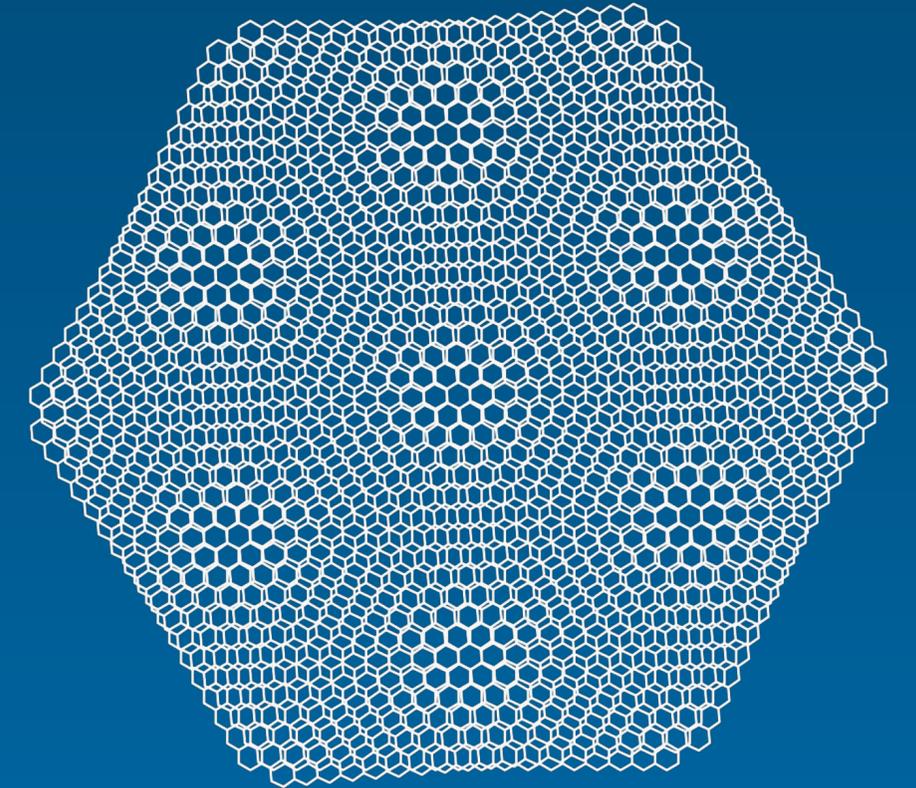


# Dirac quantum criticality in moiré systems

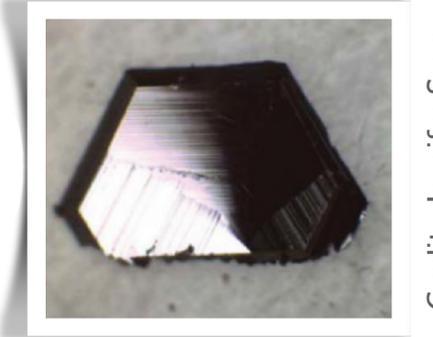
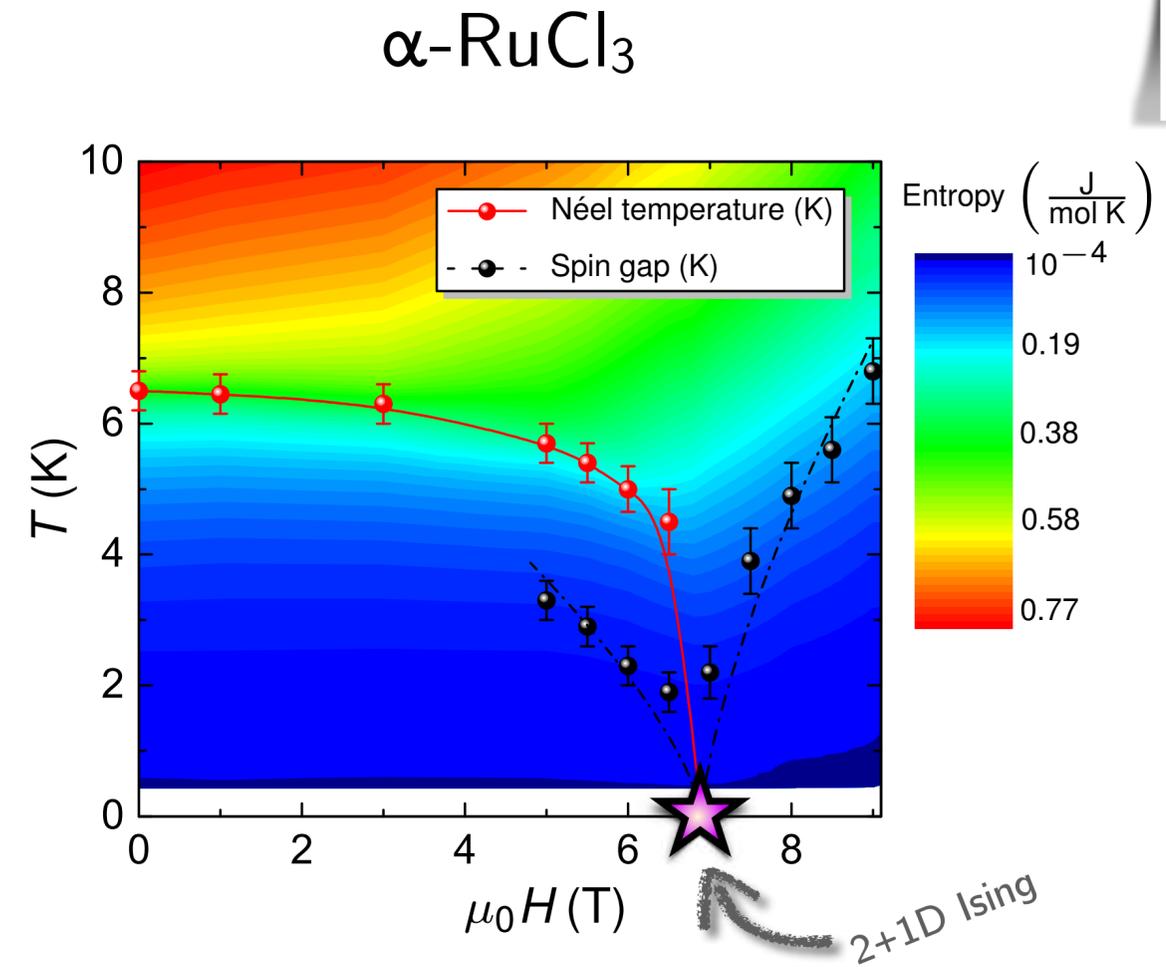
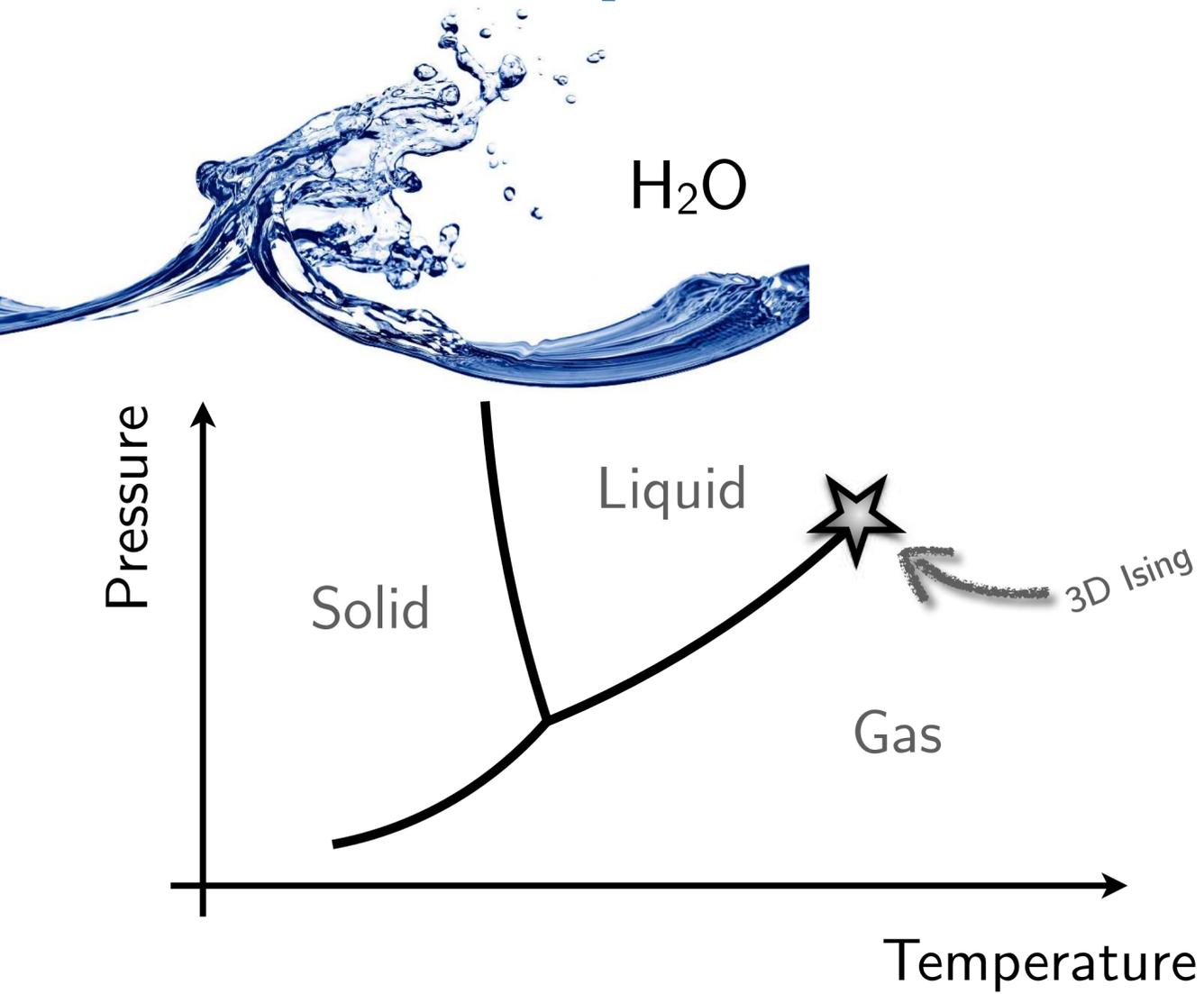
Lukas Janssen



Jan Biedermann



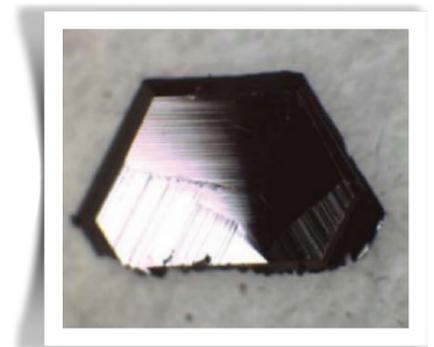
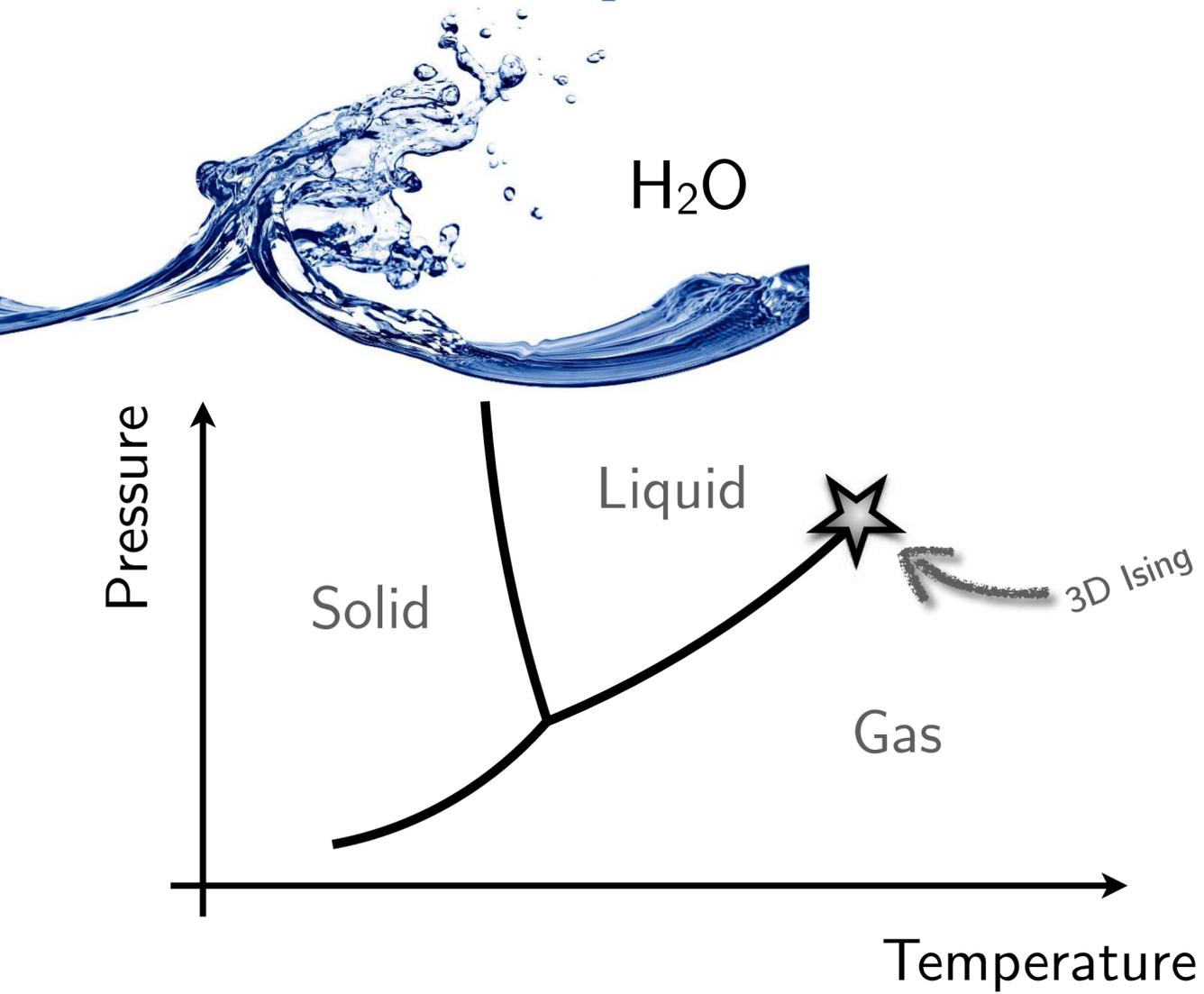
# Classical vs quantum criticality



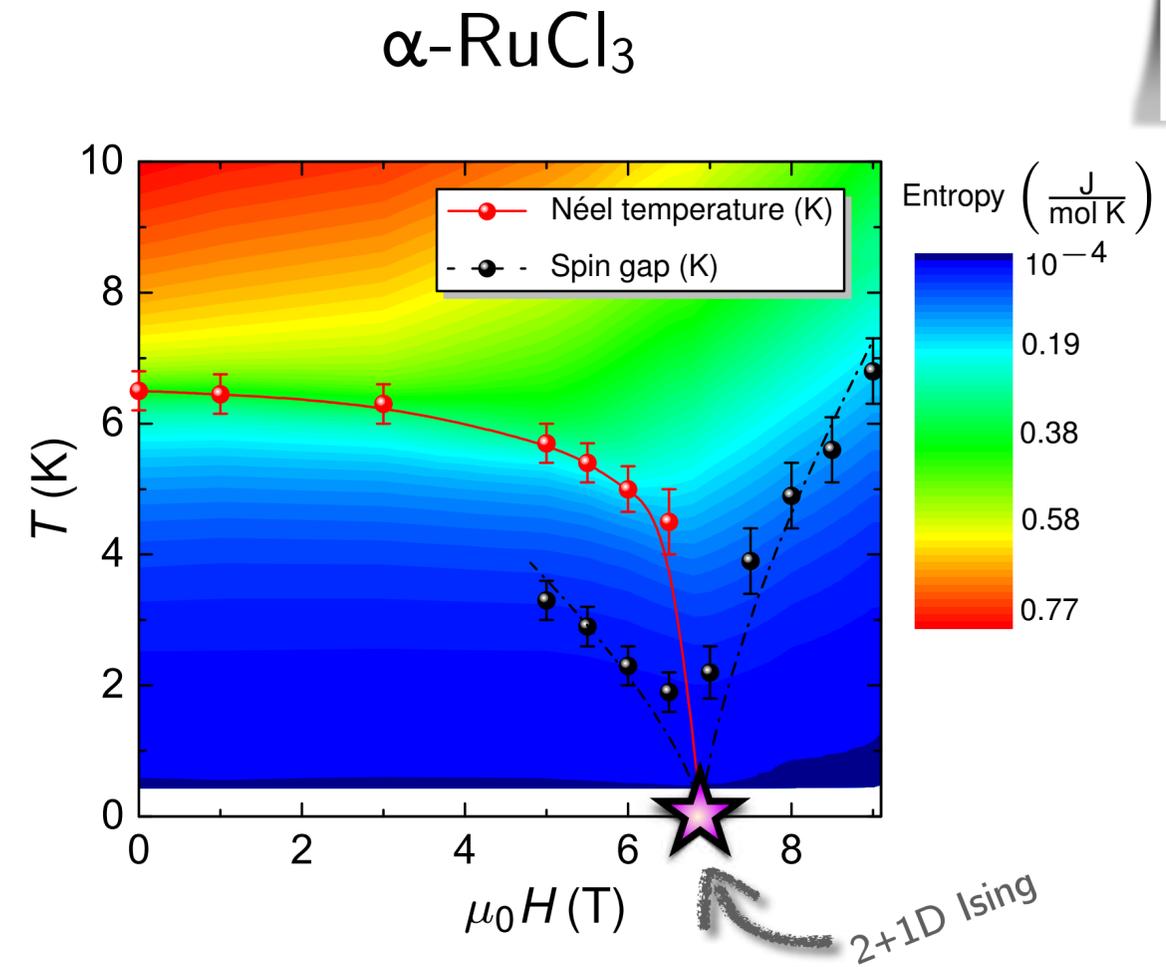
Credit: Jennifer Sears

[Wolter, Corredor, LJ, et al., PRB '17]

# Classical vs quantum criticality



Credit: Jennifer Sears



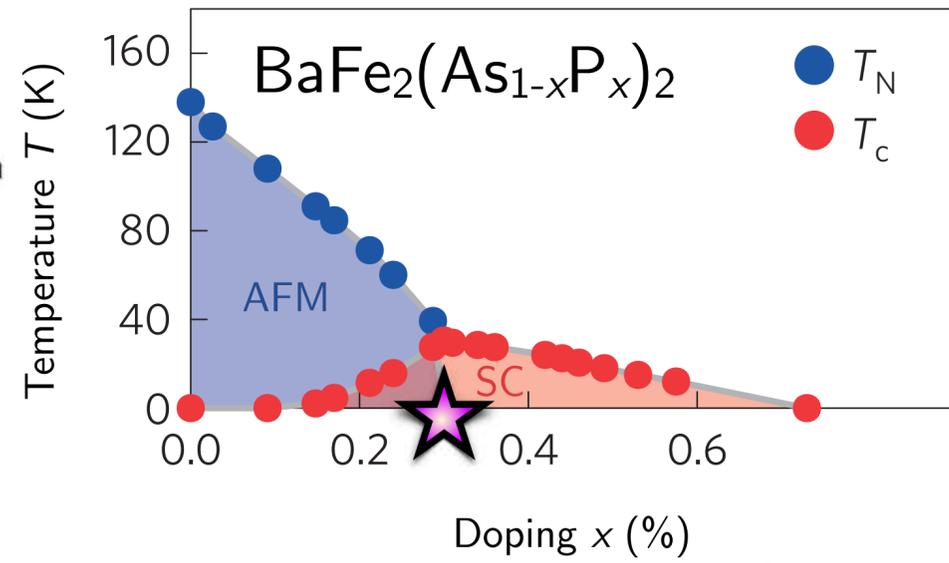
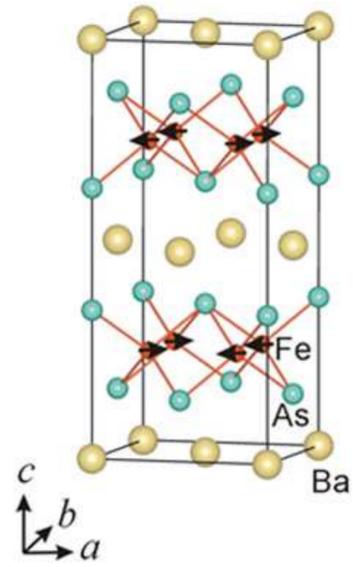
[Wolter, Corredor, LJ, et al., PRB '17]

Universal field theory:

$$\mathcal{L} = \frac{1}{2} \varphi (-\partial_\tau^2 - v_B^2 \vec{\nabla}^2) \varphi + \lambda \varphi^4 + \dots$$

