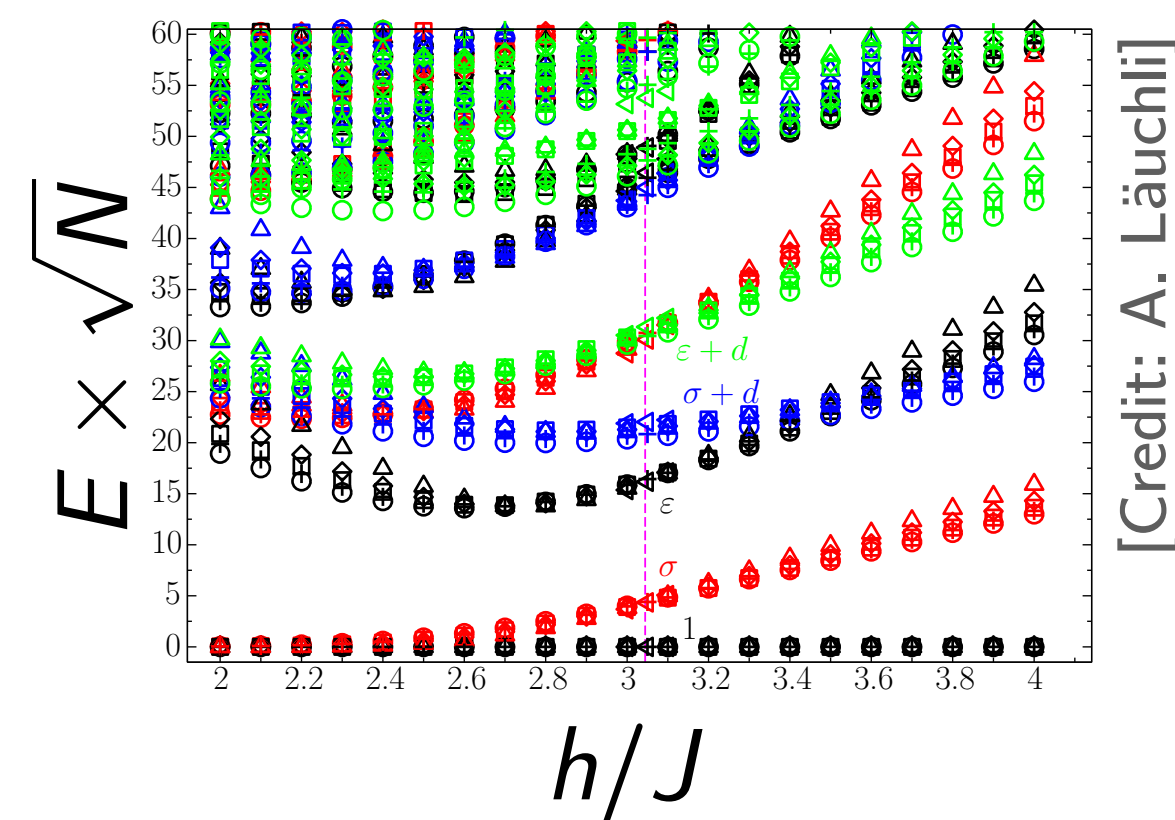
Transverse-field toric code:

$$\mathcal{H} = -J \sum_s \prod_{i \in s} \sigma_i^x - J \sum_p \prod_{i \in p} \sigma_i^z - h \sum_i \sigma_i^x$$

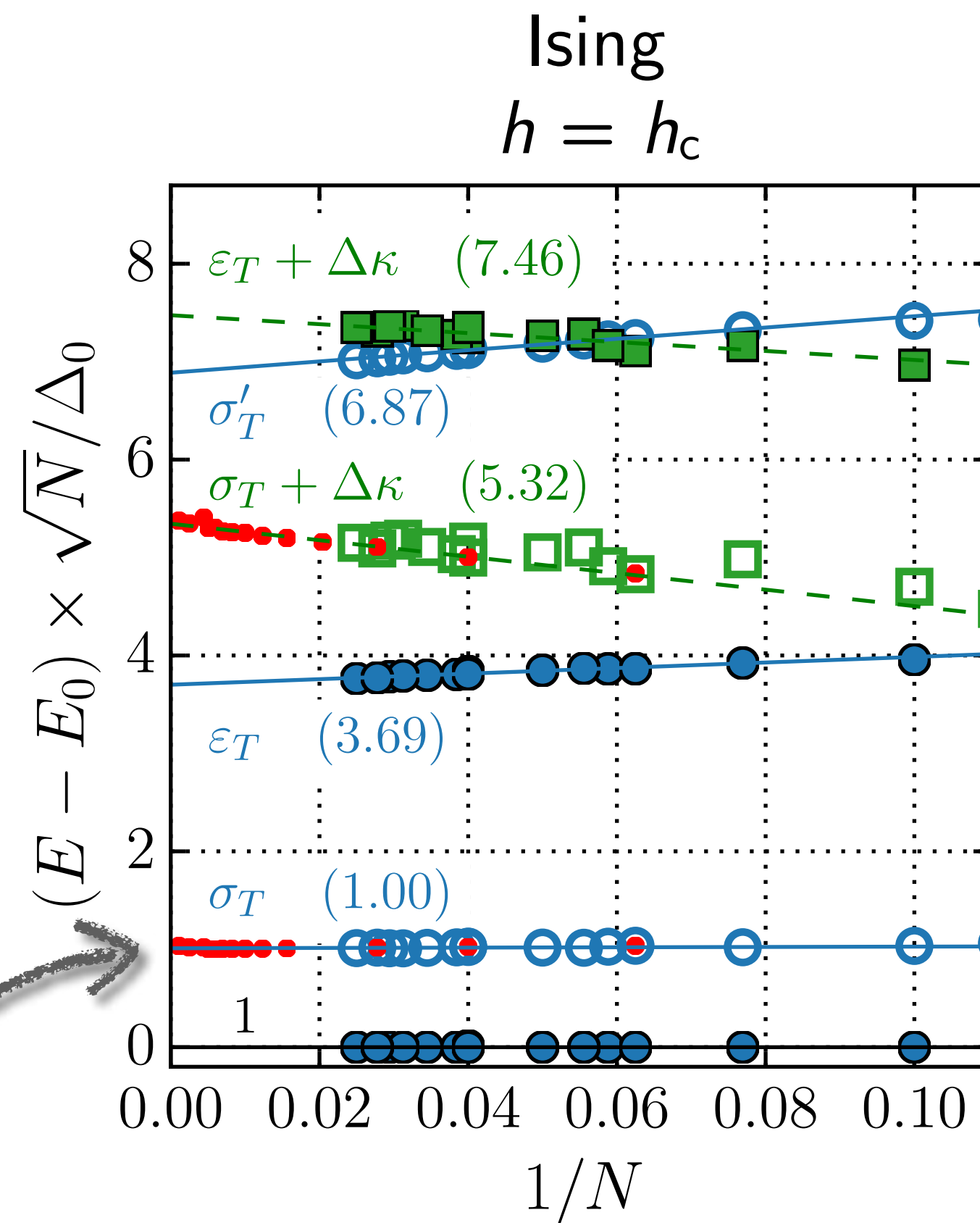
Finite-size spectroscopy: Ising vs Ising*