# Genuine *quantum* criticality?

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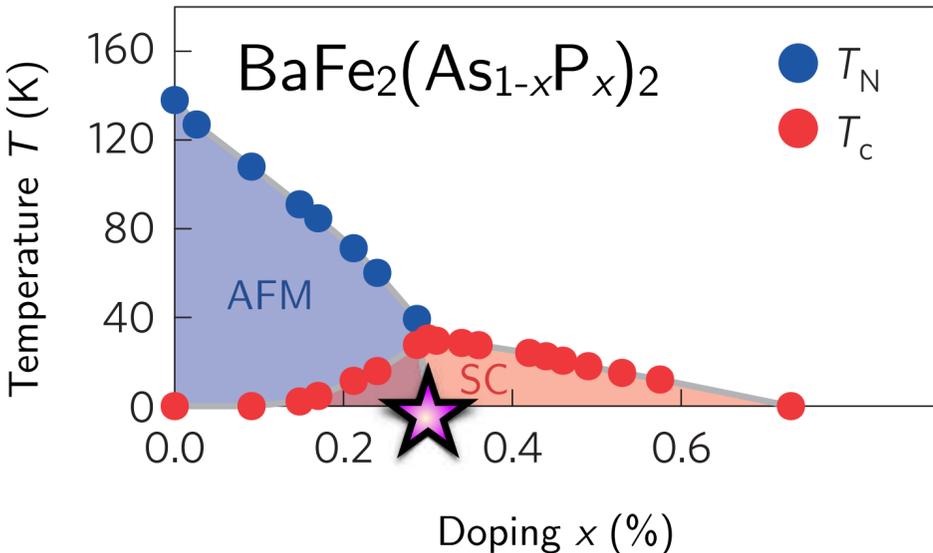
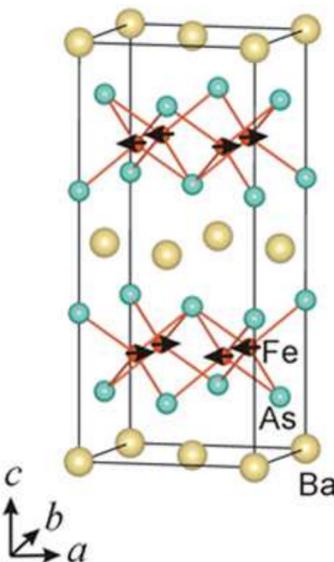


[Huang *et al.*, PRL '08]

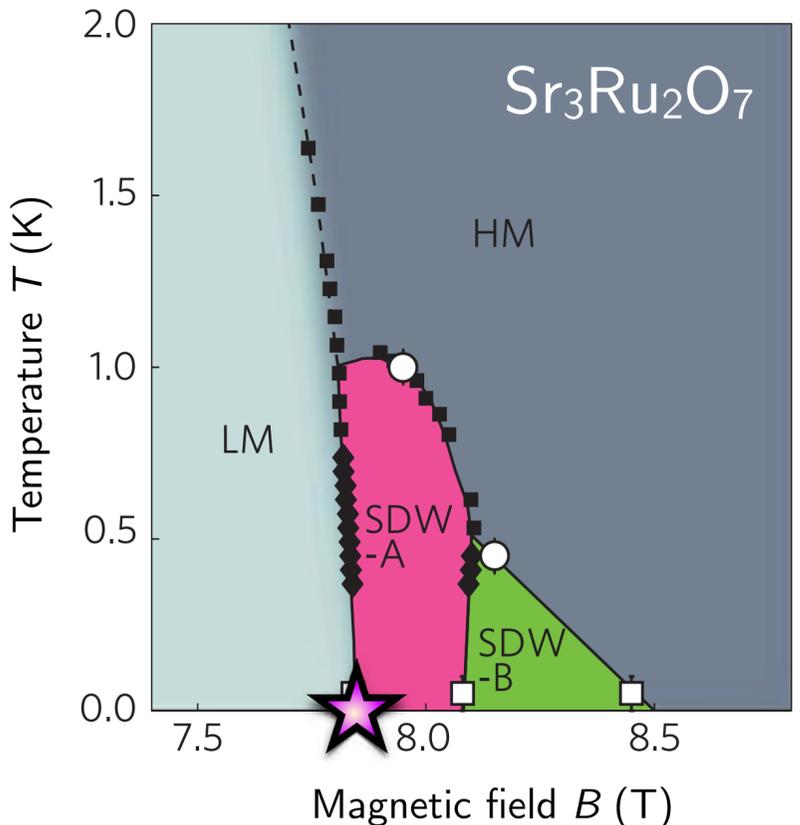
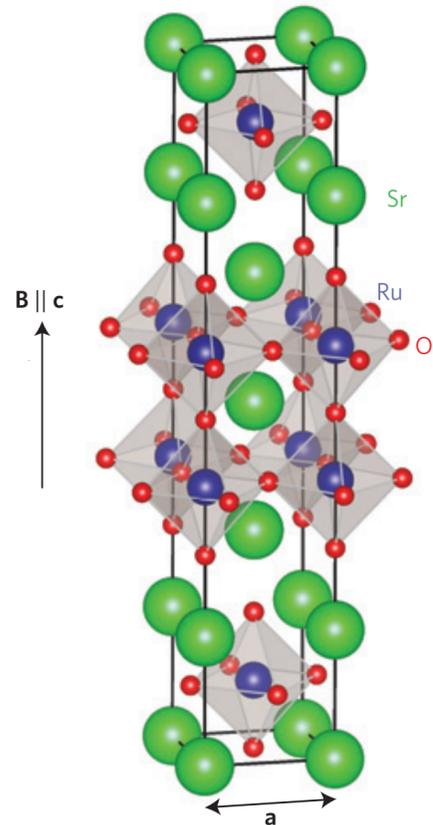
[Hashimoto *et al.*, Science '12]

[Hayes *et al.*, Nat. Phys. '16]

# Genuine *quantum* criticality?

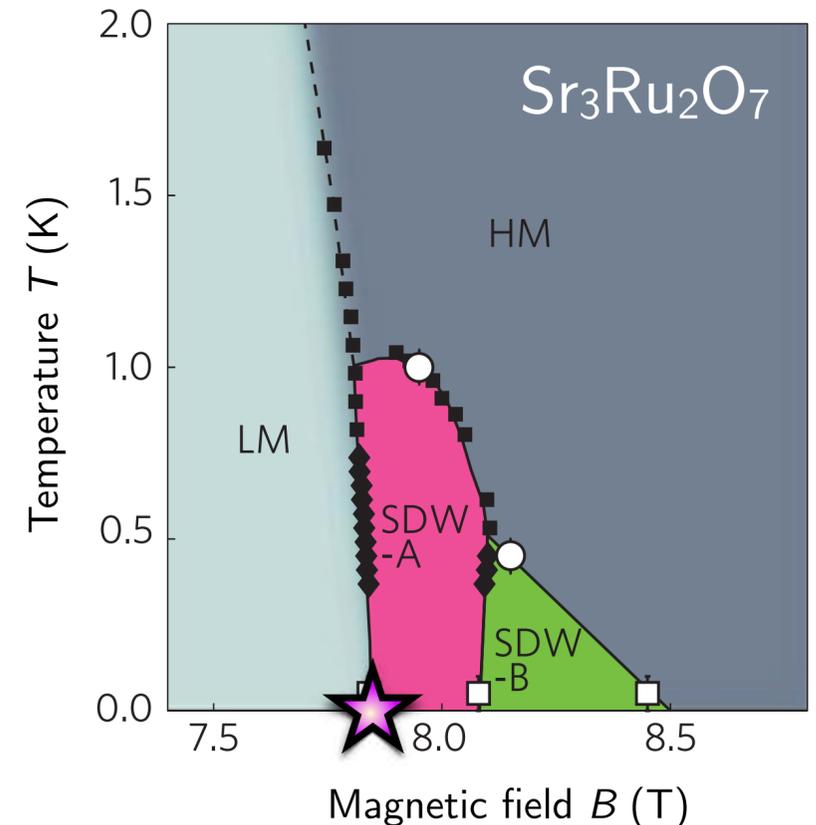
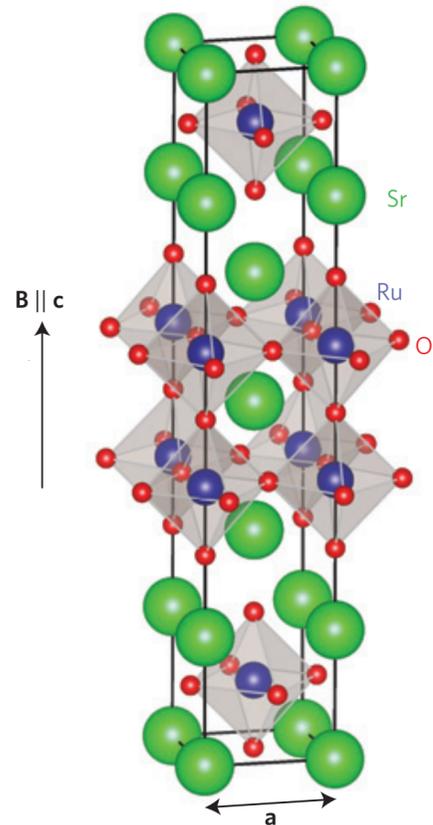
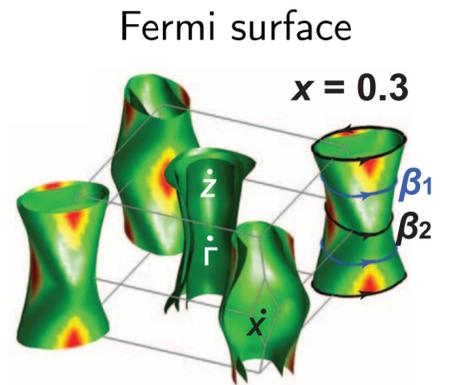
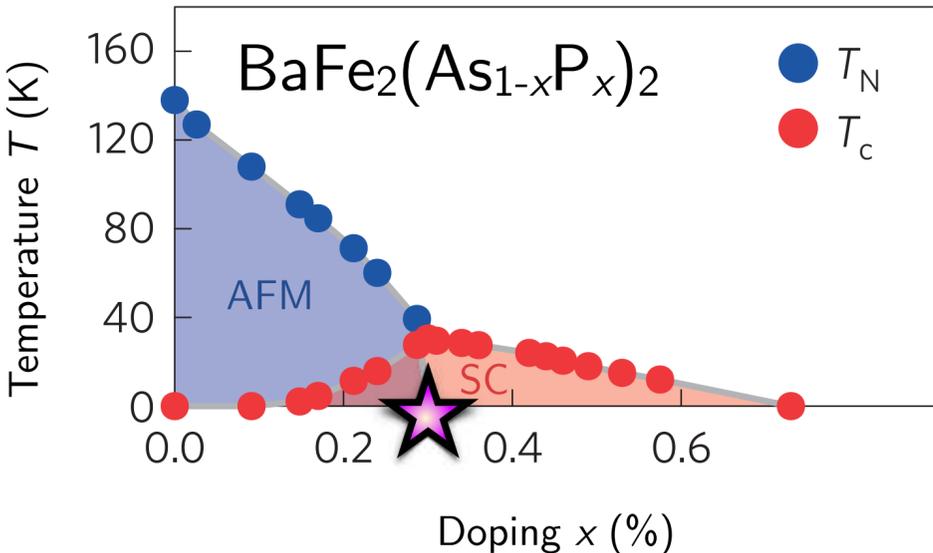
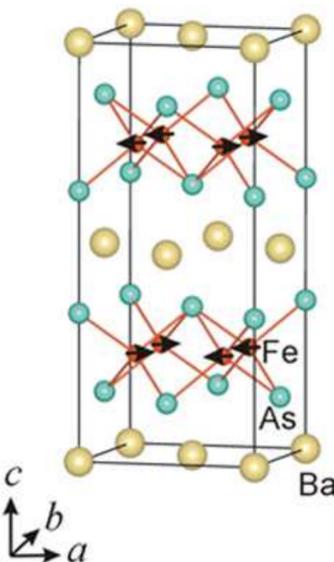


[Huang *et al.*, PRL '08]  
 [Hashimoto *et al.*, Science '12]  
 [Hayes *et al.*, Nat. Phys. '16]



[Lester *et al.*, Nat. Mat. '15]

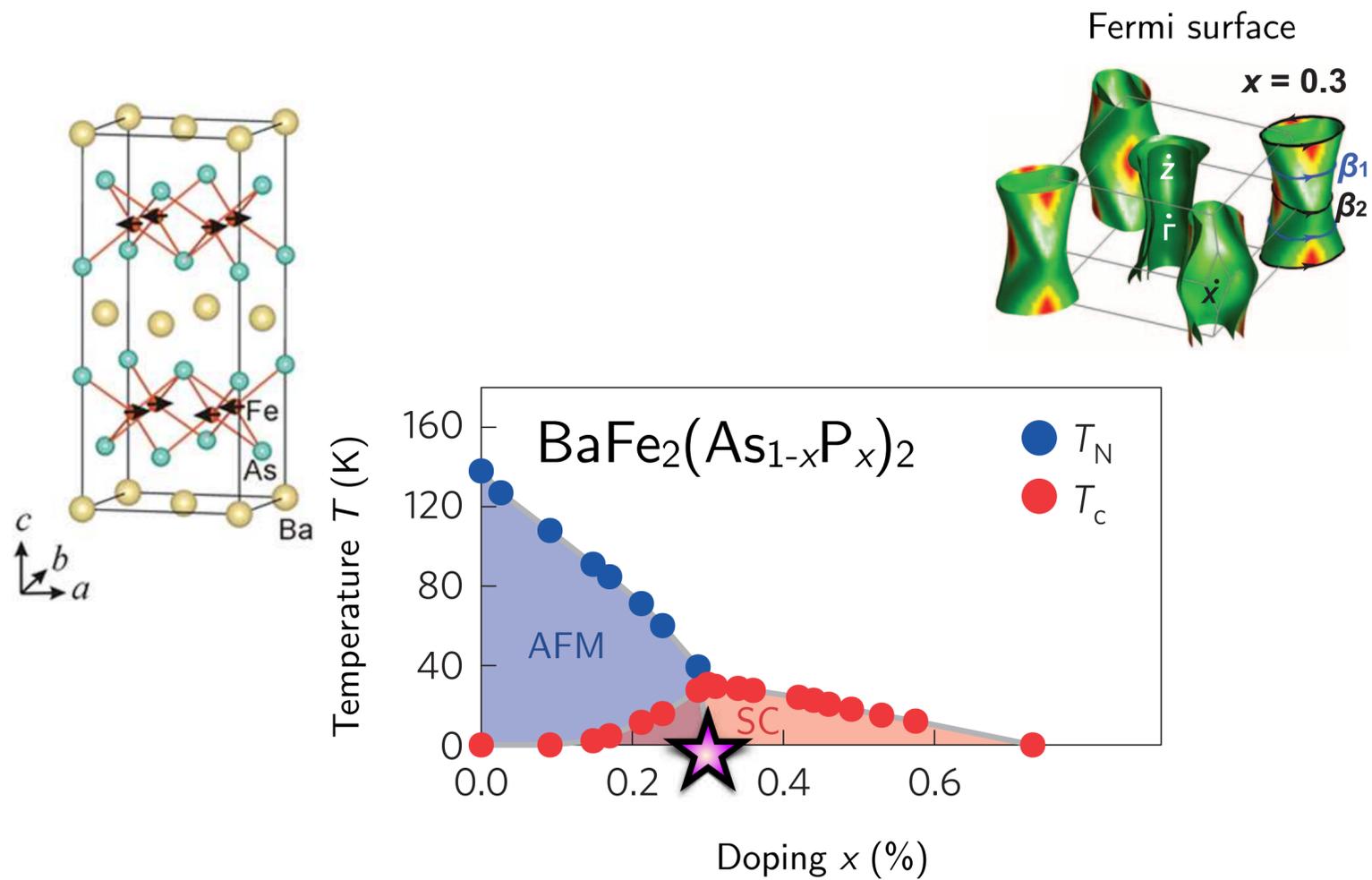
# Genuine *quantum* criticality?



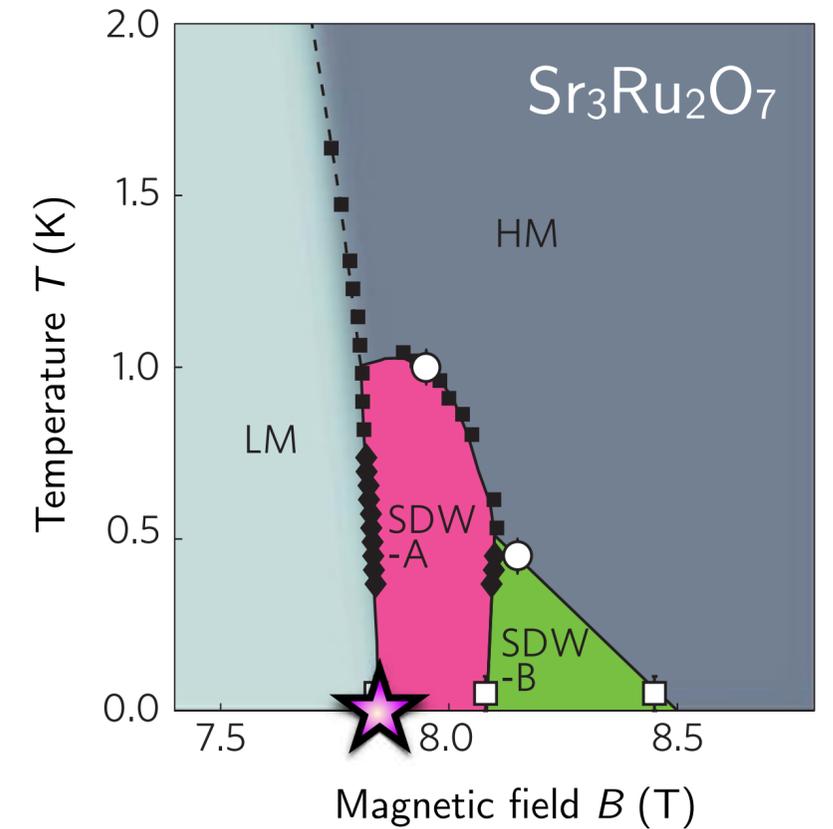
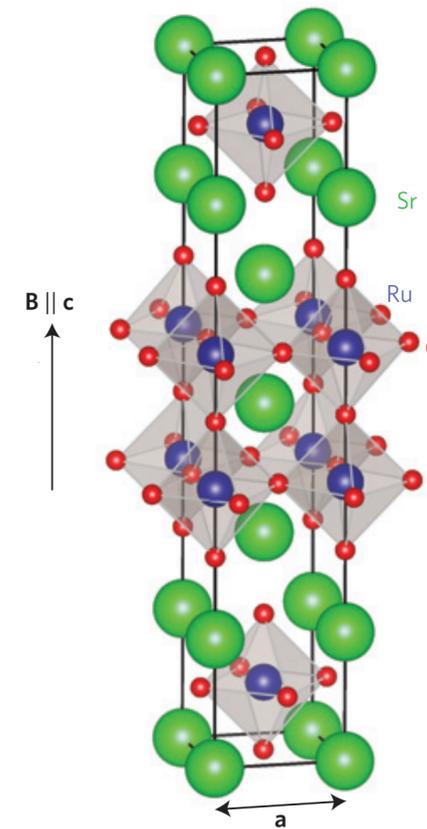
[Huang *et al.*, PRL '08]  
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# Genuine *quantum* criticality?



[Huang *et al.*, PRL '08]  
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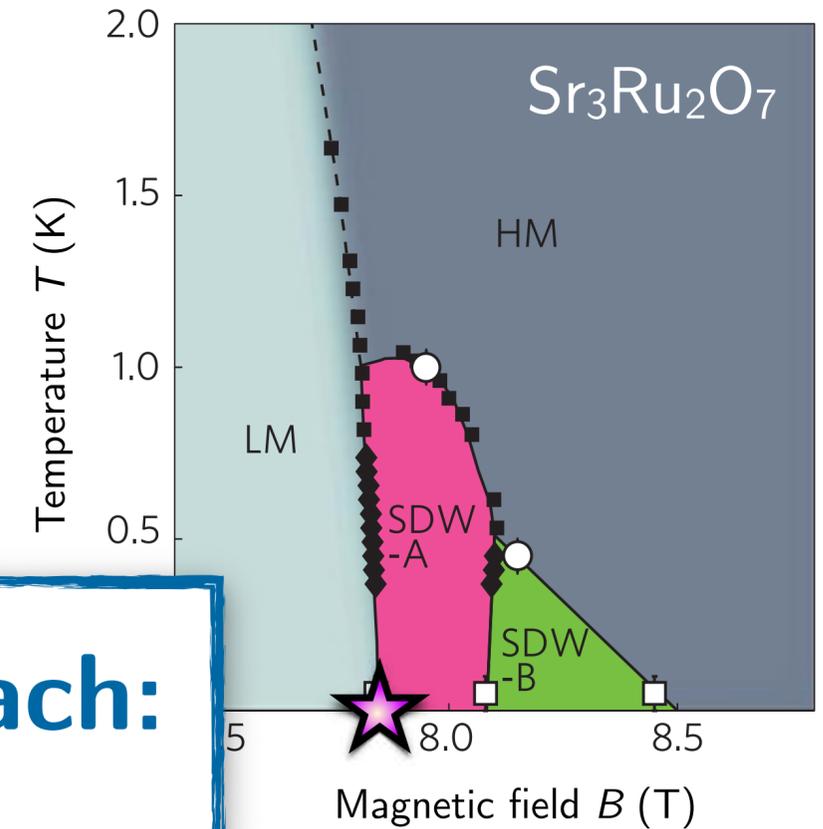
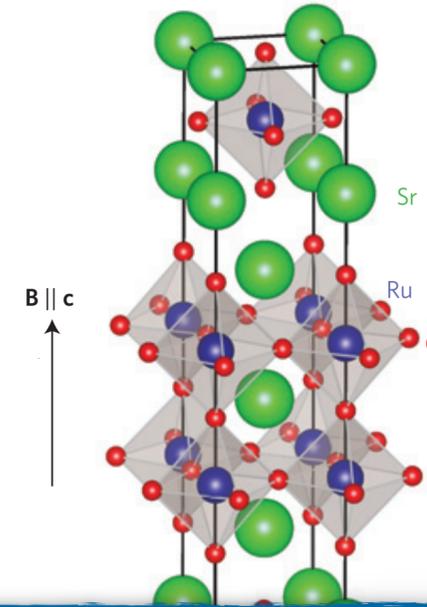
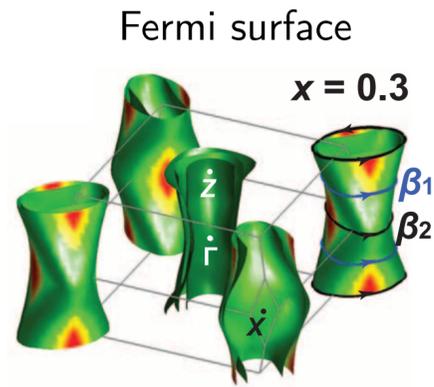
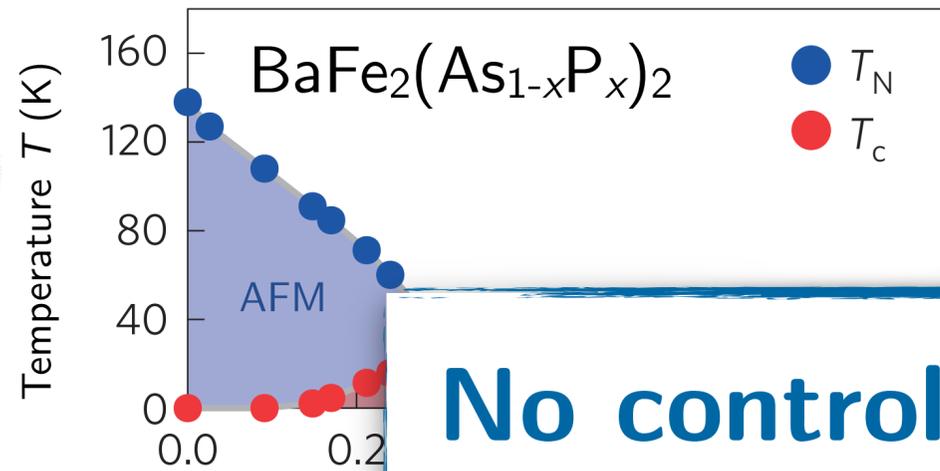
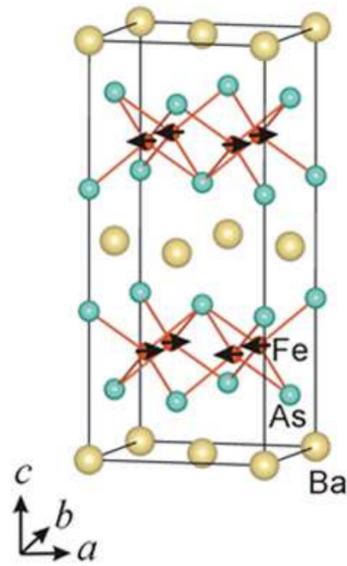
[Lester *et al.*, Nat. Mat. '15]

Universal field theory:

$$\mathcal{L} = \frac{1}{2} \vec{\varphi} \cdot (-\partial_\tau^2 - v_B^2 \vec{\nabla}^2) \vec{\varphi} + \psi^\dagger \left( \partial_\tau + \varepsilon(-i\vec{\nabla}) \right) \psi + \vec{\varphi} \cdot \psi^\dagger \vec{\sigma} \psi + \dots$$

[Metlitski & Sachdev, PRB '10]

# Genuine *quantum* criticality?



**No controlled theoretical approach:**

- ⚡ Too many gapless modes
- ⚡ Too small symmetry

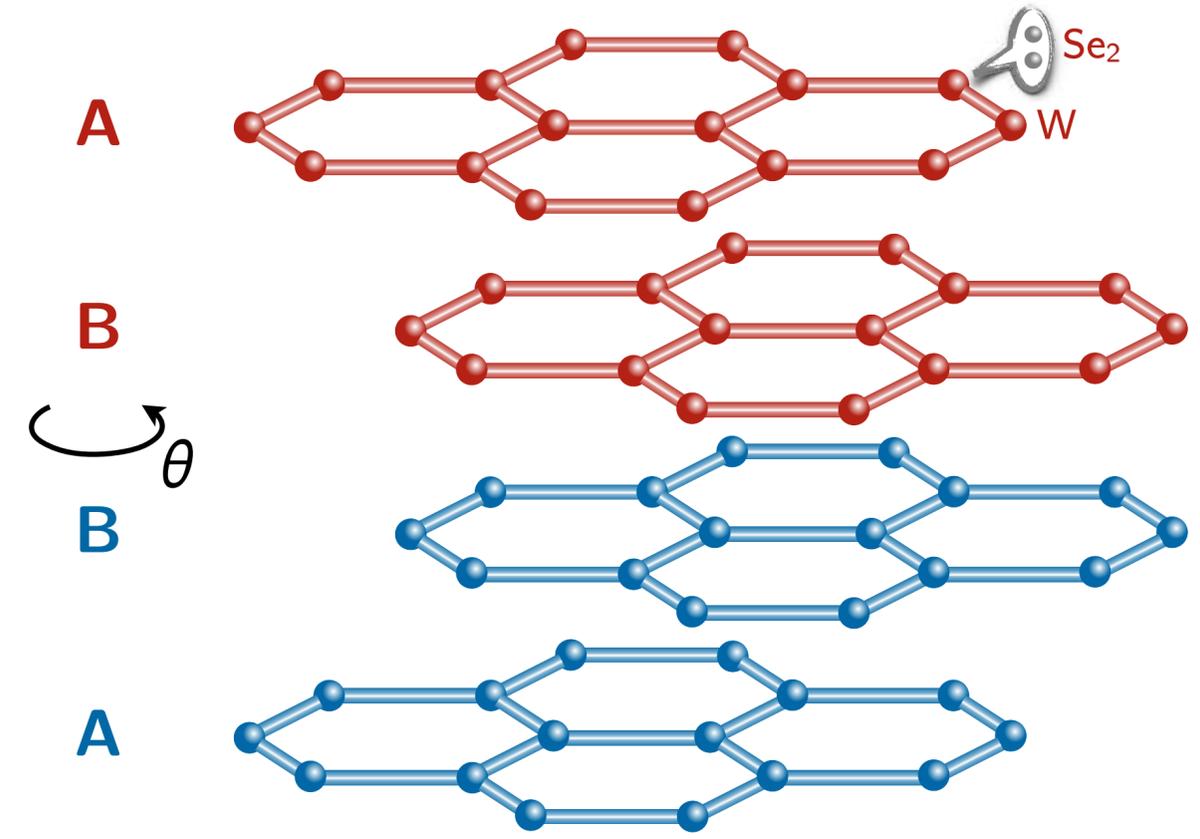
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[Lester *et al.*, Nat. Mat. '15]

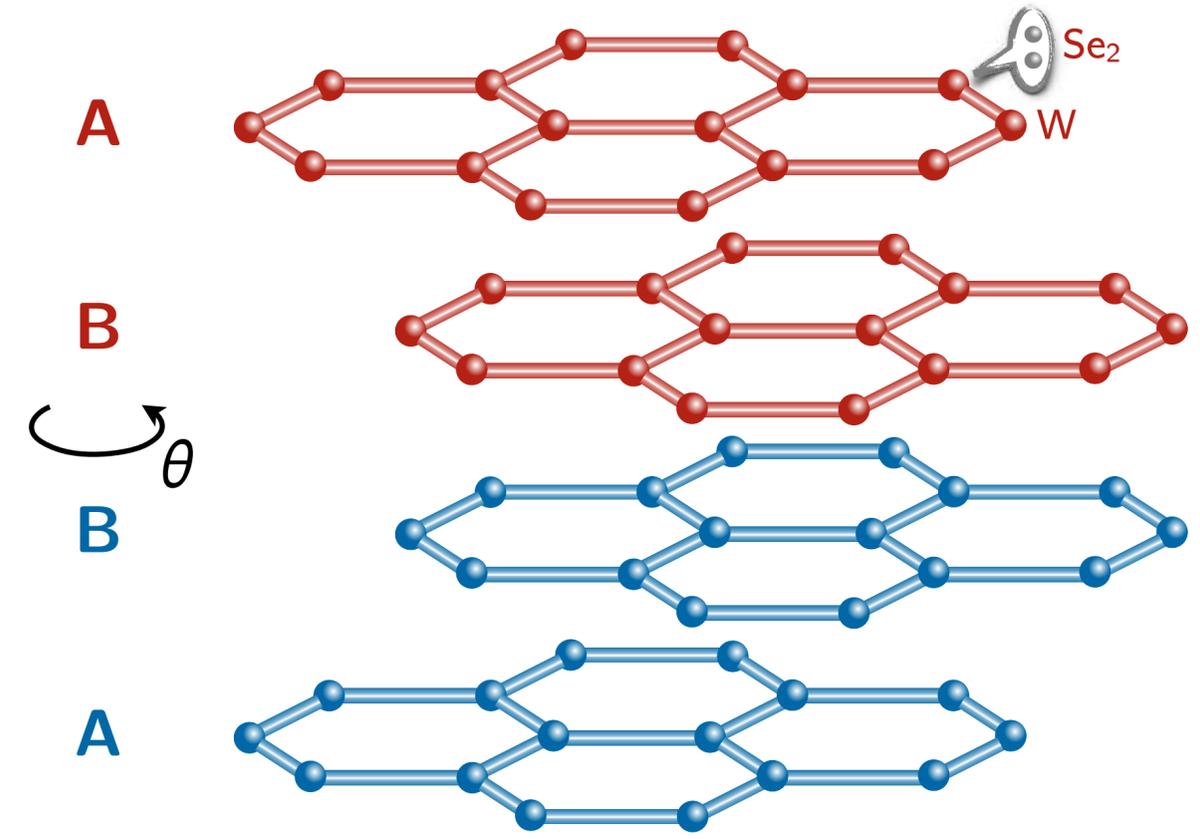
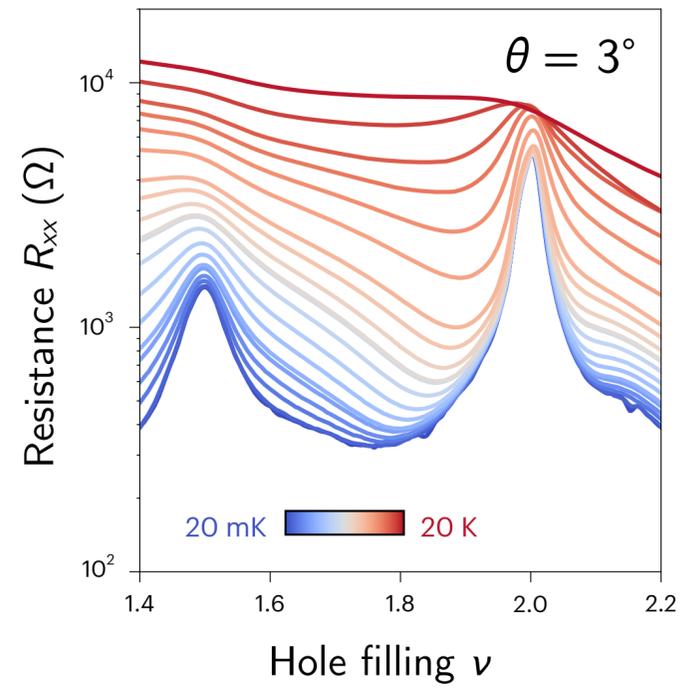
[Metlitski & Sachdev, PRB '10]

# Twisted double bilayer WSe<sub>2</sub>: Experiments



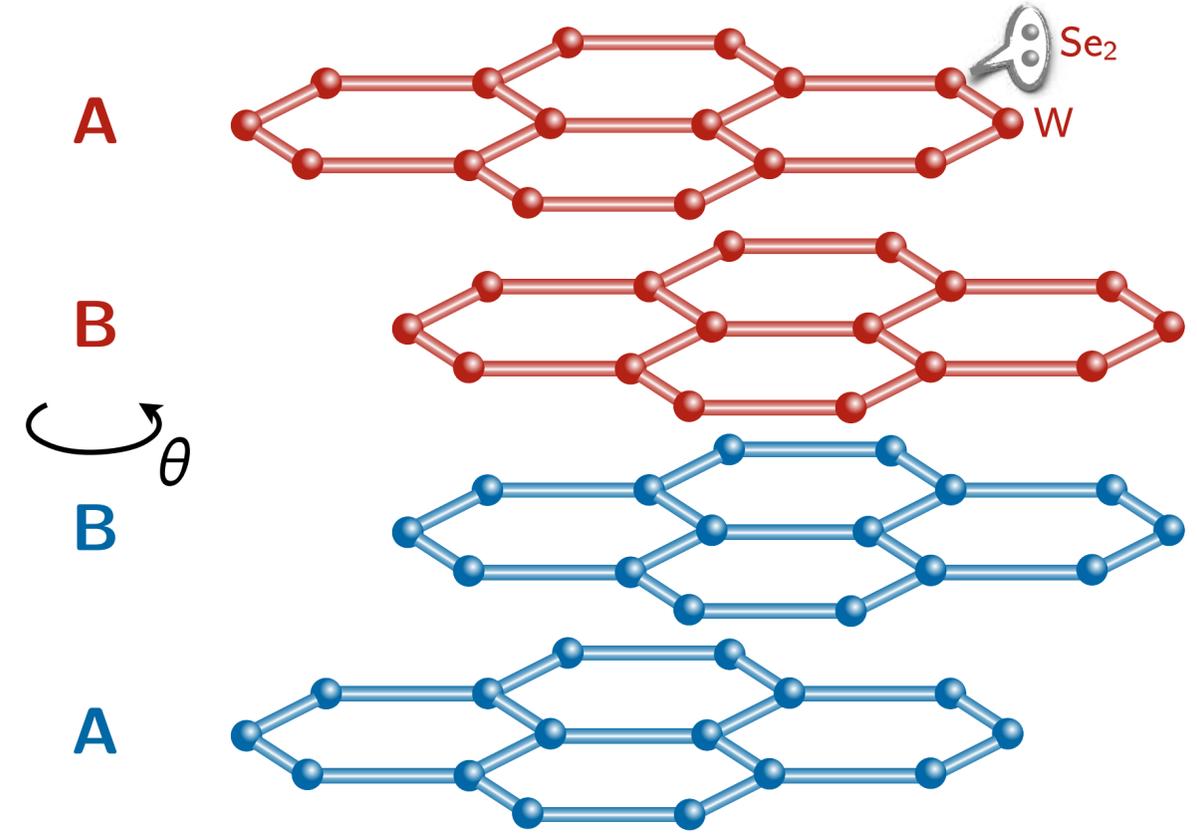
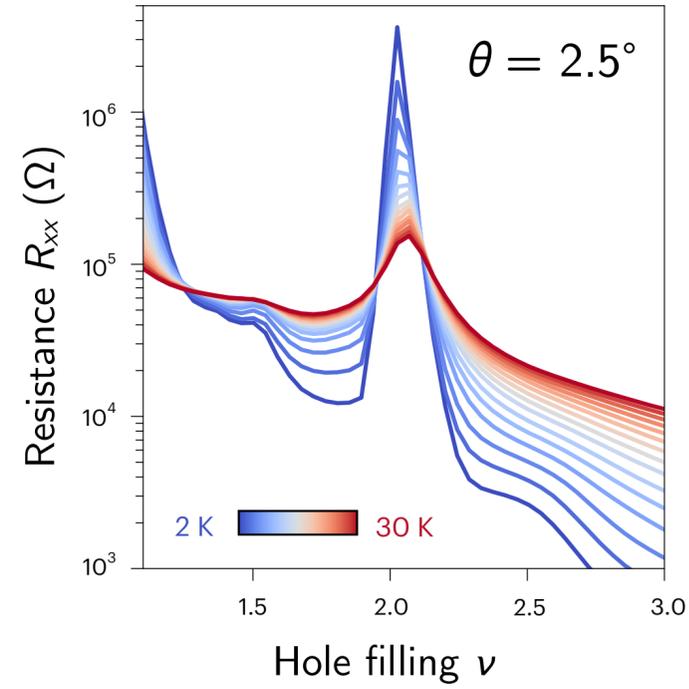
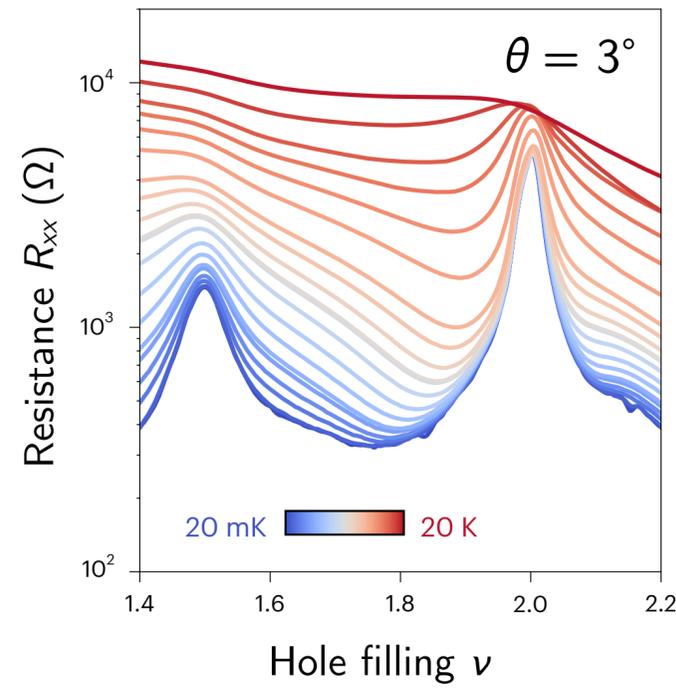
# Twisted double bilayer $\text{WSe}_2$ : Experiments

Transport:



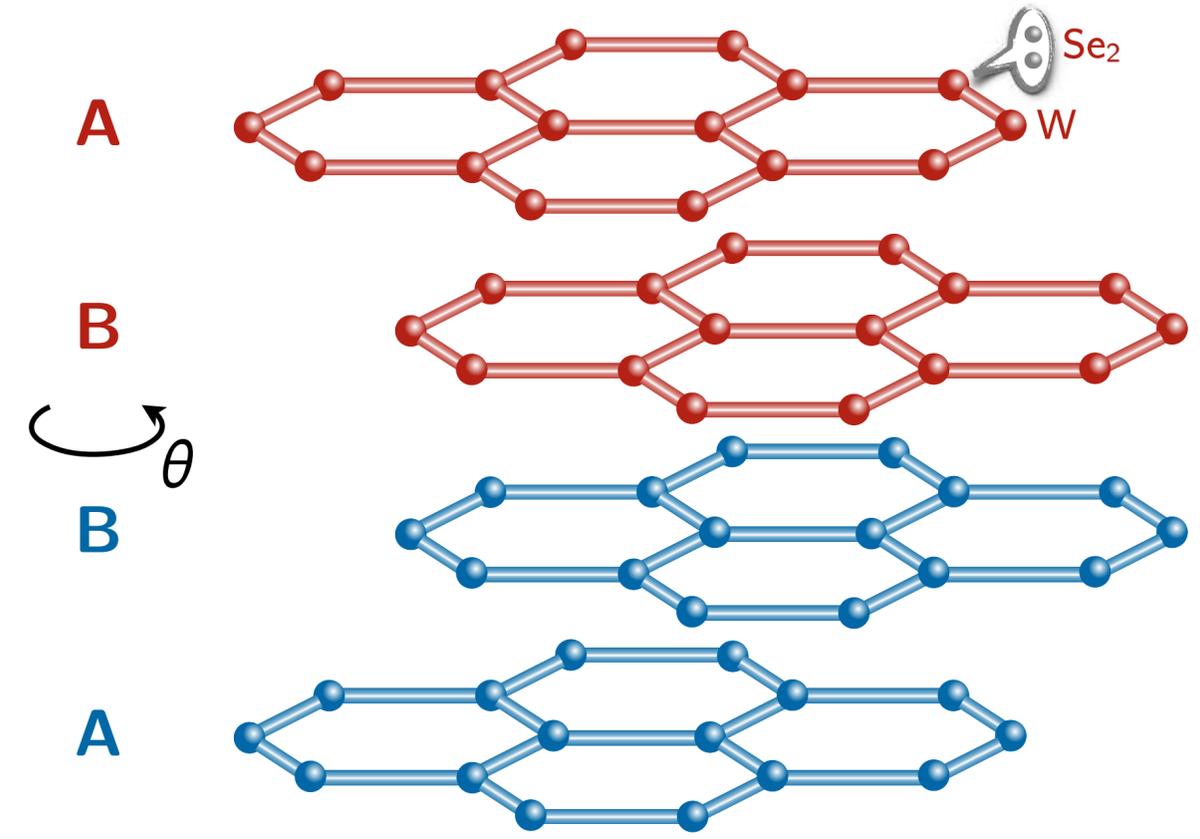
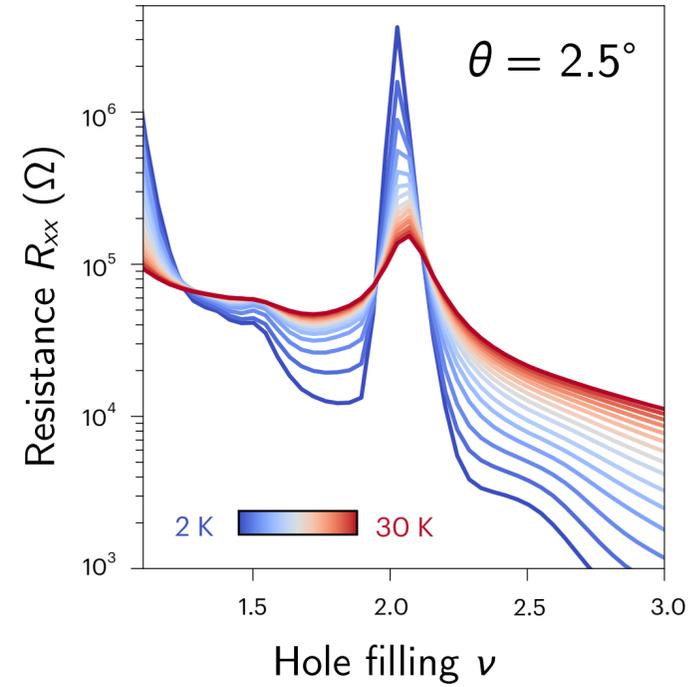
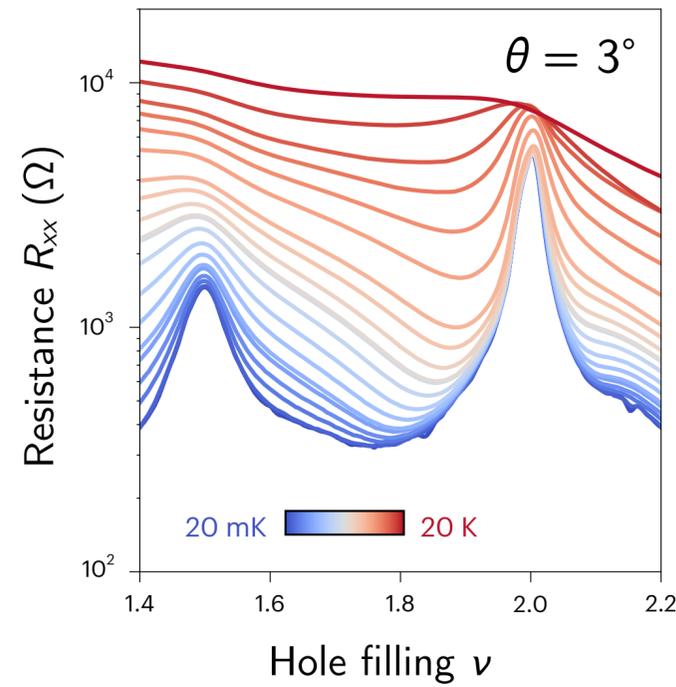
# Twisted double bilayer $WSe_2$ : Experiments