Transverse-field Ising:

$$\mathcal{H} = -J \sum_{\langle ij \rangle} \sigma_i^z \sigma_j^z - h \sum_i \sigma_i^x$$



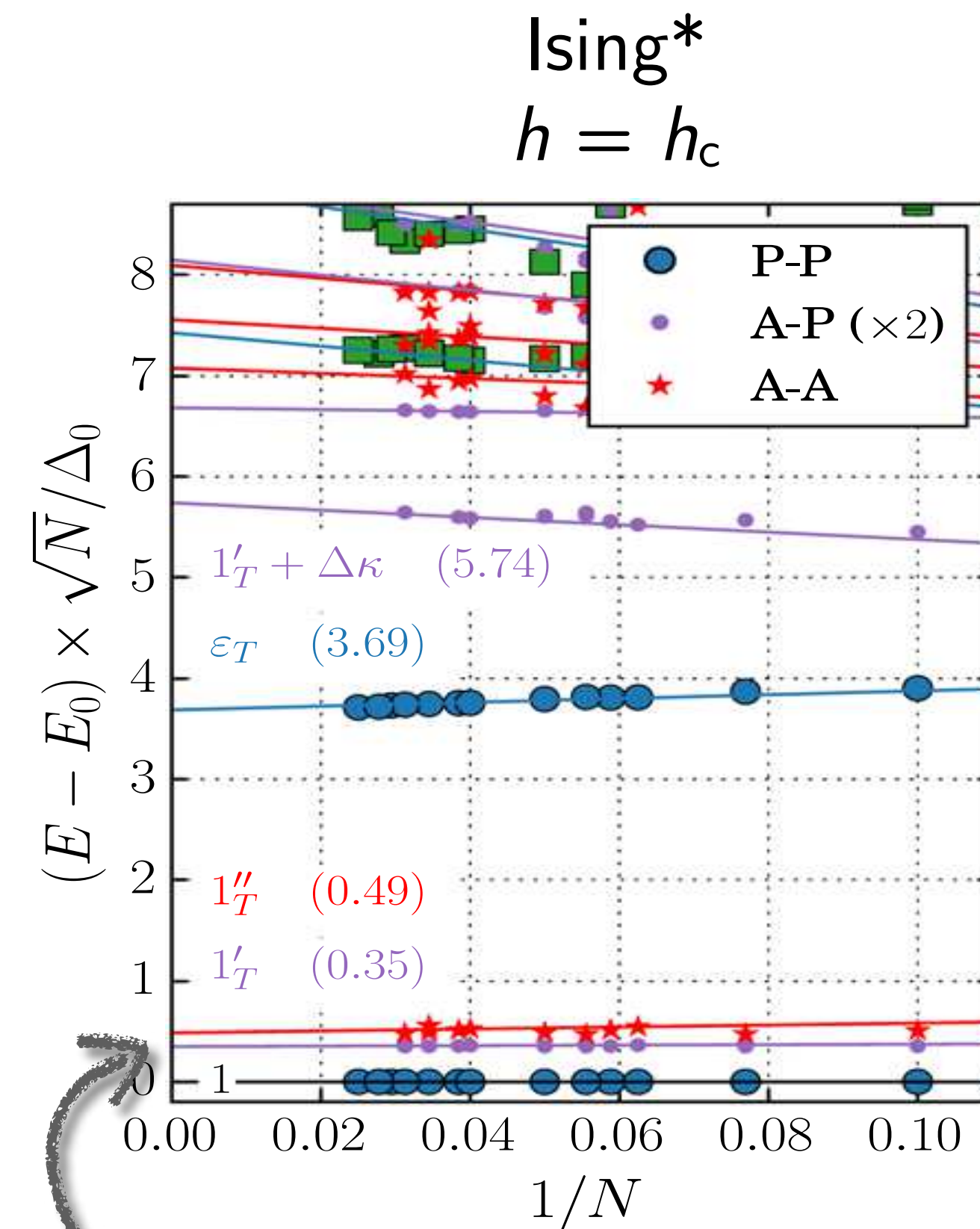
[Credit: A. Läuchli]