Transport:

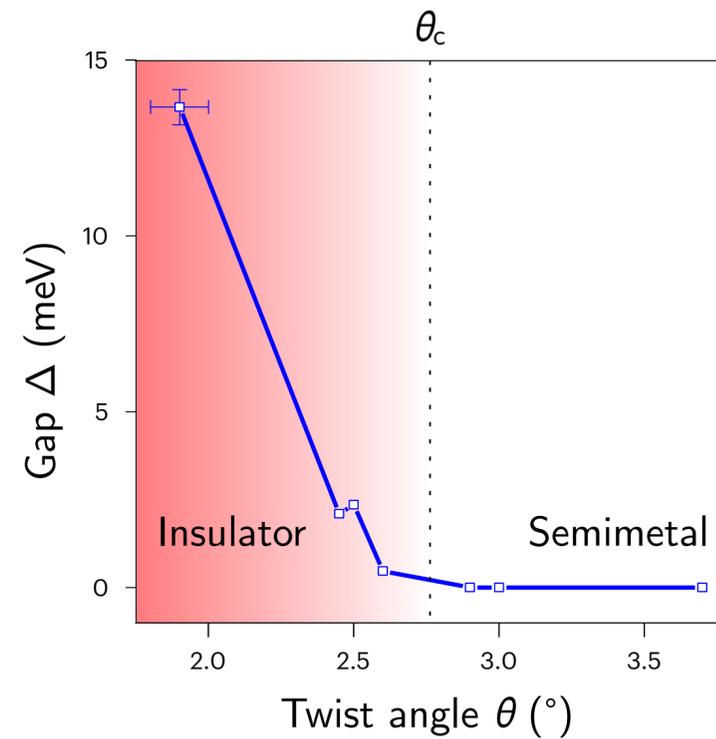


# Twisted double bilayer $WSe_2$ : Experiments

Transport:



Semimetal-to-insulator transition:

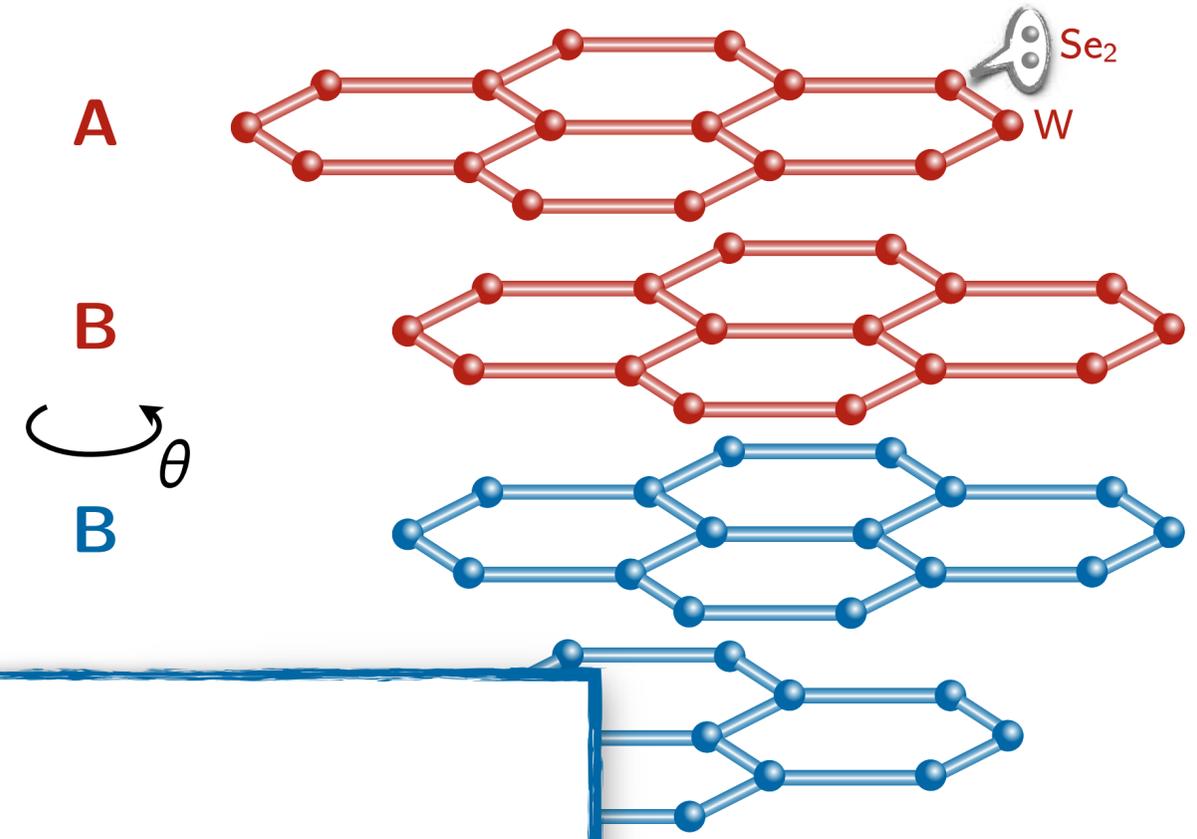
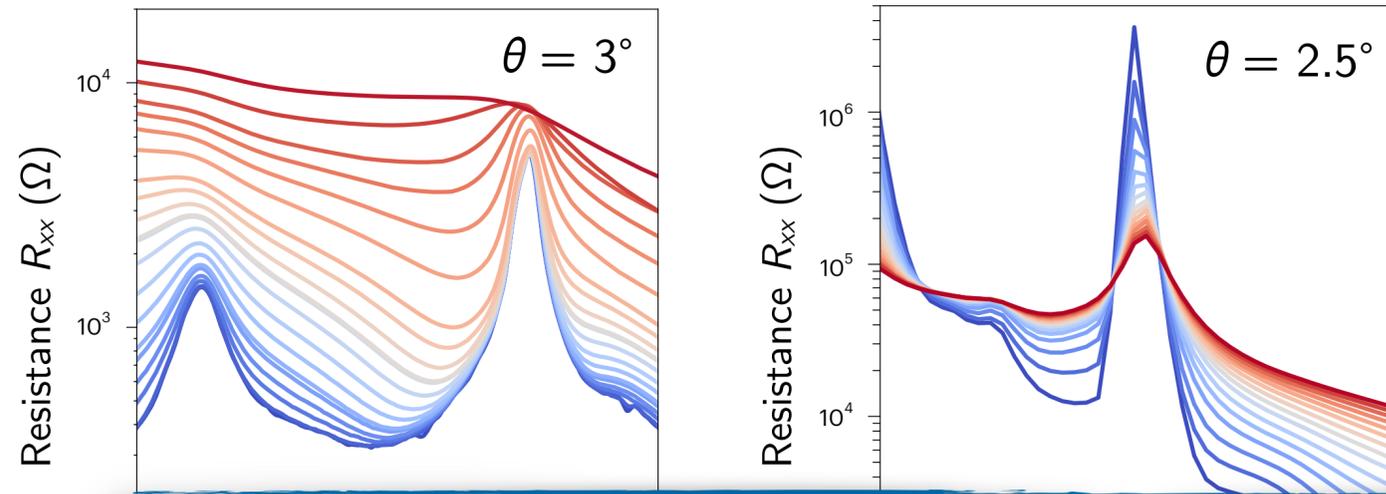


... for  $\nu = 2$

[Ma *et al.*, Nat. Mater. '25]

# Twisted double bilayer $WSe_2$ : Experiments

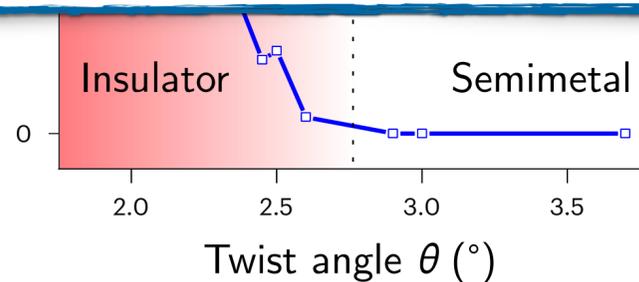
Transport:



Semimetal-to-insulator

**Today:**

- (1) Dirac semimetal ( $\theta > 2.7^\circ$ )
- (2) Antiferromagnetic insulator ( $\theta < 2.7^\circ$ )
- (3) Emergent Lorentz symmetry ( $\theta \simeq 2.7^\circ$ )

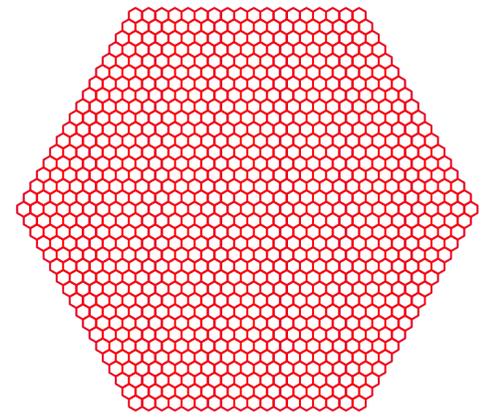
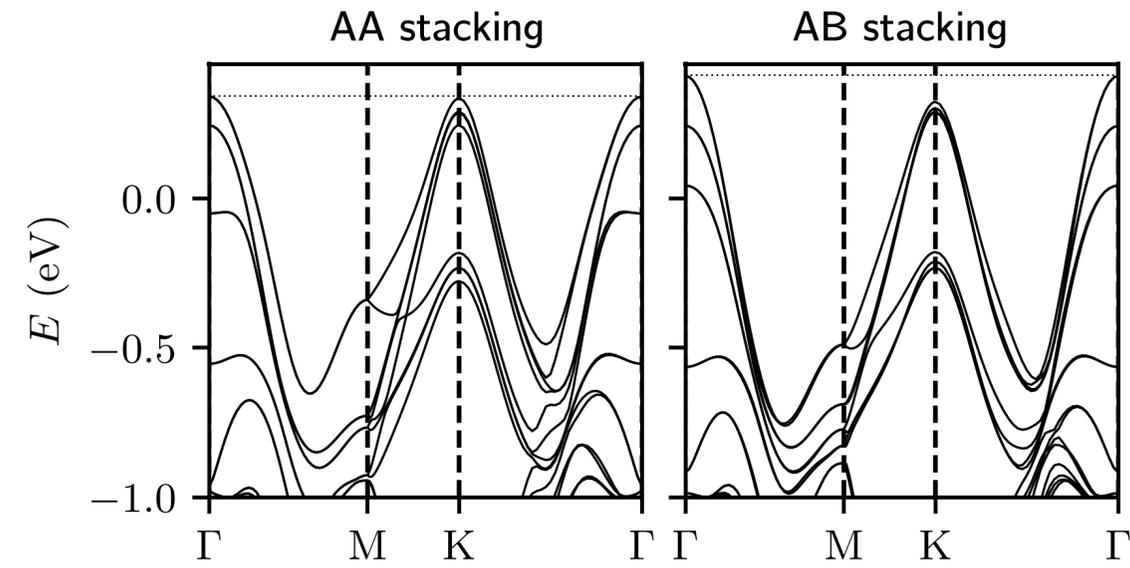


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# Twisted double bilayer $WSe_2$ : Noninteracting spectrum

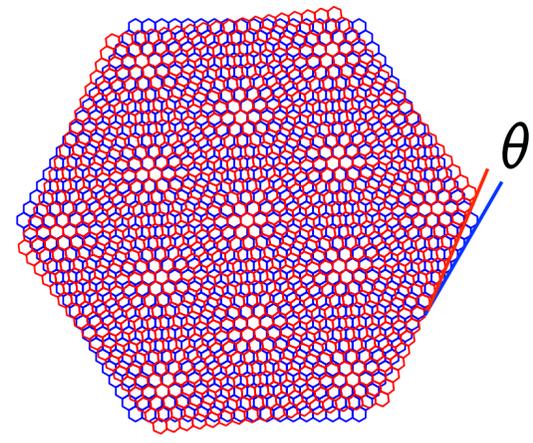
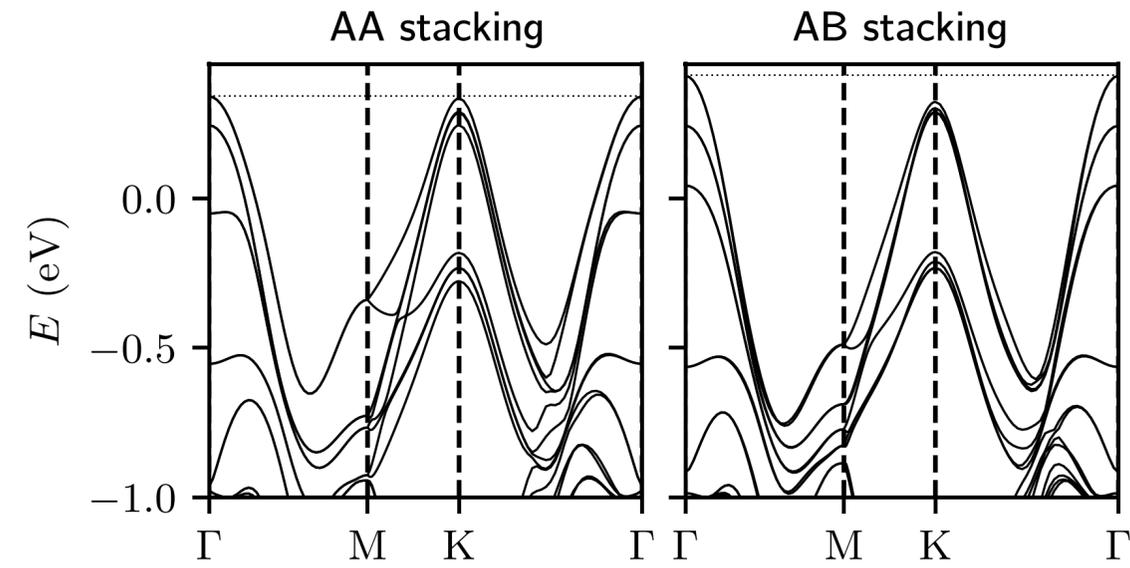
Untwisted spectrum:



[Pan *et al.*, PRRes '23]

# Twisted double bilayer $WSe_2$ : Noninteracting spectrum

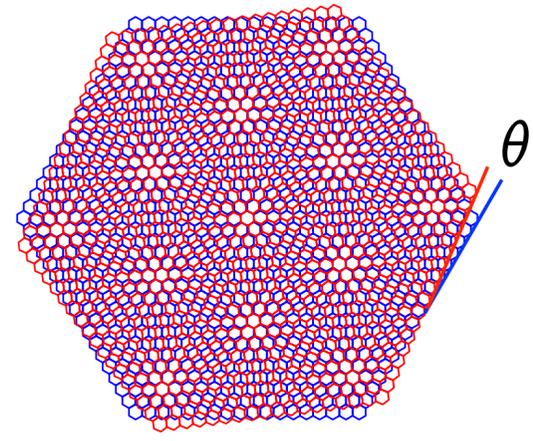
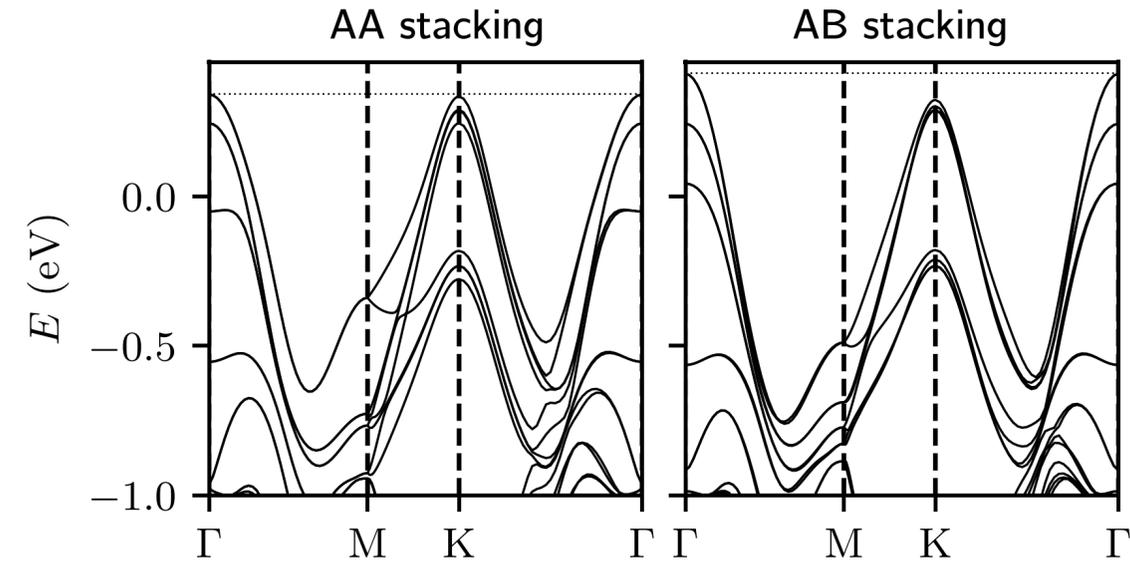
Untwisted spectrum:



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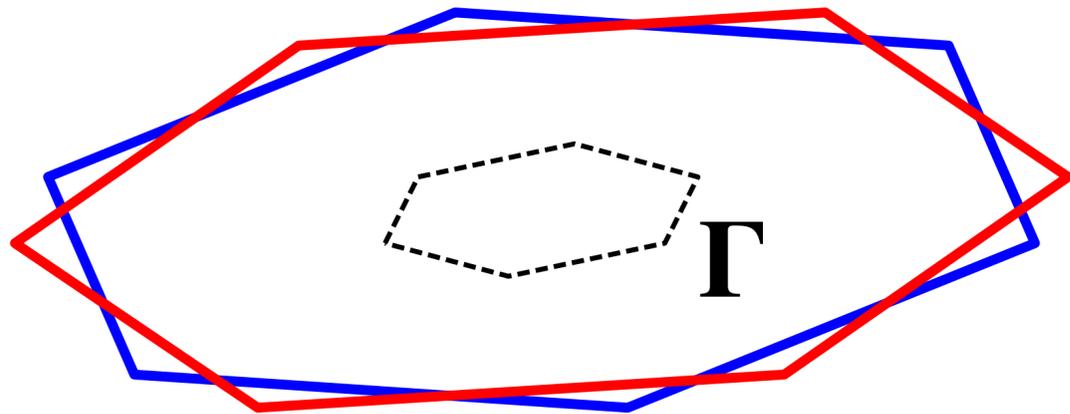
# Twisted double bilayer $WSe_2$ : Noninteracting spectrum

Untwisted spectrum:

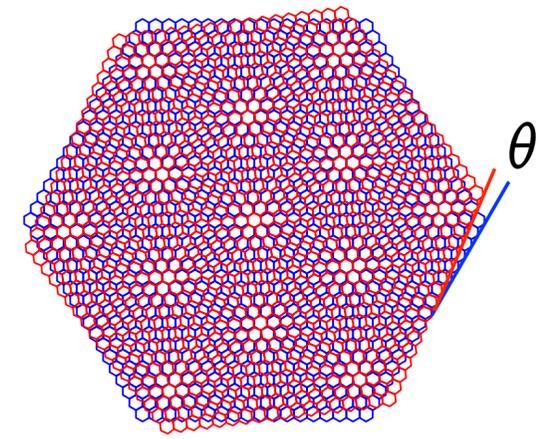


[Pan *et al.*, PRRes '23]