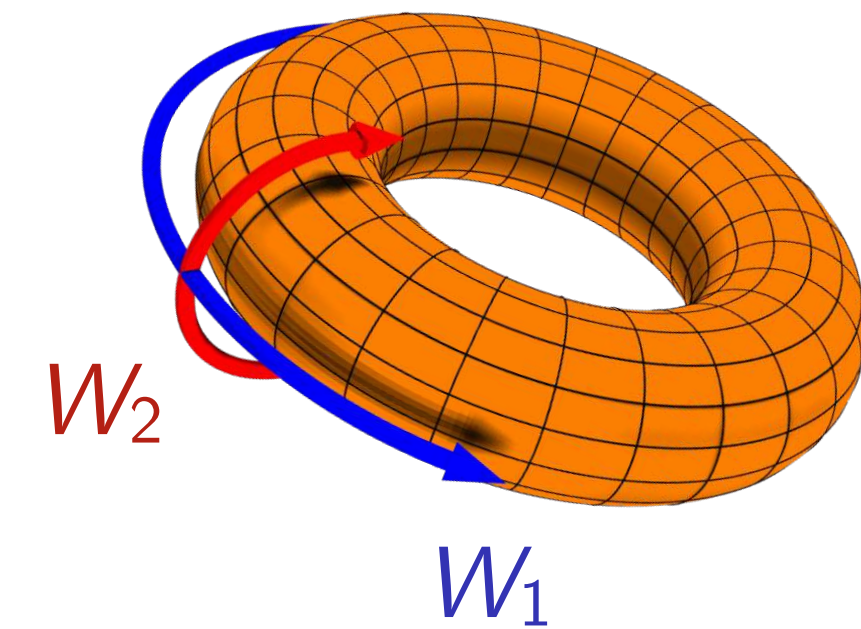
missing in Ising*

Transverse-field toric code:

$$\mathcal{H} = -J \sum_s \prod_{i \in s} \sigma_i^x - J \sum_p \prod_{i \in p} \sigma_i^z - h \sum_i \sigma_i^x$$



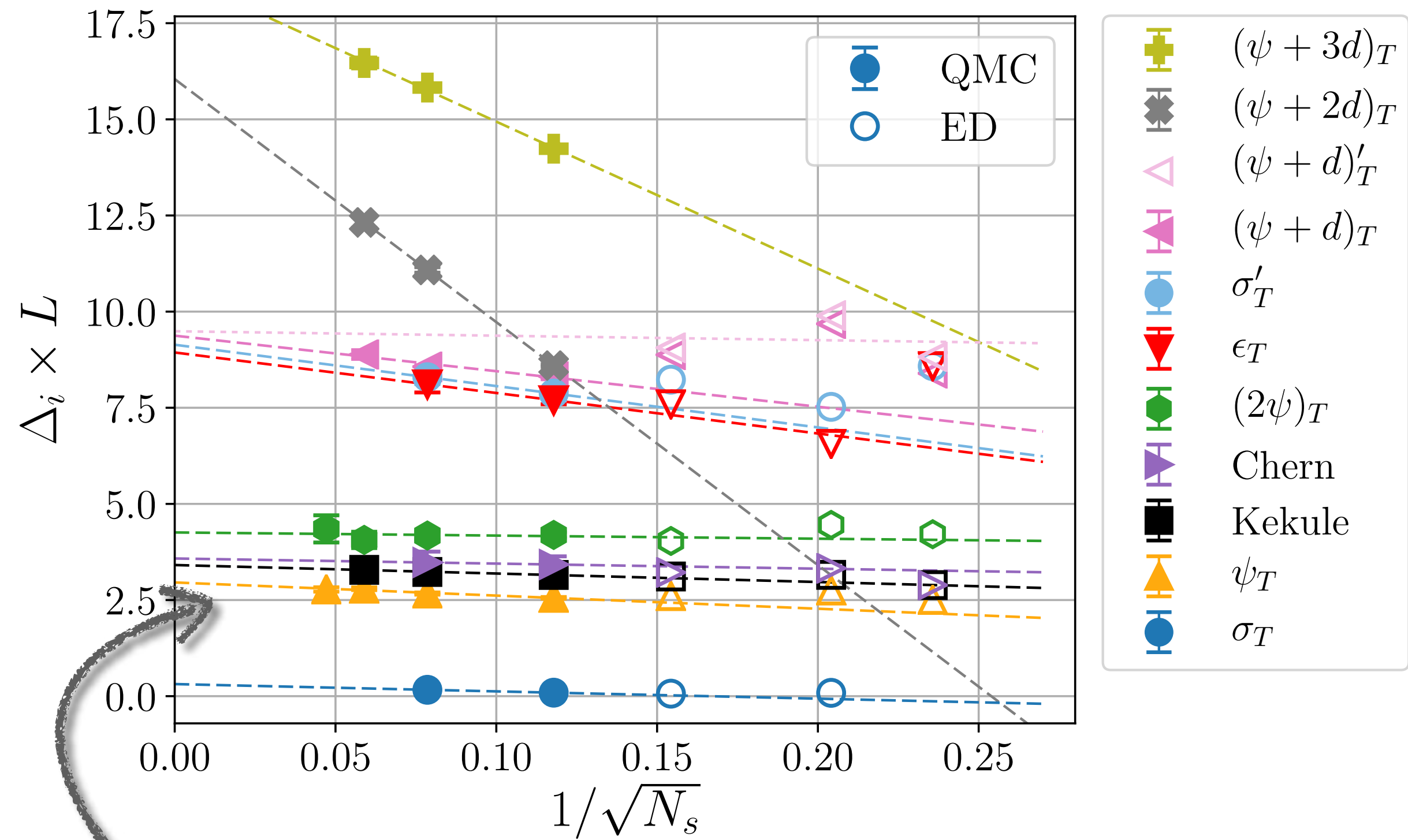
topological "copies"



[Schuler, Whitsitt, Henry, Sachdev, Läuchli, PRL '16]