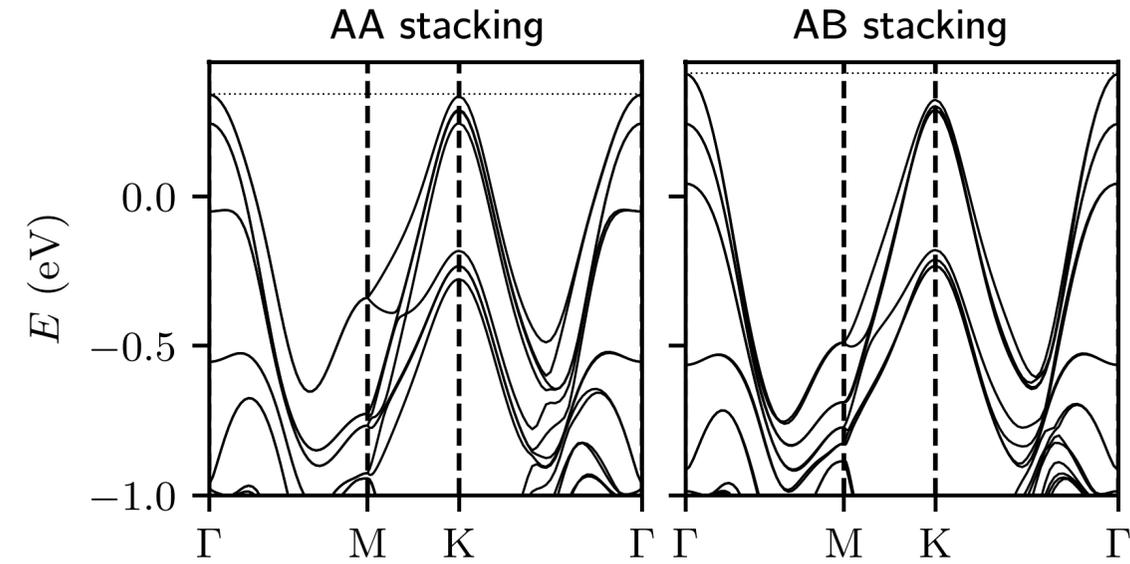
Moiré Brillouin zone:



# Twisted double bilayer WSe<sub>2</sub>: Noninteracting spectrum

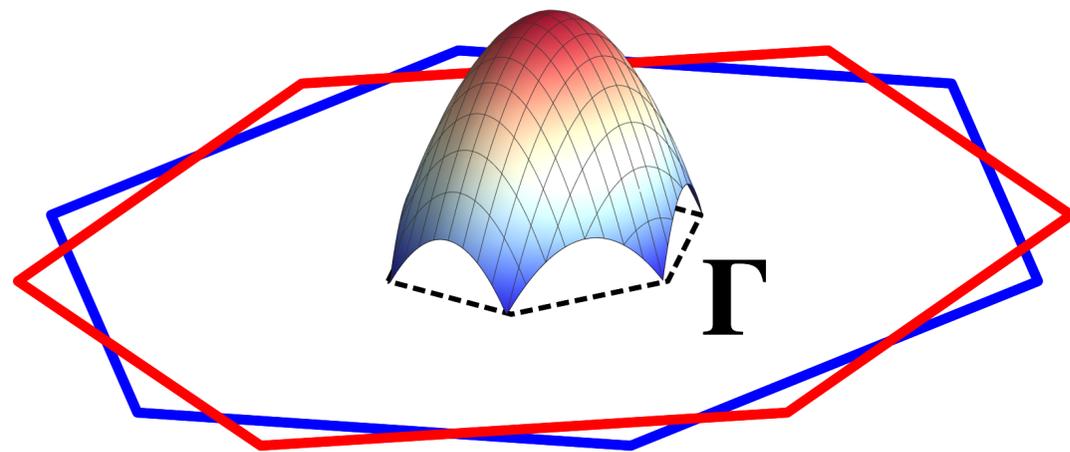


Untwisted spectrum:

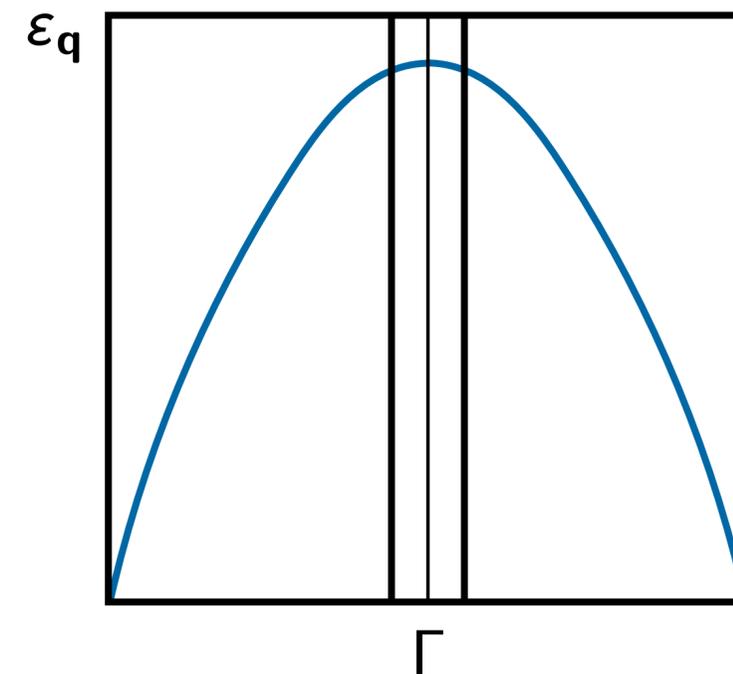


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Moiré Brillouin zone:

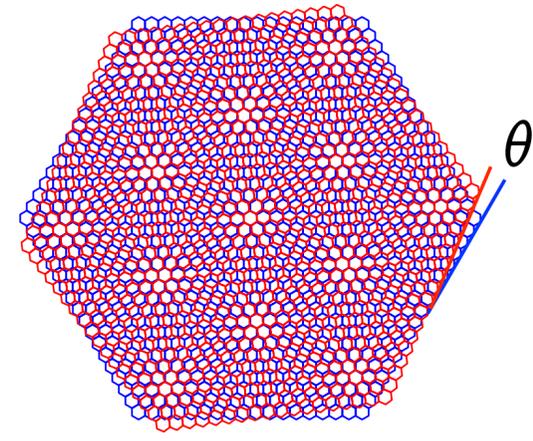


Twisted spectrum:

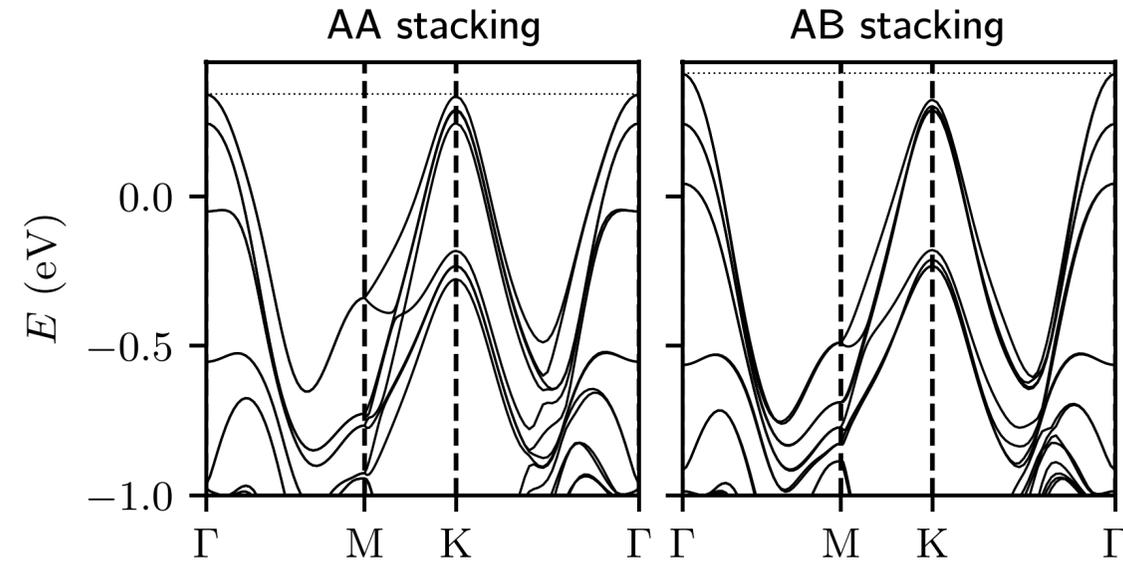


$$\epsilon_{\mathbf{q}} = -\frac{\mathbf{q}^2}{2m^*} + \mathcal{O}(q^3)$$

# Twisted double bilayer WSe<sub>2</sub>: Noninteracting spectrum

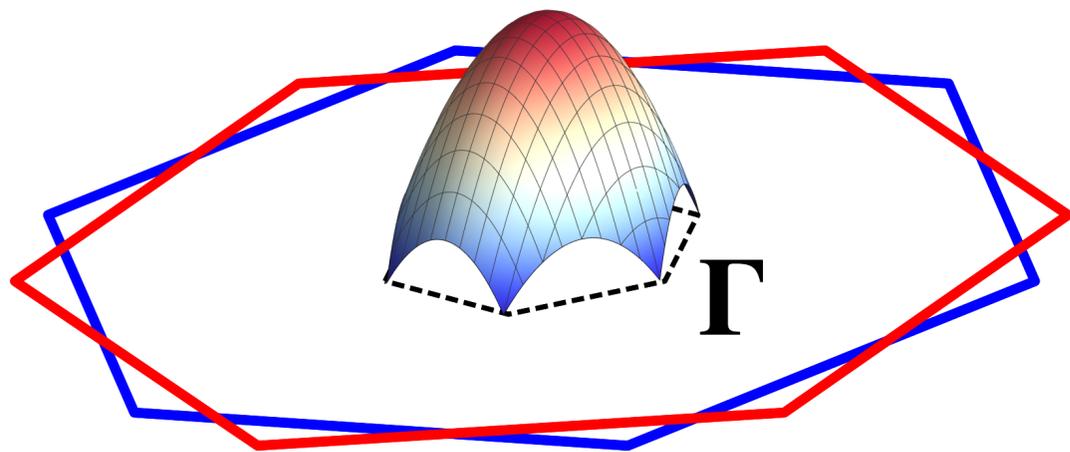


Untwisted spectrum:

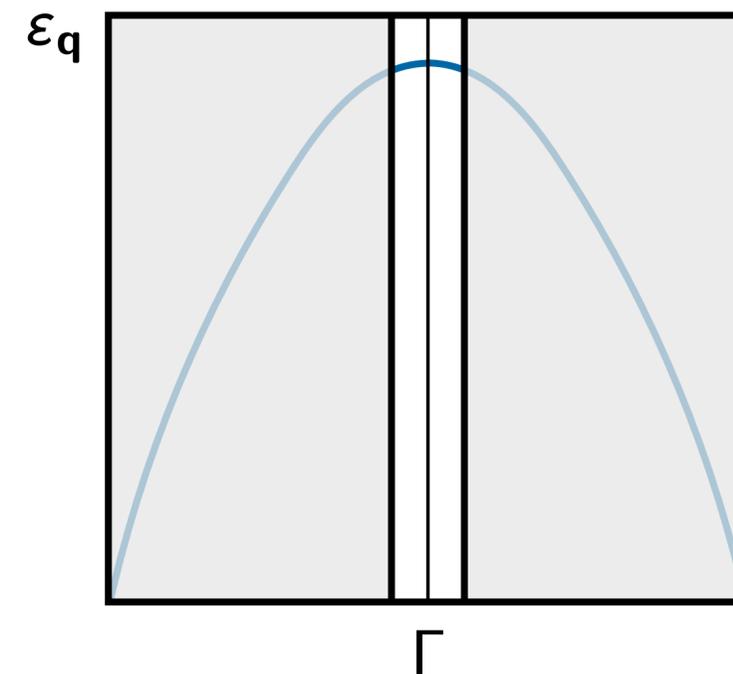


[Pan *et al.*, PRRes '23]

Moiré Brillouin zone:

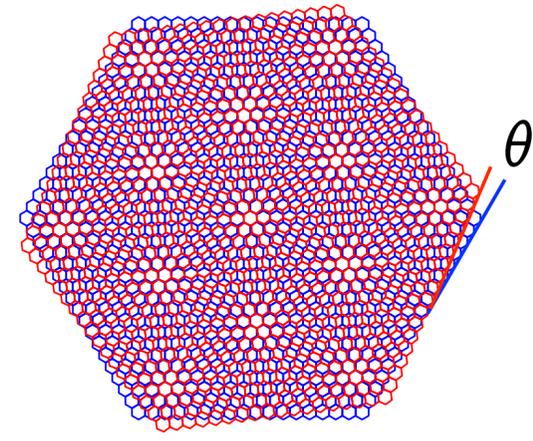


Twisted spectrum:

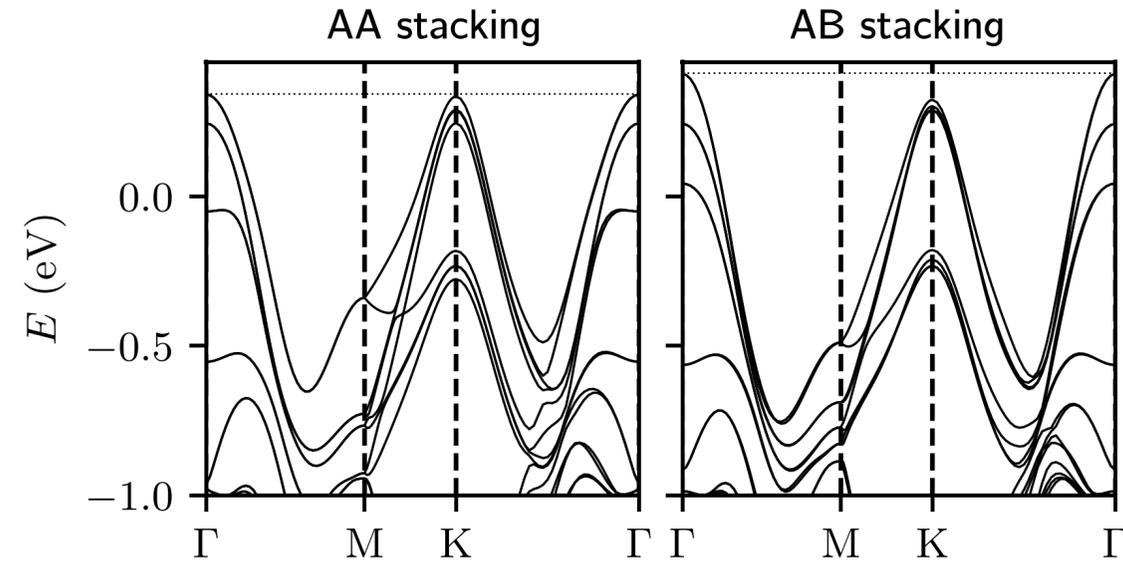


$$\epsilon_{\mathbf{q}} = -\frac{\mathbf{q}^2}{2m^*} + \mathcal{O}(q^3)$$

# Twisted double bilayer WSe<sub>2</sub>: Noninteracting spectrum

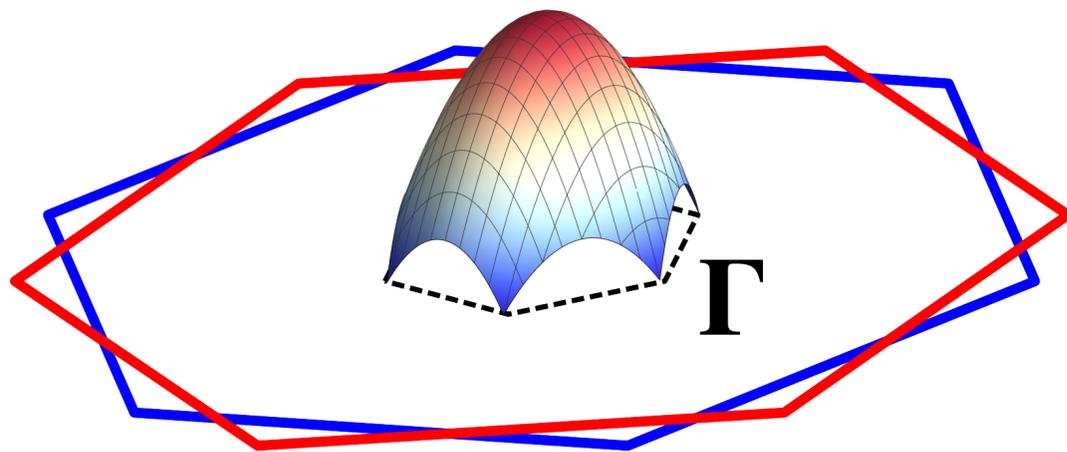


Untwisted spectrum:

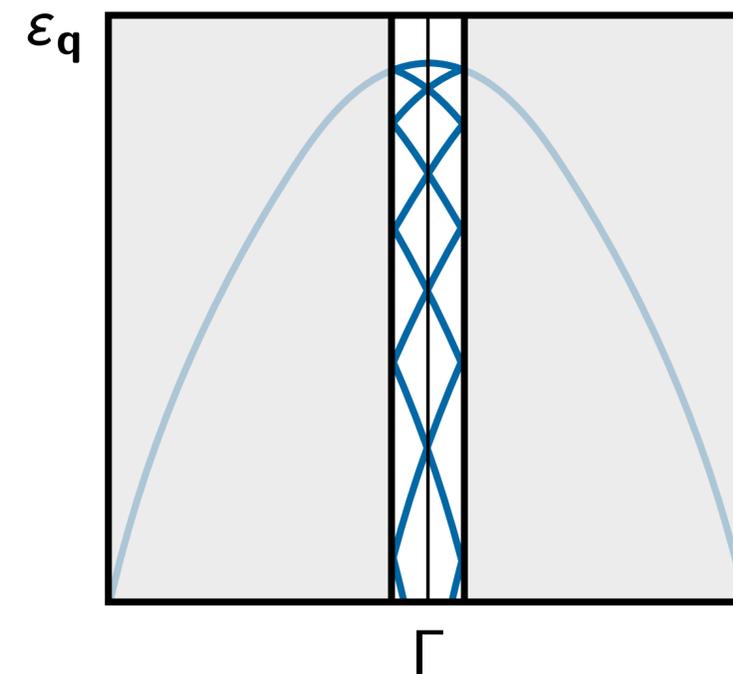


[Pan *et al.*, PRRes '23]

Moiré Brillouin zone:



Twisted spectrum:



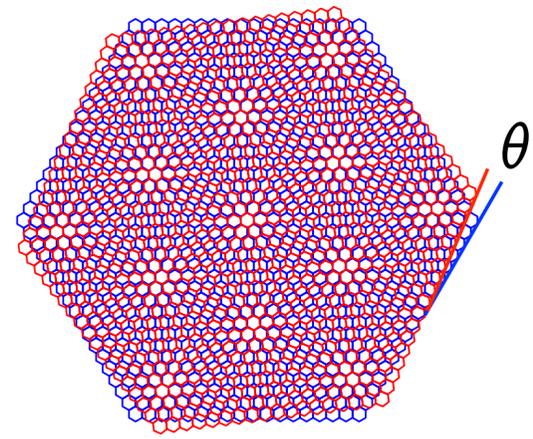
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→ Talk by  
M. Ünzelmann

# Twisted double bilayer WSe<sub>2</sub>: Theory

Interacting Bistritzer-MacDonald model:

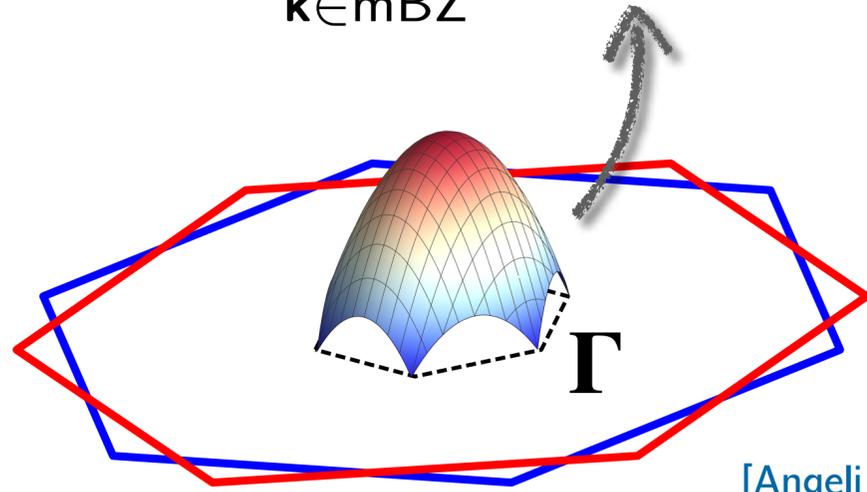
$$\mathcal{H} = \sum_{\mathbf{k} \in \text{mBZ}} c_{\mathbf{k}}^{\dagger} h(\mathbf{k}) c_{\mathbf{k}} - \frac{1}{2A} \sum_{\mathbf{q}} V_{\mathbf{q}} : \rho_{\mathbf{q}} \rho_{-\mathbf{q}} :$$



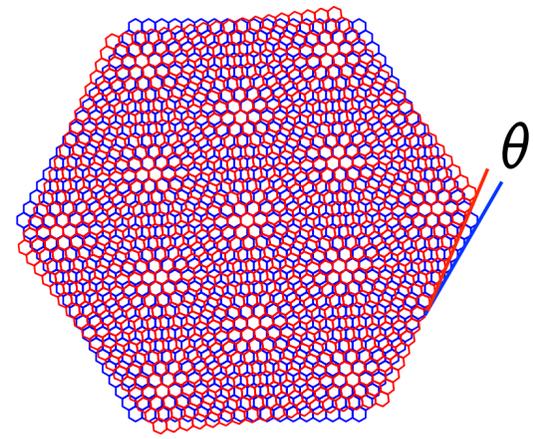
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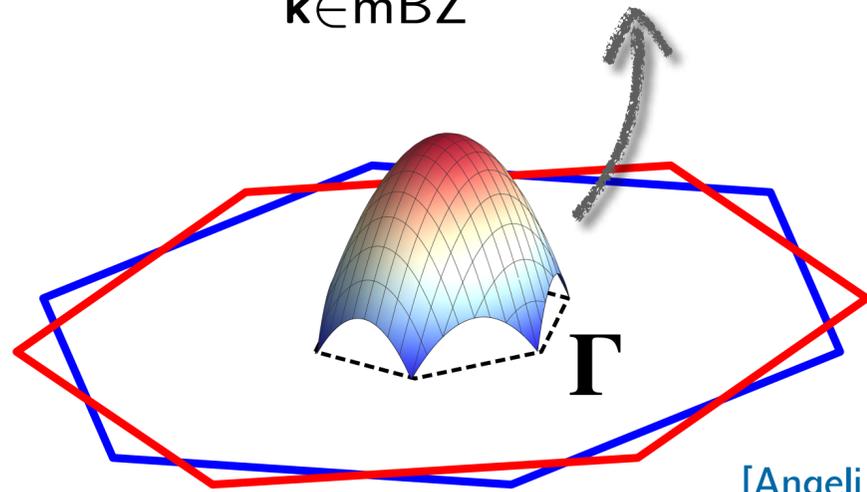
[Angeli & MacDonald, PNAS '21]



# Twisted double bilayer WSe<sub>2</sub>: Theory

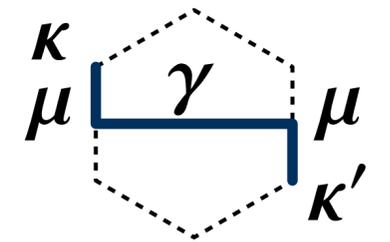
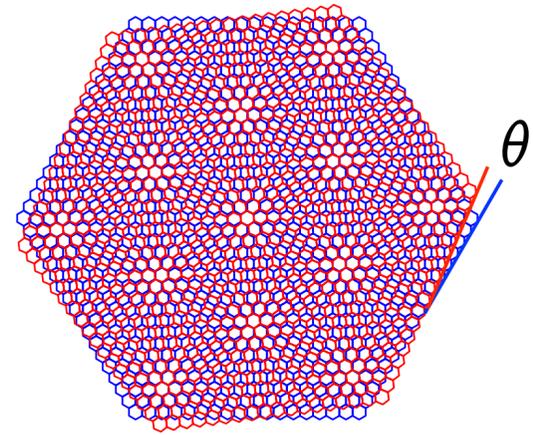
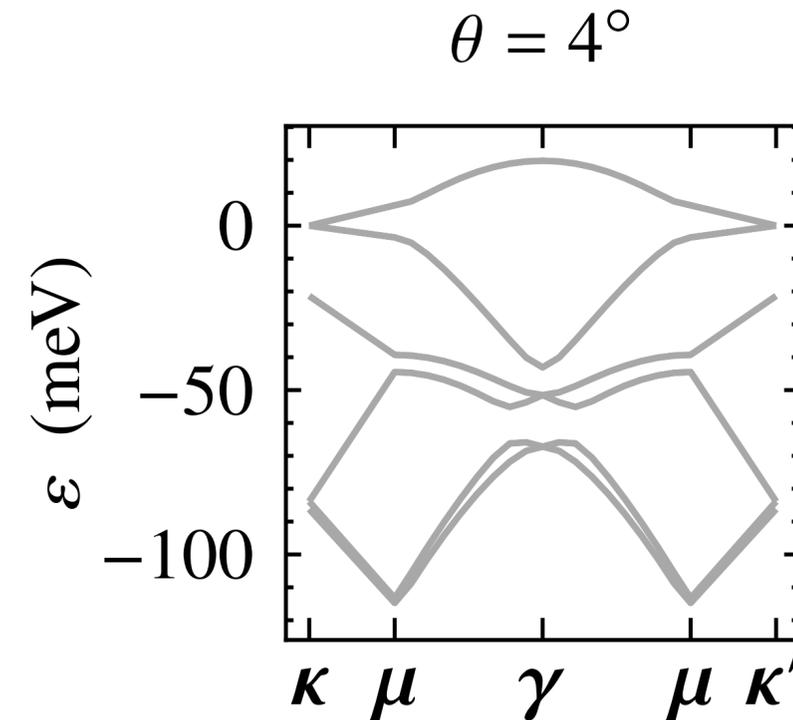
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[Angeli & MacDonald, PNAS '21]

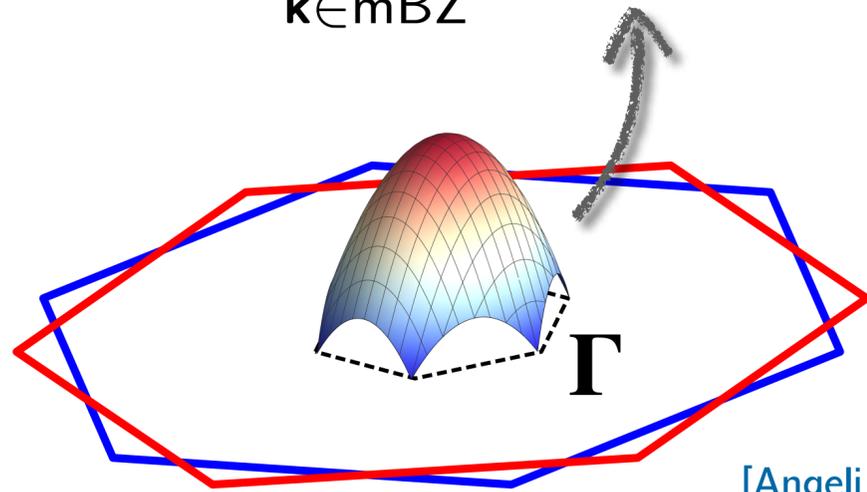
Noninteracting spectrum:



# Twisted double bilayer WSe<sub>2</sub>: Theory

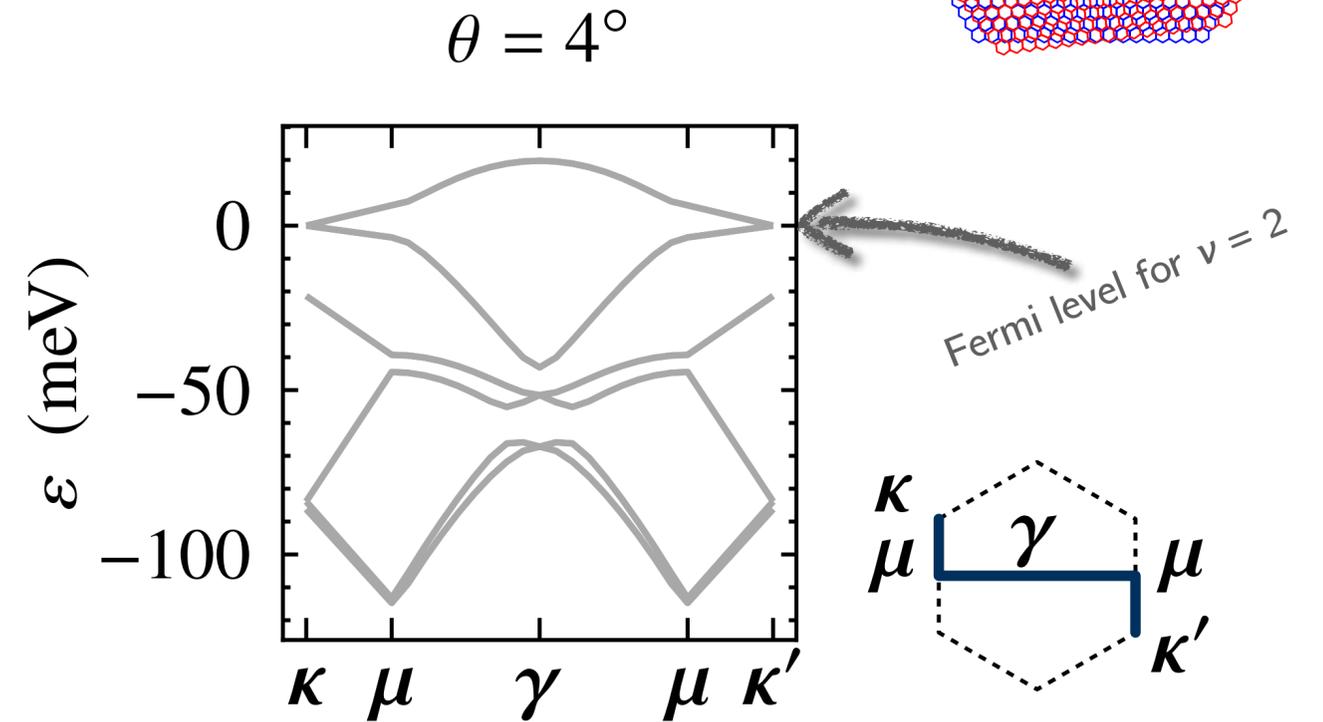
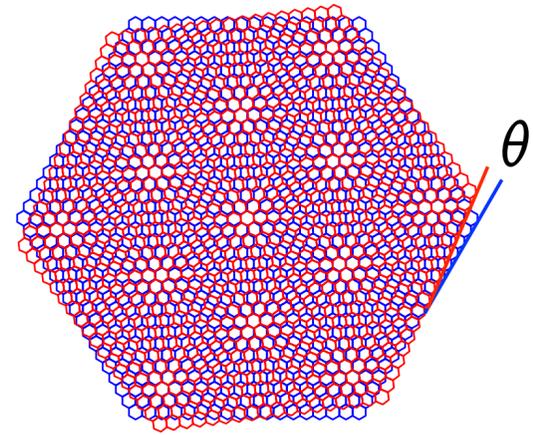
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[Angeli & MacDonald, PNAS '21]

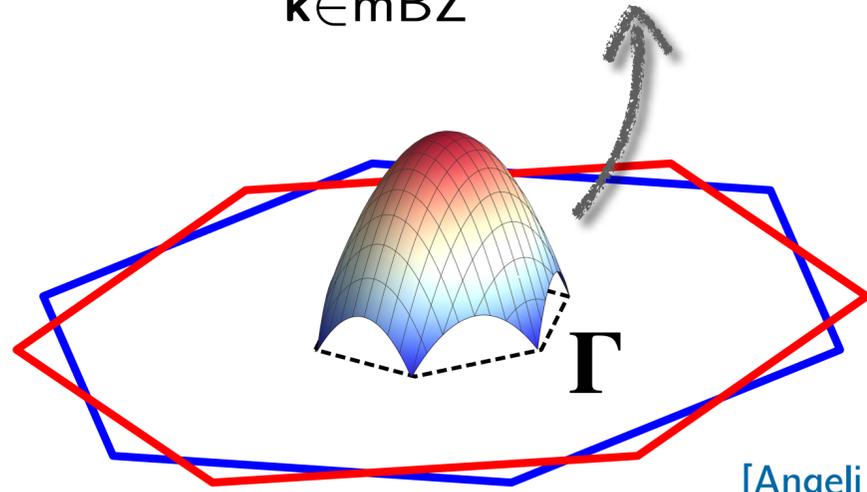
Noninteracting spectrum:



# Twisted double bilayer WSe<sub>2</sub>: Theory

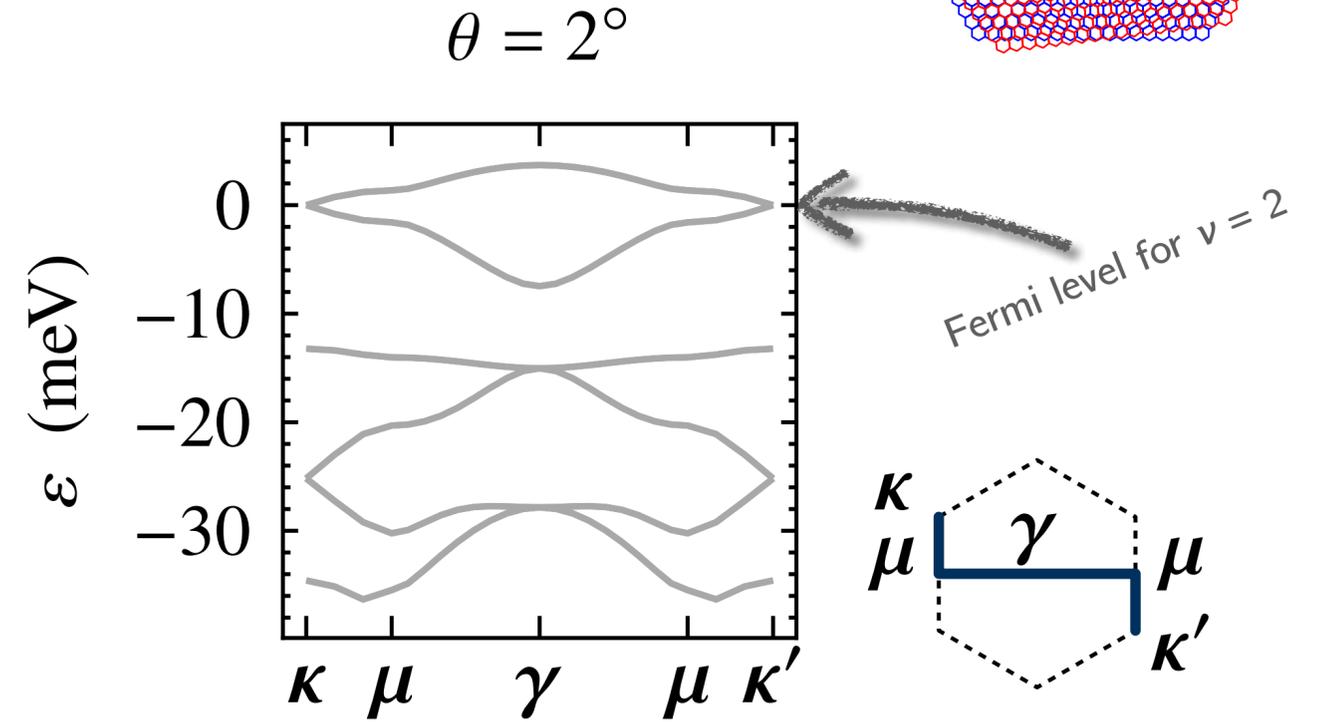
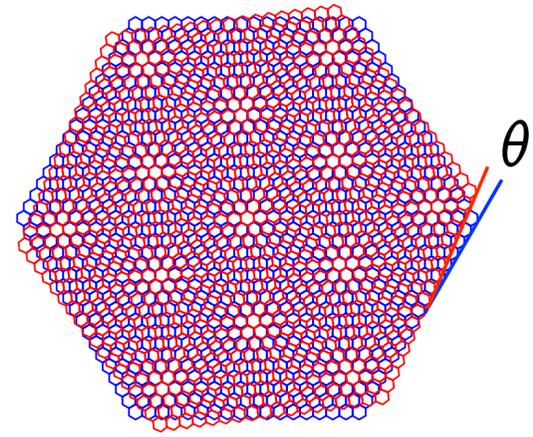
Interacting Bistritzer-MacDonald model:

$$\mathcal{H} = \sum_{\mathbf{k} \in \text{mBZ}} c_{\mathbf{k}}^\dagger h(\mathbf{k}) c_{\mathbf{k}} - \frac{1}{2A} \sum_{\mathbf{q}} V_{\mathbf{q}} : \rho_{\mathbf{q}} \rho_{-\mathbf{q}} :$$



[Angeli & MacDonald, PNAS '21]

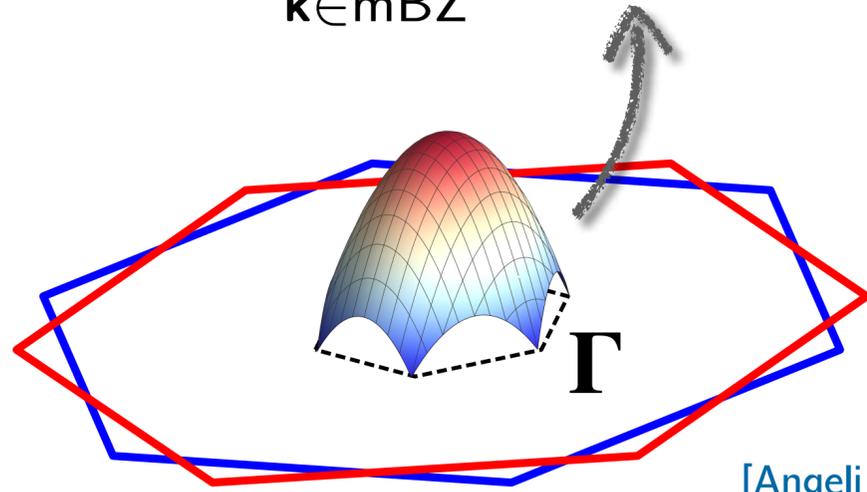
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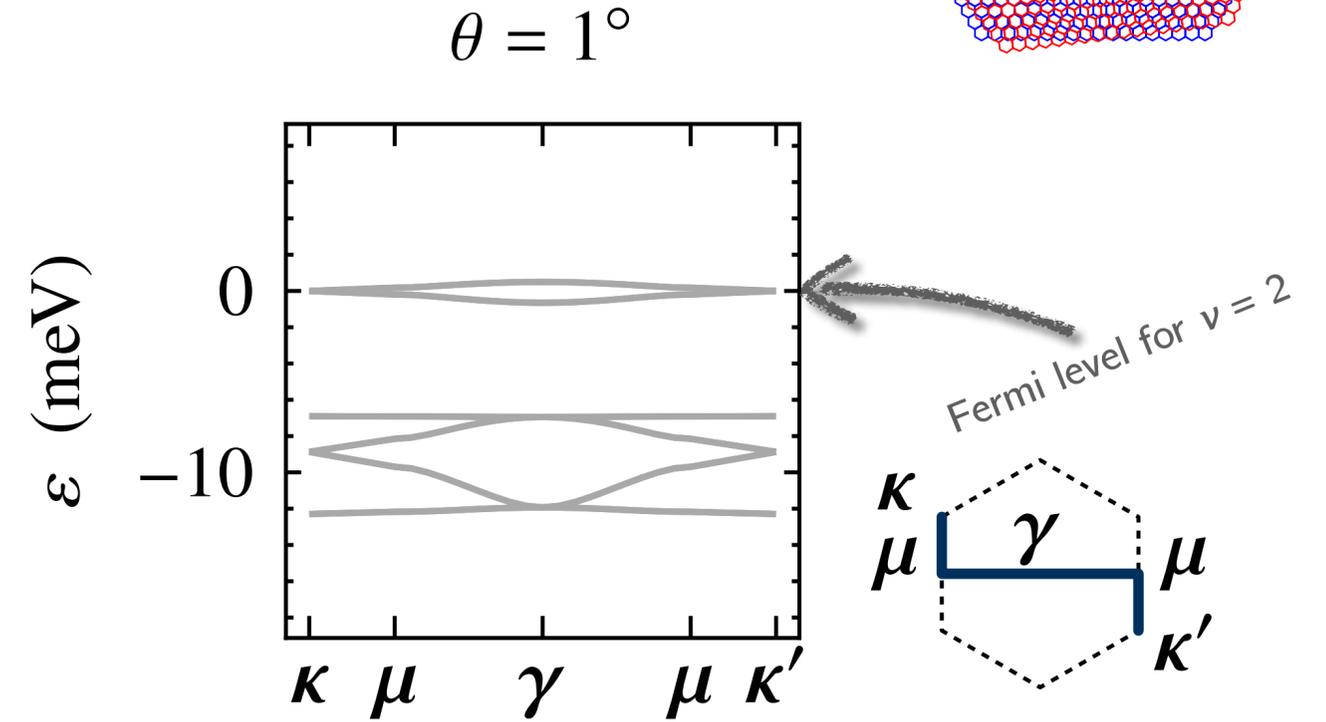
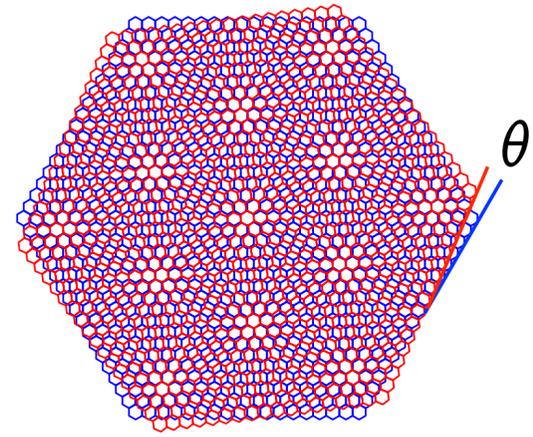
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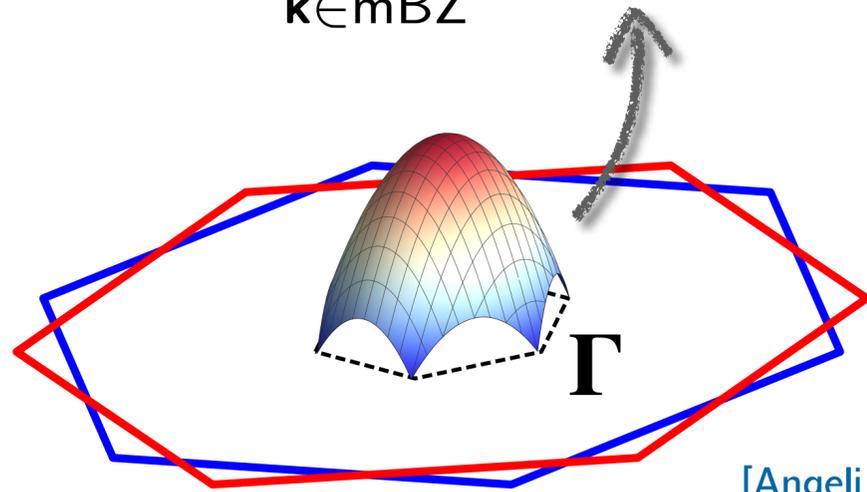
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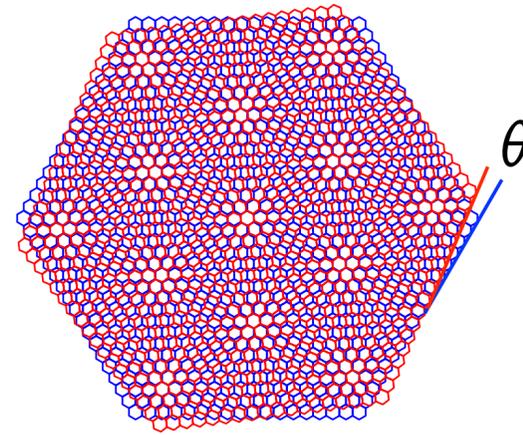
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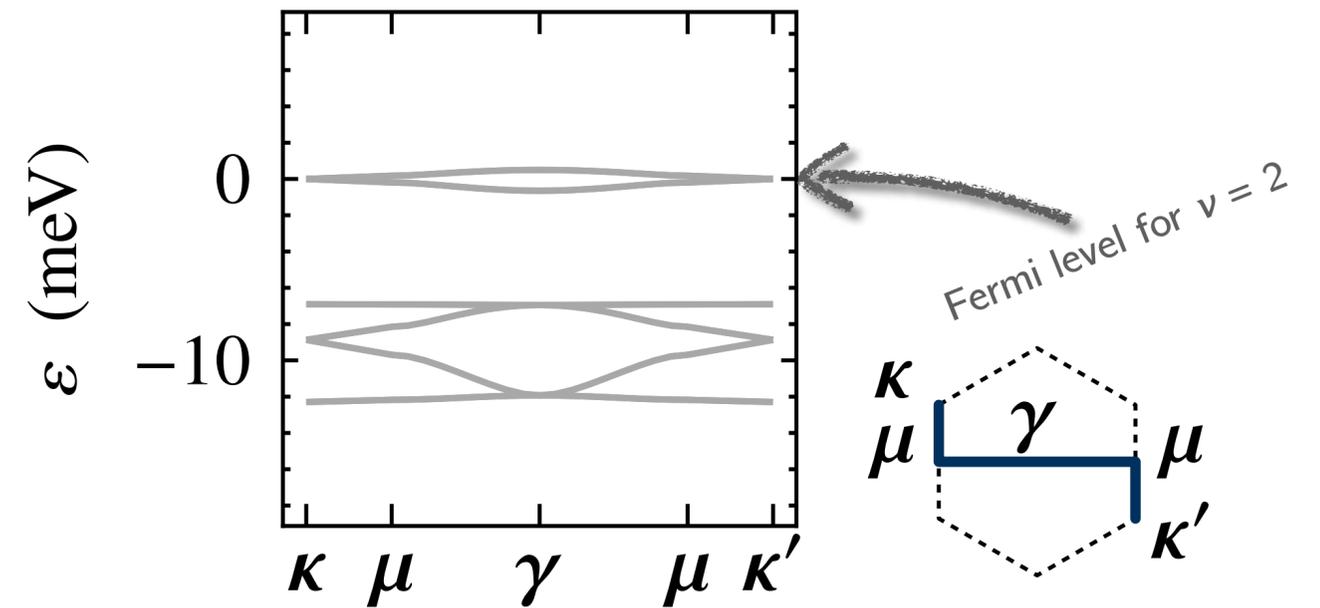