Gross-Neveu vs Gross-Neveu*

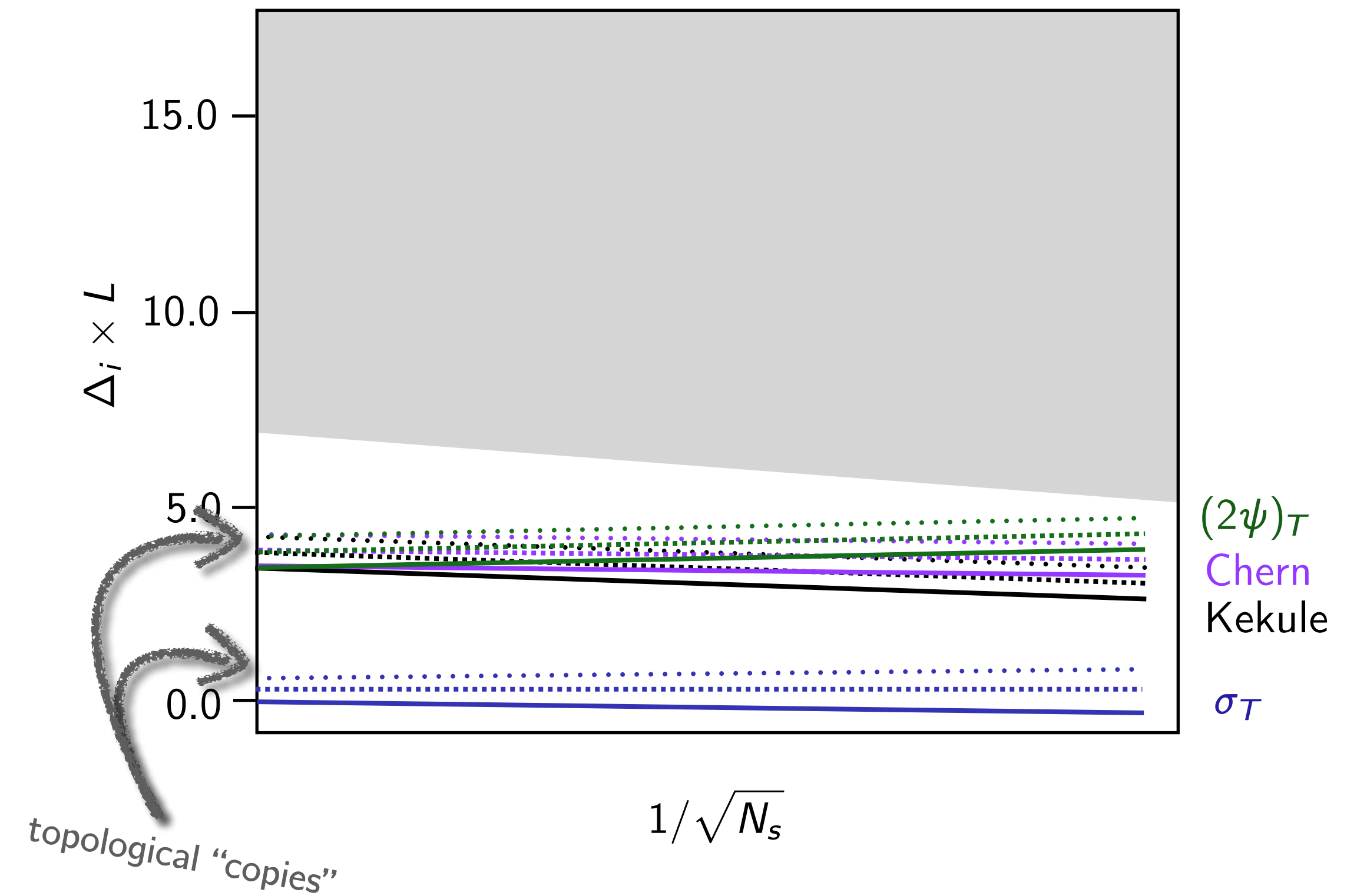
Gross-Neveu- \mathbb{Z}_2



missing in GN*

[Schuler, Hesselmann, Whitsitt, Lang, Wessel, Läuchli, PRB '21]

Gross-Neveu- \mathbb{Z}_2^* (schematic)



... testable in future simulations

Gross-Neveu-SO(3) criticality

Field theory:

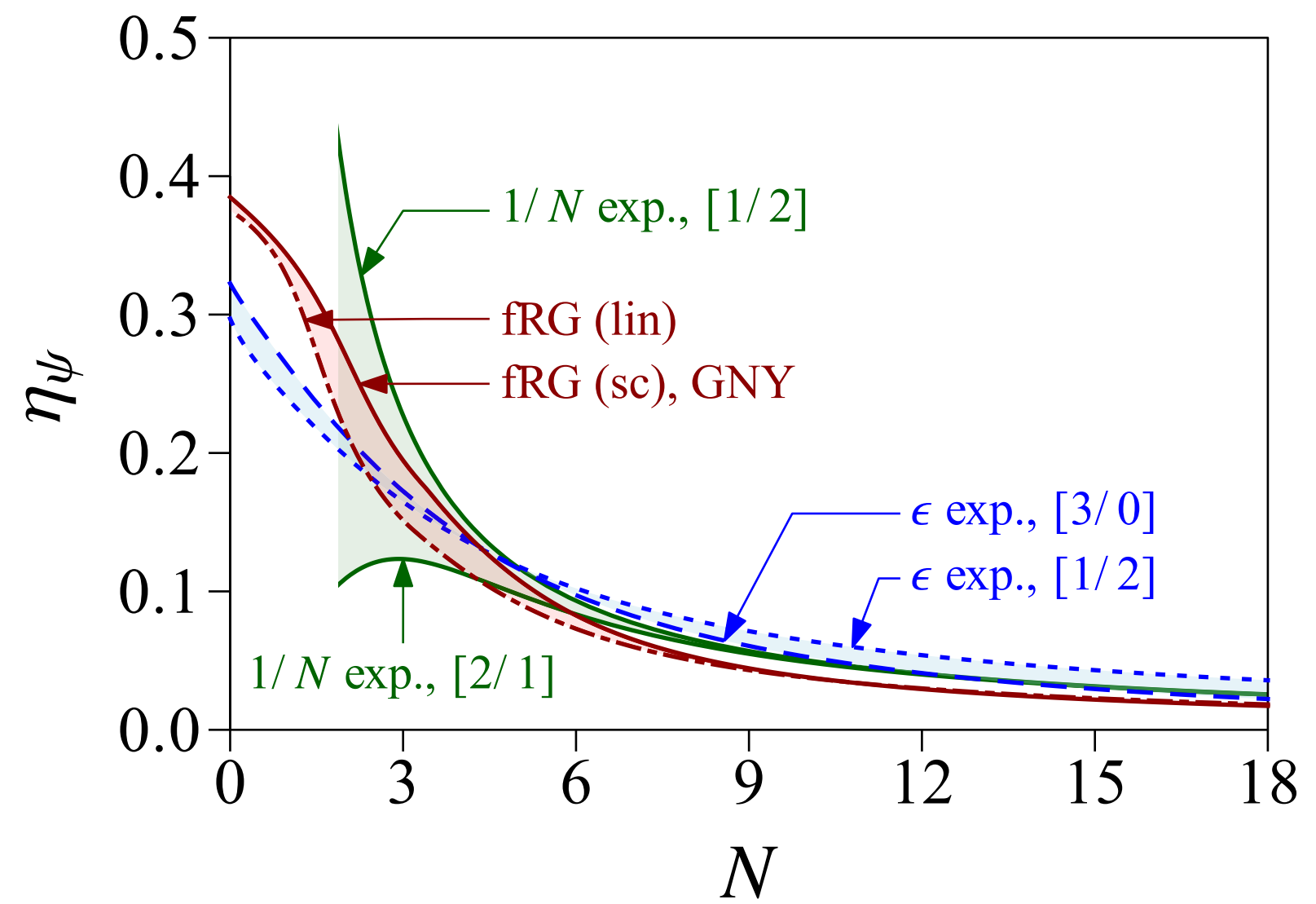
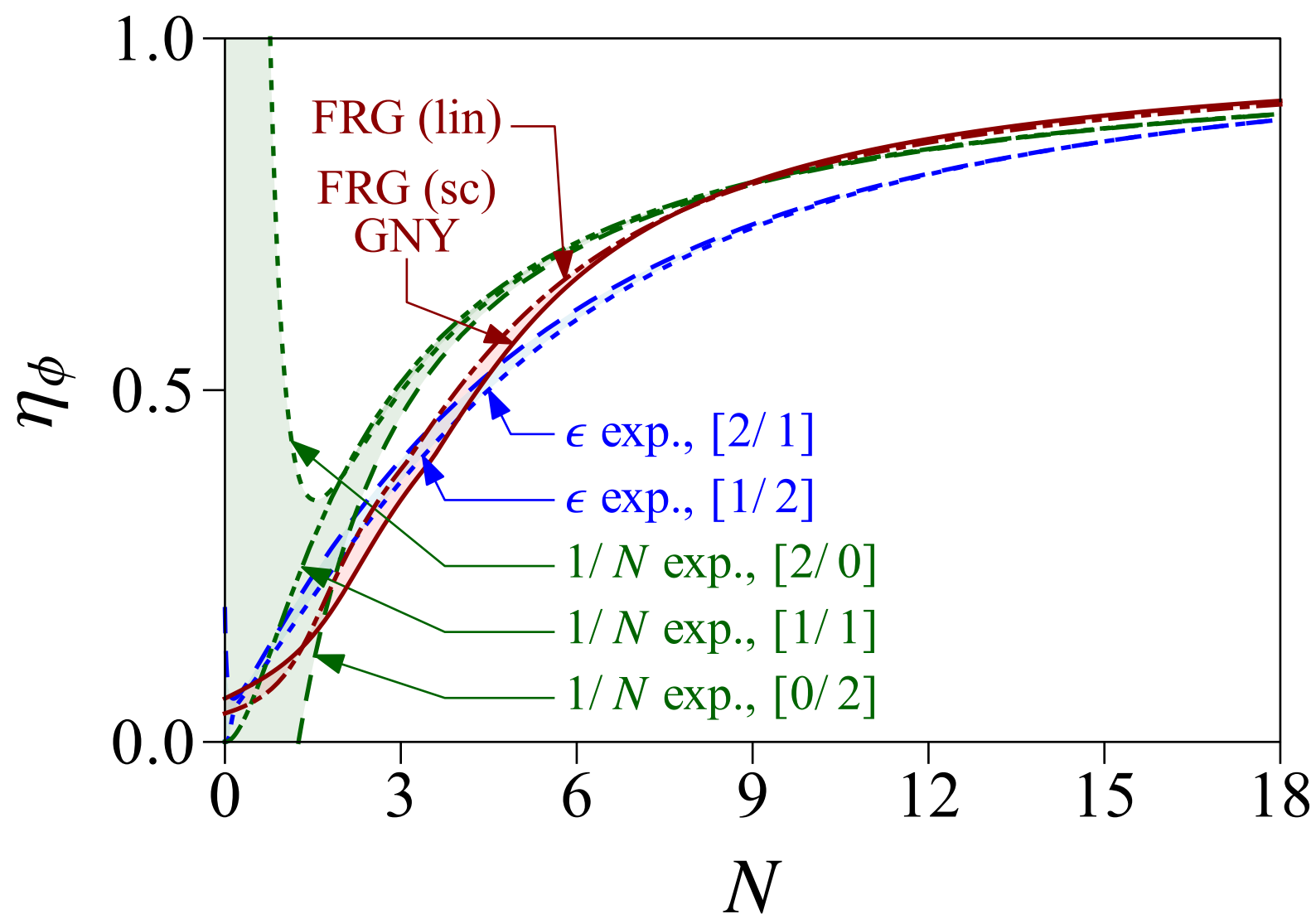
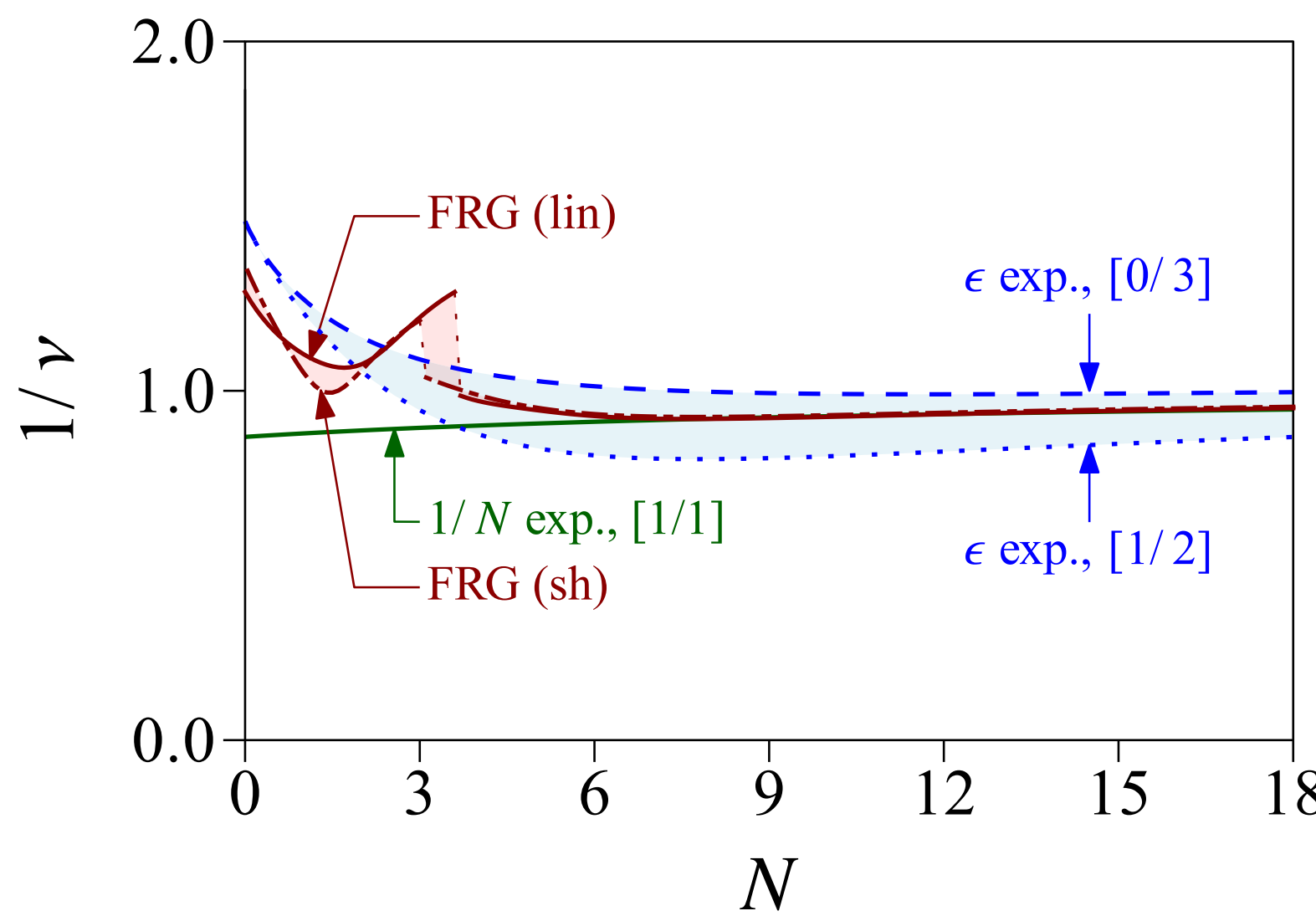
$$\mathcal{S} = \int d^2\vec{x}d\tau \left[\bar{\psi}\gamma^\mu\partial_\mu\psi + g\vec{\varphi} \cdot \bar{\psi}(\mathbb{1}_2 \otimes \vec{L})\psi + \frac{1}{2}\vec{\varphi}(-\partial_\mu^2 + m^2)\vec{\varphi} + \lambda(\vec{\varphi} \cdot \vec{\varphi})^2 \right]$$