[Angeli & MacDonald, PNAS '21]

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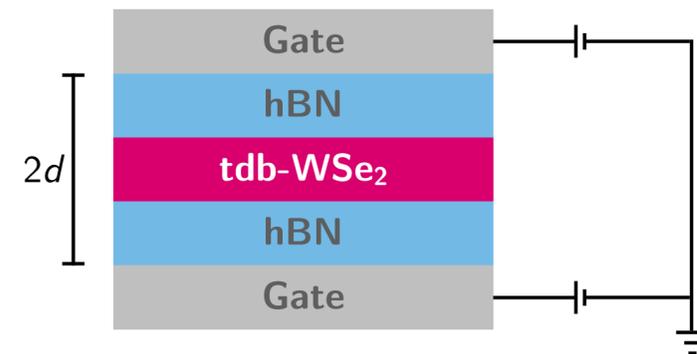


$\theta = 1^\circ$



Coulomb interaction:

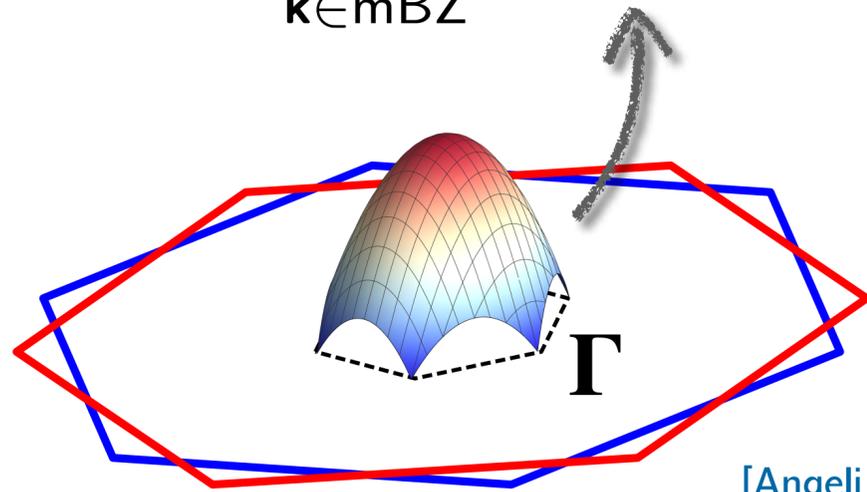
$$V_{\mathbf{q}} = \frac{e^2}{2\epsilon_0 \epsilon_{\text{eff}} |\mathbf{q}|} \tanh(|\mathbf{q}|d)$$



# Twisted double bilayer WSe<sub>2</sub>: Theory

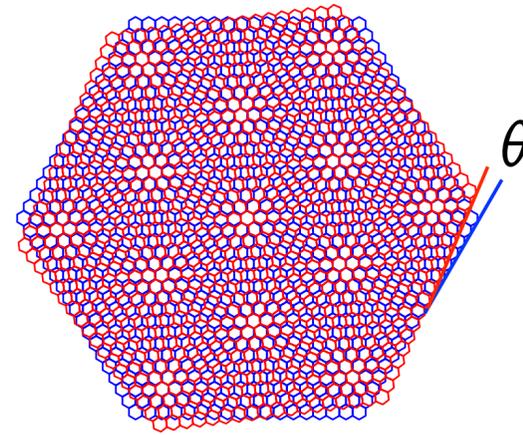
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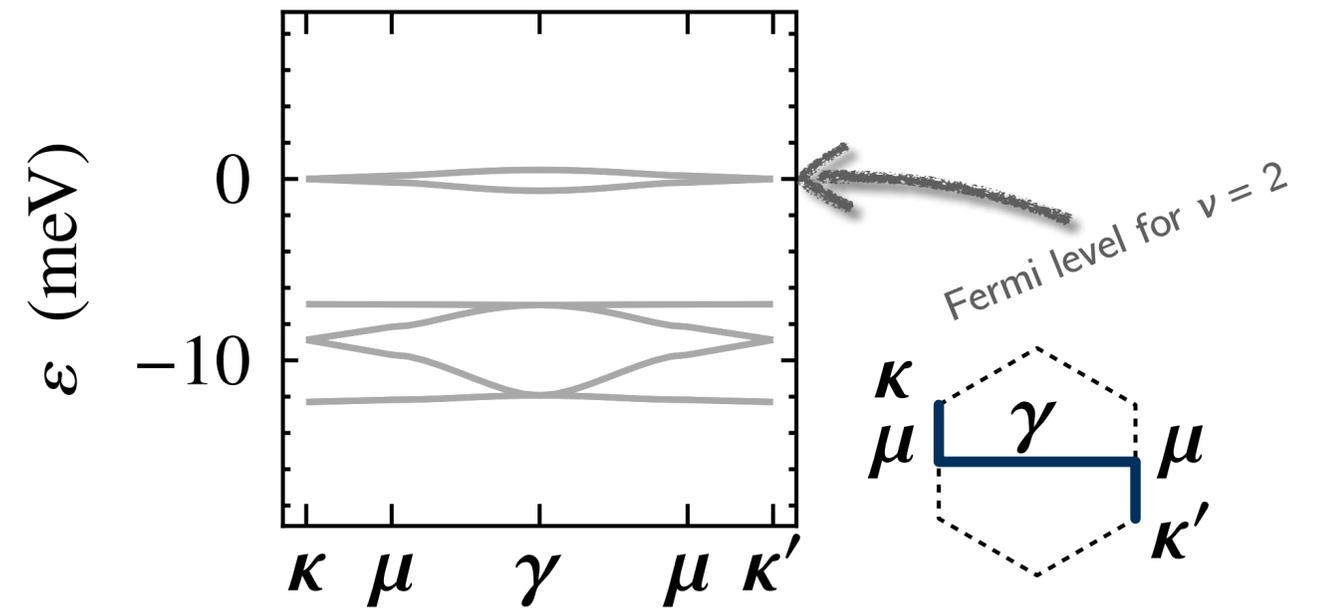


[Angeli & MacDonald, PNAS '21]

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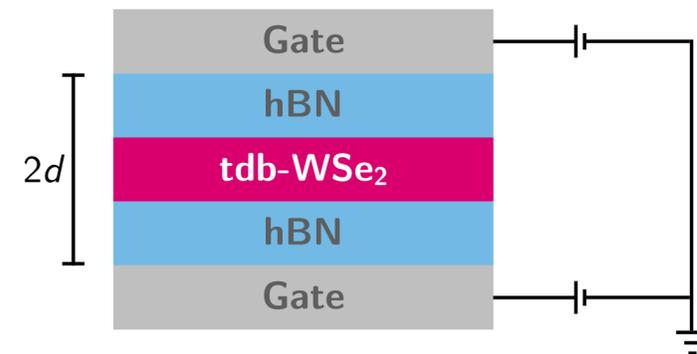
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Coulomb interaction:

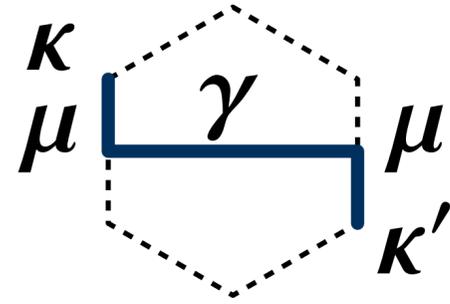
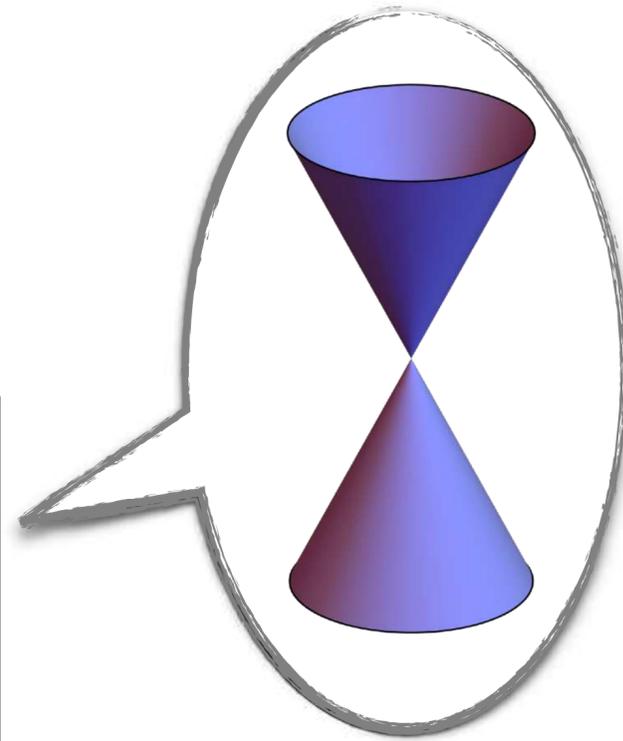
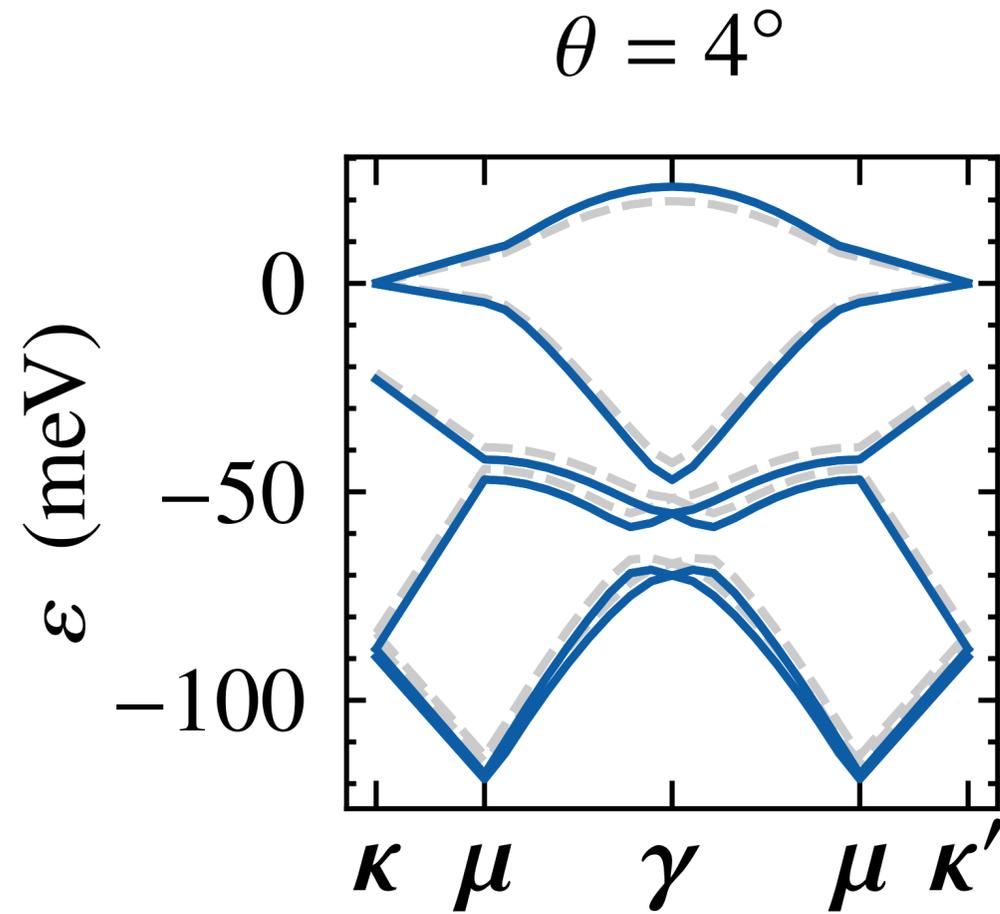
$$V_{\mathbf{q}} = \frac{e^2}{2\epsilon_0 \epsilon_{\text{eff}} |\mathbf{q}|} \tanh(|\mathbf{q}|d)$$

Effective permittivity  
 $\epsilon_{\text{eff}} \lesssim \mathcal{O}(100)$



# Large $\theta$ : Dirac semimetal

Interacting spectrum:



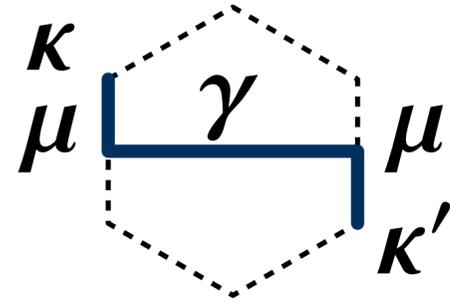
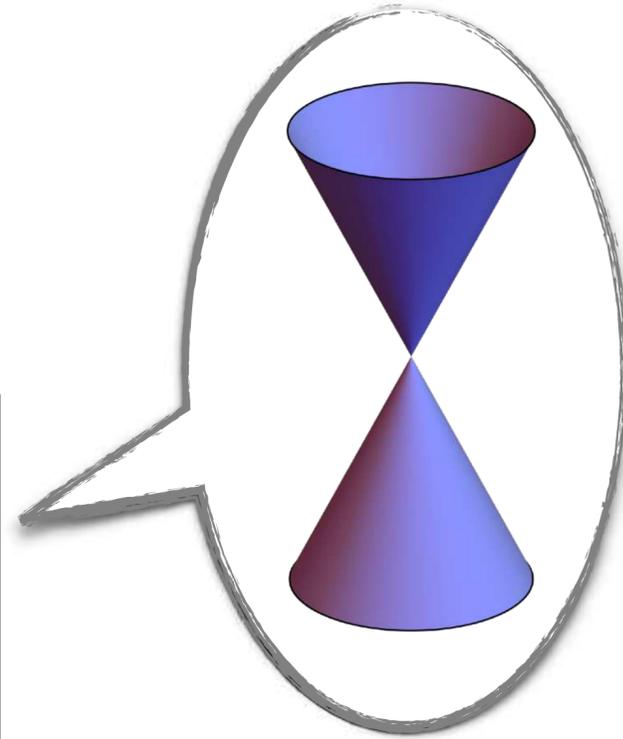
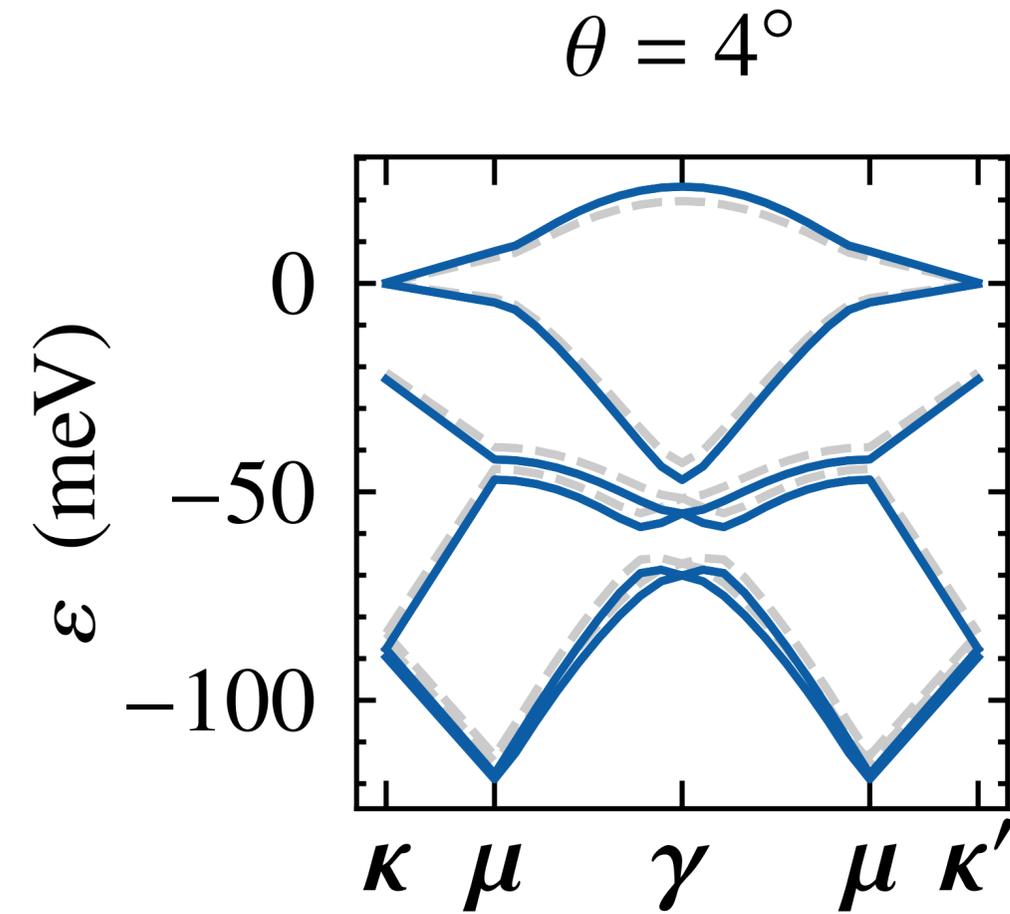
... from selfconsistent Hartree-Fock analysis