Gross-Neveu-SO(3) criticality

Field theory:

$$\mathcal{S} = \int d^2\vec{x}d\tau \left[\bar{\psi}\gamma^\mu\partial_\mu\psi + g\vec{\varphi} \cdot \bar{\psi}(\mathbb{1}_2 \otimes \vec{L})\psi + \frac{1}{2}\vec{\varphi}(-\partial_\mu^2 + m^2)\vec{\varphi} + \lambda(\vec{\varphi} \cdot \vec{\varphi})^2 \right]$$

Critical exponents:



... from three-loop ϵ expansion,
second/third-order $1/N$ expansion,
functional RG in local potential approximation

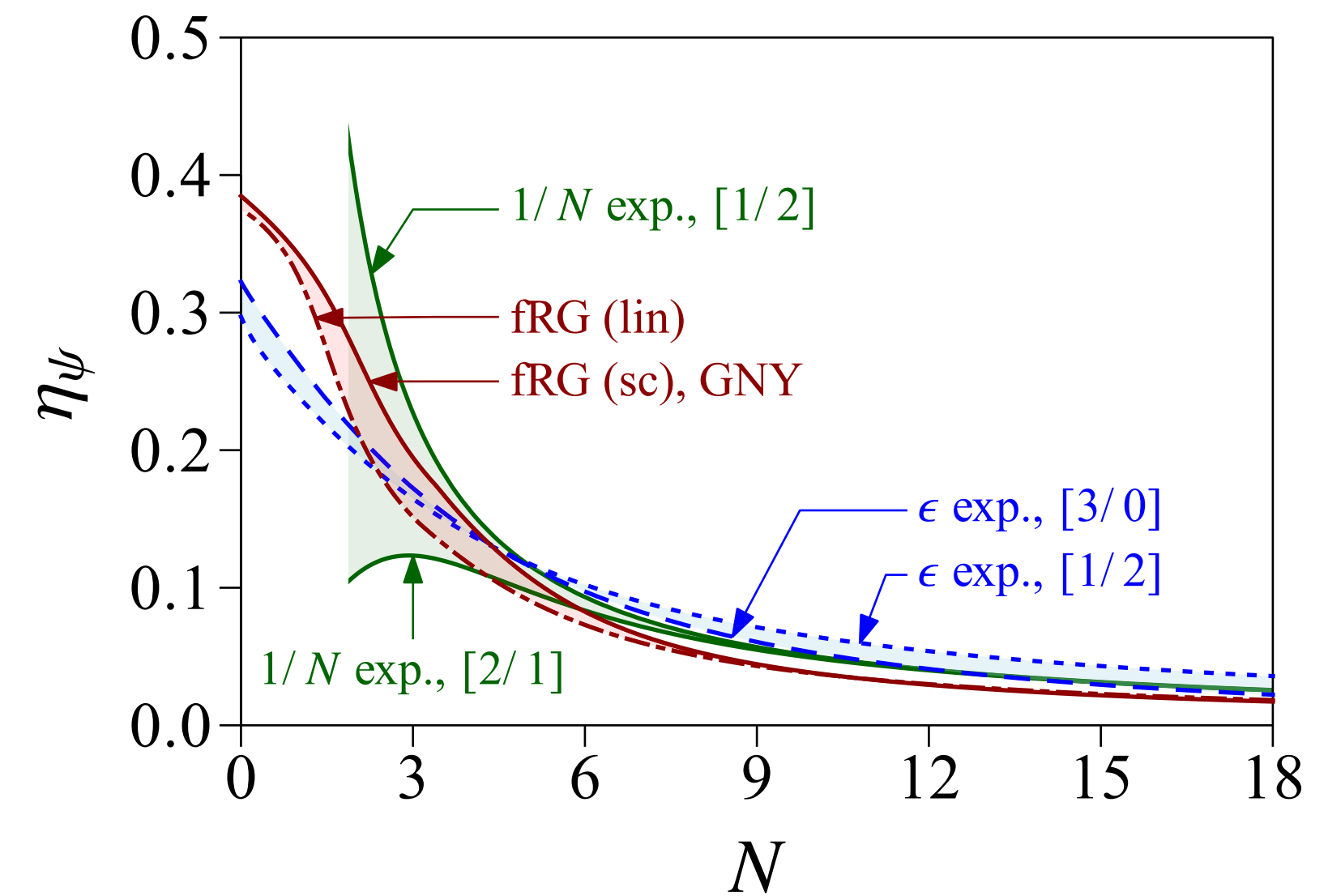
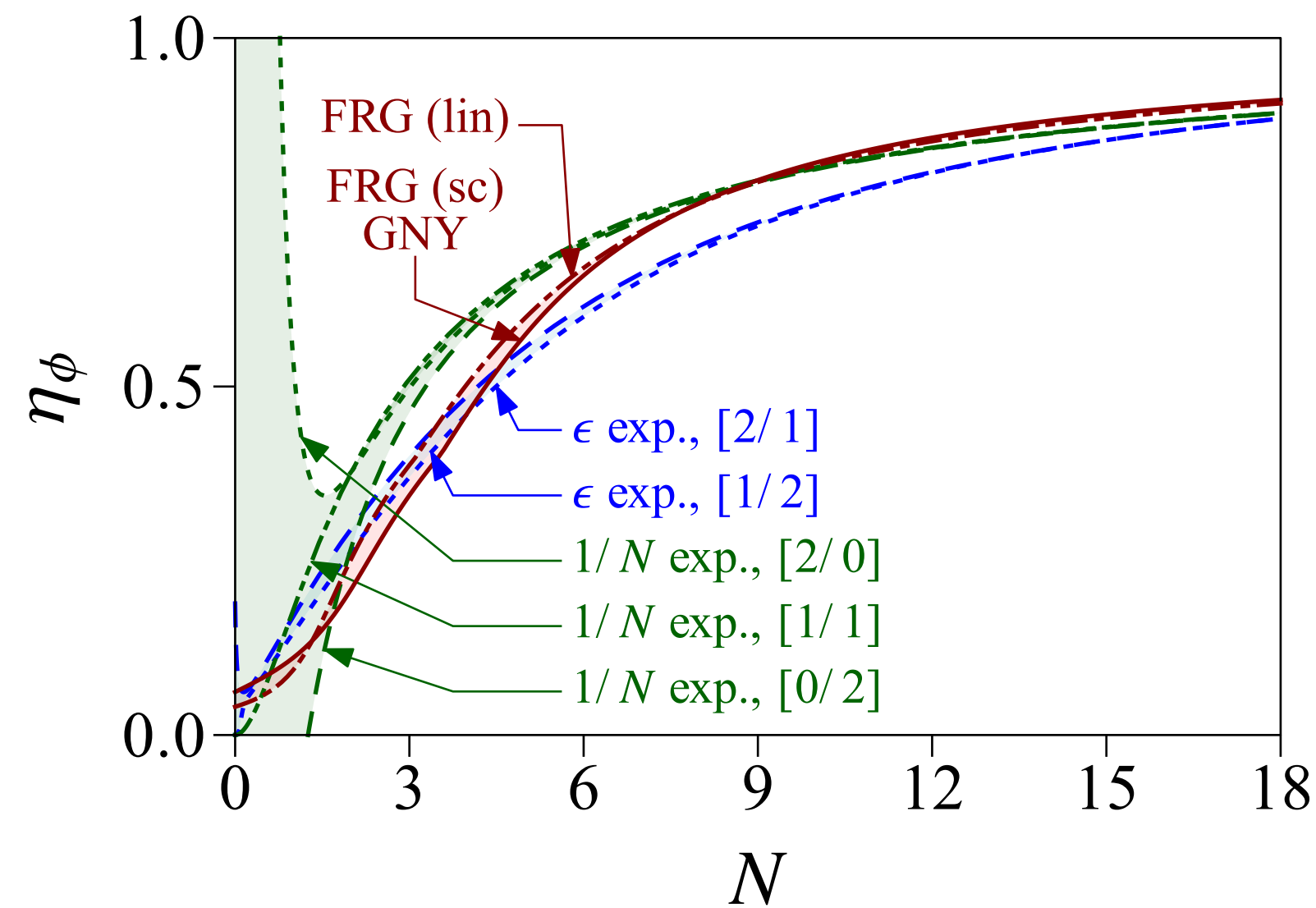
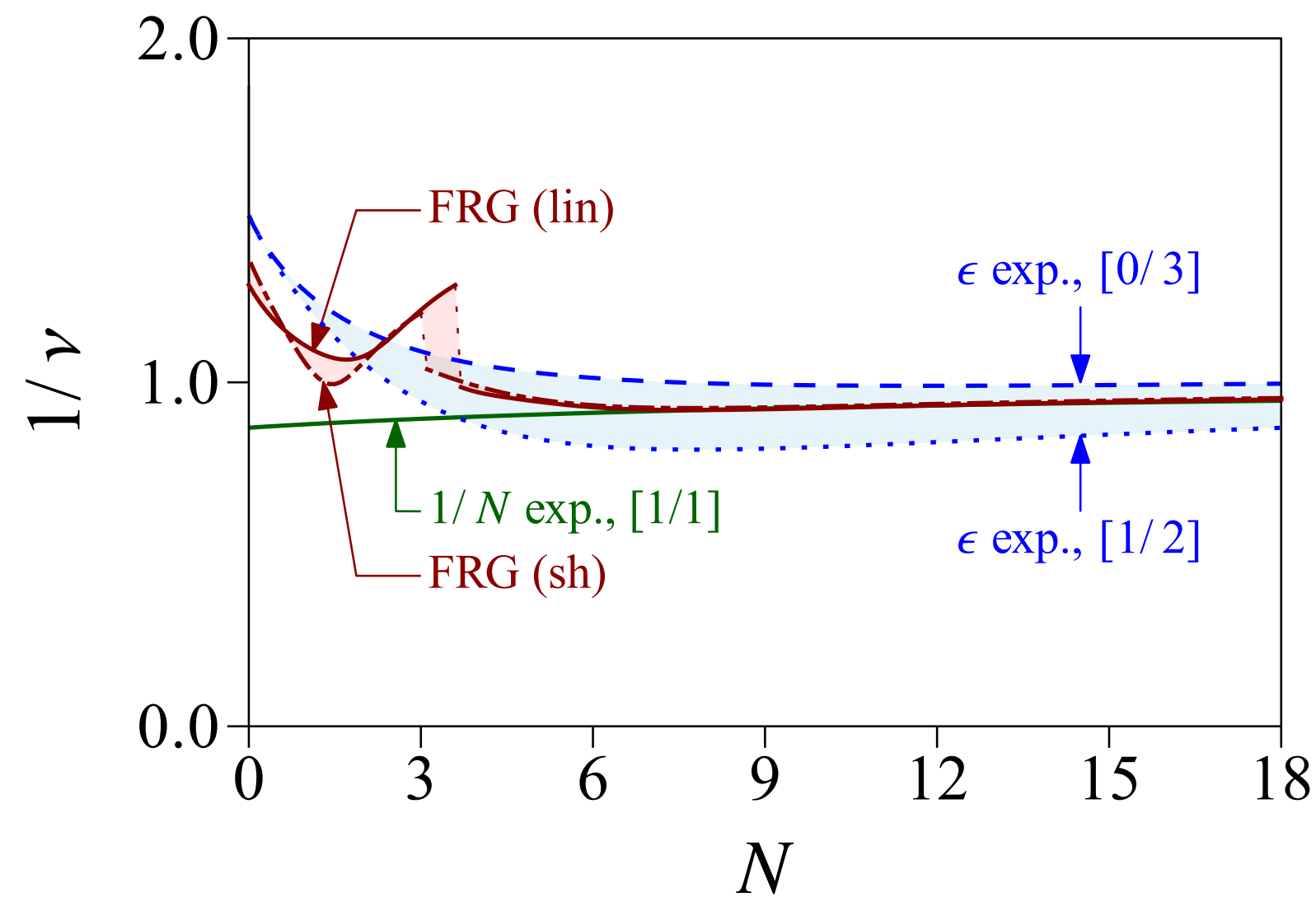
[Ray, Ihrig, Kruti, Gracey, Scherer, LJ, PRB '21]

Gross-Neveu-SO(3) criticality

Field theory:

$$\mathcal{S} = \int d^2\vec{x}d\tau \left[\bar{\psi}\gamma^\mu\partial_\mu\psi + g\vec{\varphi} \cdot \bar{\psi}(\mathbb{1}_2 \otimes \vec{L})\psi + \frac{1}{2}\vec{\varphi}(-\partial_\mu^2 + m^2)\vec{\varphi} + \lambda(\vec{\varphi} \cdot \vec{\varphi})^2 \right]$$

Critical exponents:



Spin-orbital realization ($N = 3$):

$$1/\nu = 1.03(15)$$

$$\eta_\phi = 0.42(7)$$

$$\eta_\psi = 0.180(10)$$

... from three-loop ϵ expansion,
second/third-order $1/N$ expansion,
functional RG in local potential approximation

... different from Gross-Neveu-Heisenberg

[Ray, Ihrig, Kruti, Gracey, Scherer, LJ, PRB '21]

Comparison of exponents: Gross-Neveu-Heisenberg

$N_f = 2$	Year	$1/\nu$	η_ϕ
Interpolation (this work)	2022	0.83(12)	1.01(6)
$4 - \varepsilon$ expansion, $\mathcal{O}(\varepsilon^4)$ [13]	2017	0.64	0.98
$1/N_f$ expansion, $\mathcal{O}(1/N_f^{2,3})$ [39]	2018	0.85	1.18
functional RG, NLO [38]	2018	0.80	1.03
functional RG, LPA' [16]	2014	0.77	1.01
DQMC, $\tau \sim L \leq 40$ [31]	2020	0.95(5)	0.75(4)
DQMC, $\tau \sim L \leq 40$ [30]	2016	0.98(1)	0.47(7)
DQMC, $\beta = L \leq 24$ [29]	2021	1.11(4)	0.80(9)
DQMC, $\beta = L \leq 21$ [28]	2019	1.14(9)	0.79(5)
DQMC, $\tau = 60, L \leq 18$ [27]	2015	1.19(6)	0.70(15)
DQMC, $\beta = L \leq 20$ [32]	2021	1.01(8)	0.55(2)
HMC, $\beta \leq 12, L \leq 102$ [36]	2021	0.84(4)	0.52(1)
HMC, $\beta \leq 12, L \leq 102$ [35]	2020	0.84(4)	0.85(13)
HMC, $\beta = 21, L \leq 24$ [34]	2019	1.08	0.62
HMC, $\beta = 21, L \leq 18$ [33]	2018	0.86	0.87(2)