... for  $\nu = 2$  holes per unit cell

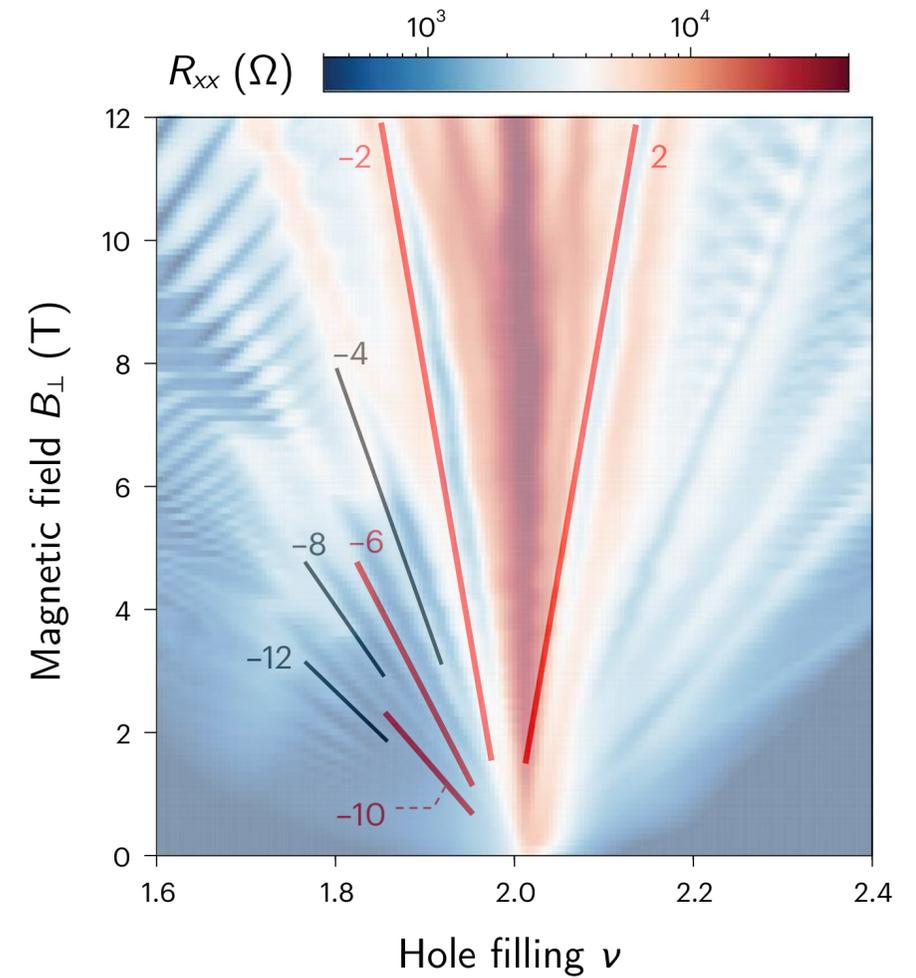
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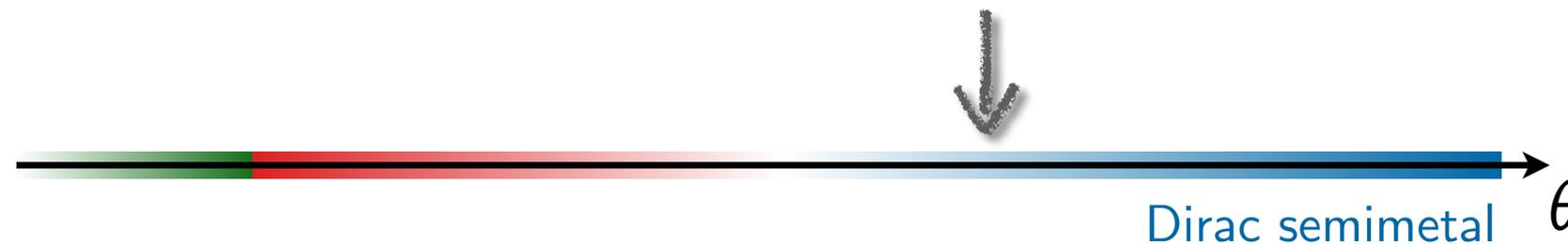


... from selfconsistent Hartree-Fock analysis

Magnetotransport:



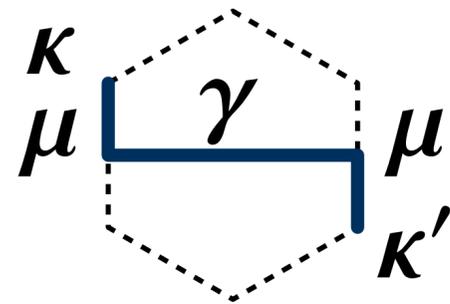
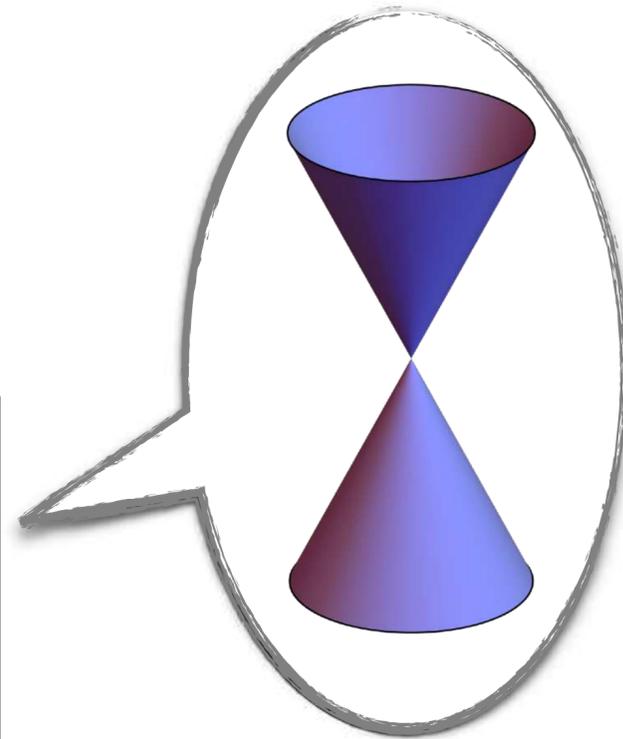
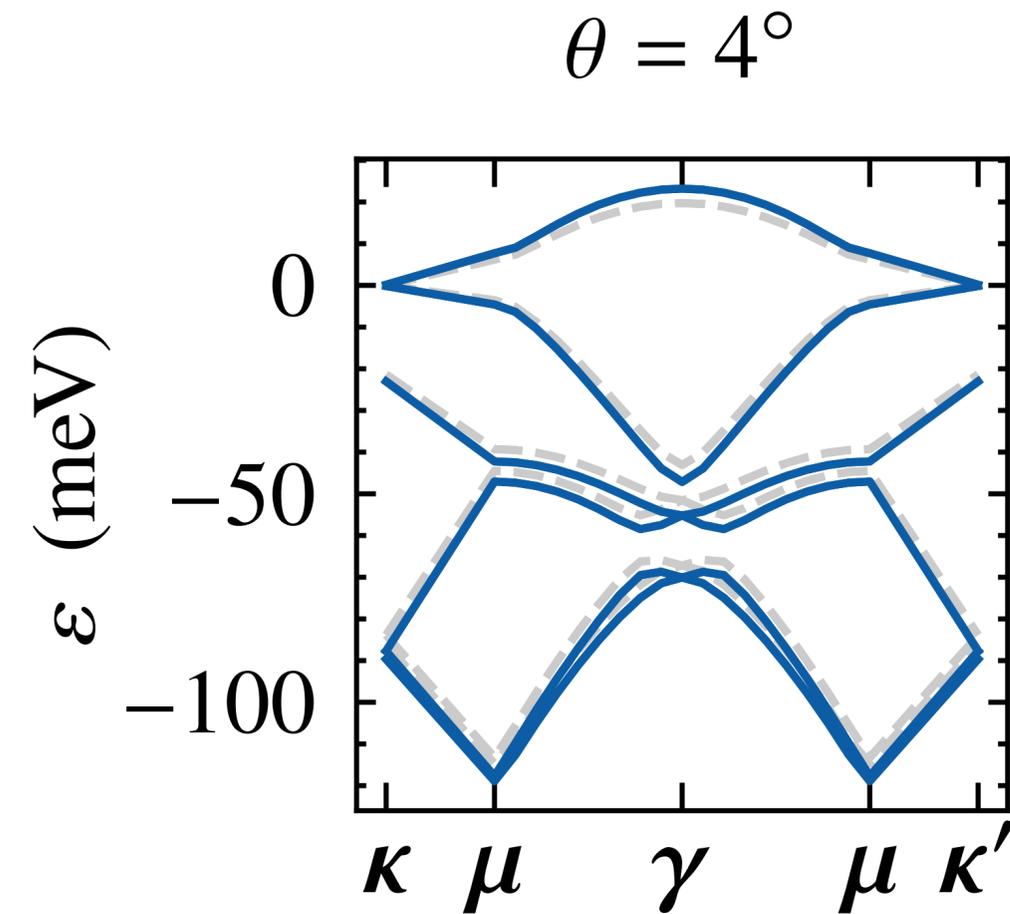
[Ma *et al.*, Nat. Mater. '25]



... for  $\nu = 2$  holes per unit cell

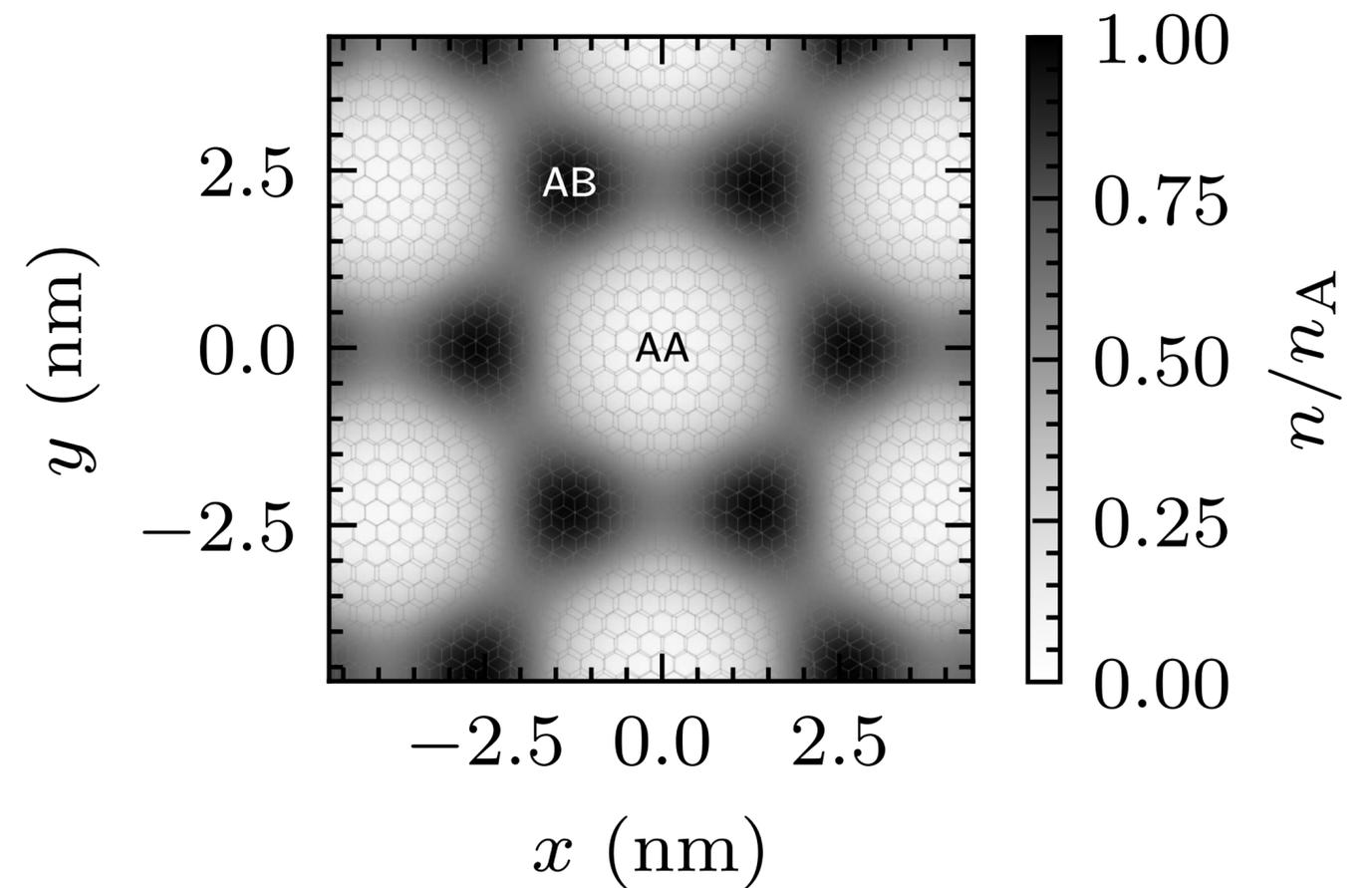
# Large $\theta$ : Dirac semimetal

Interacting spectrum:



... from selfconsistent Hartree-Fock analysis

Charge density:

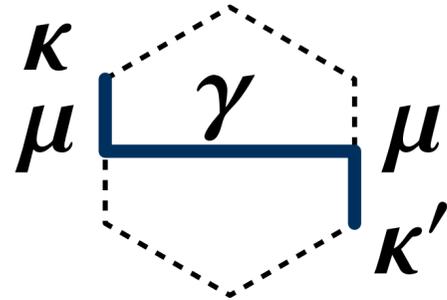
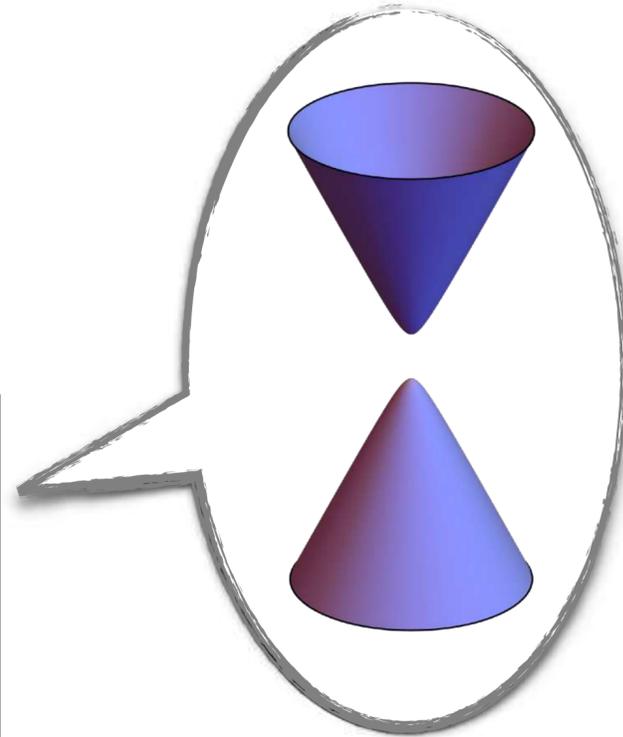
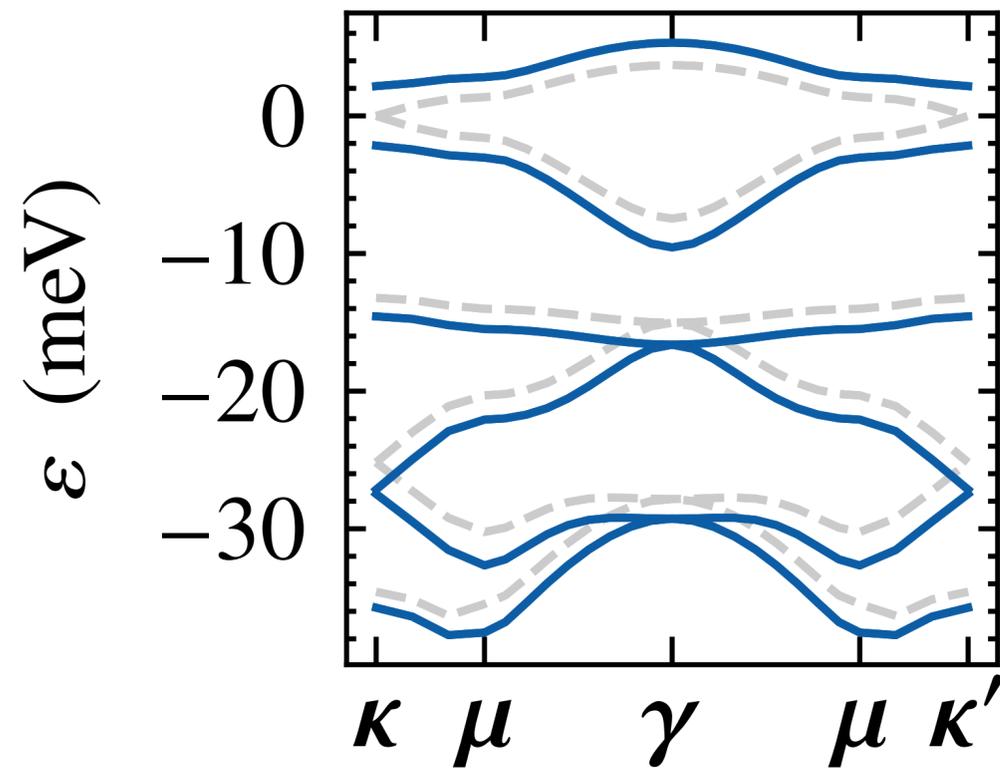


... for  $\nu = 2$  holes per unit cell

# Small $\theta$ : Antiferromagnetic insulator

Interacting spectrum:

$$\theta = 2^\circ$$



... from selfconsistent Hartree-Fock analysis

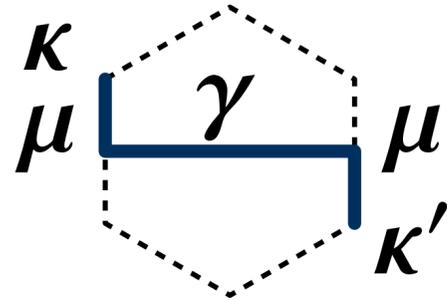
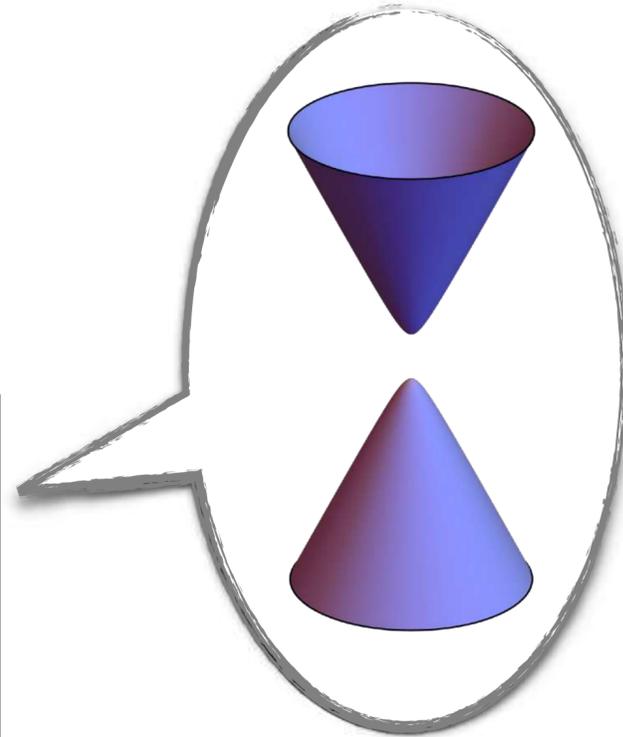
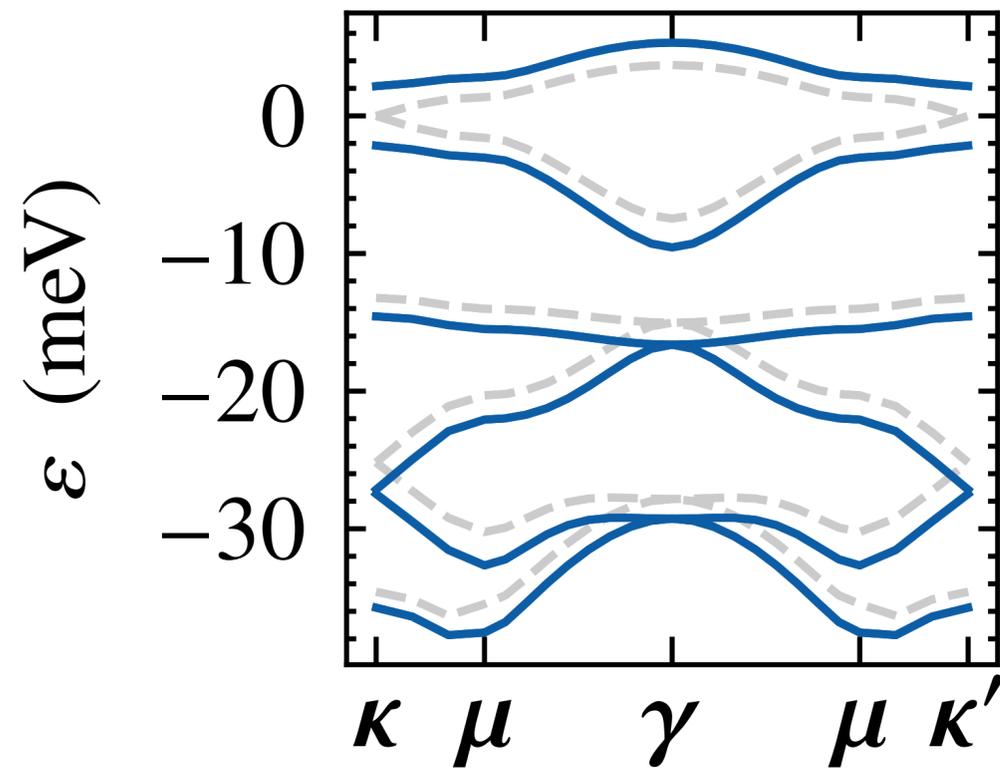


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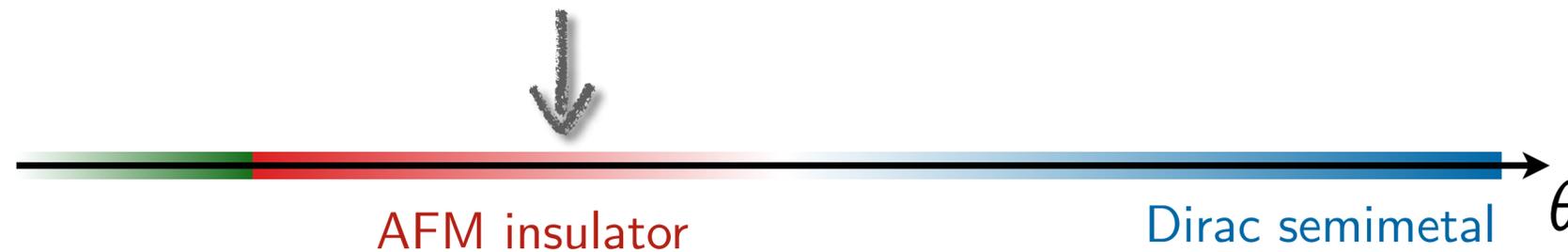
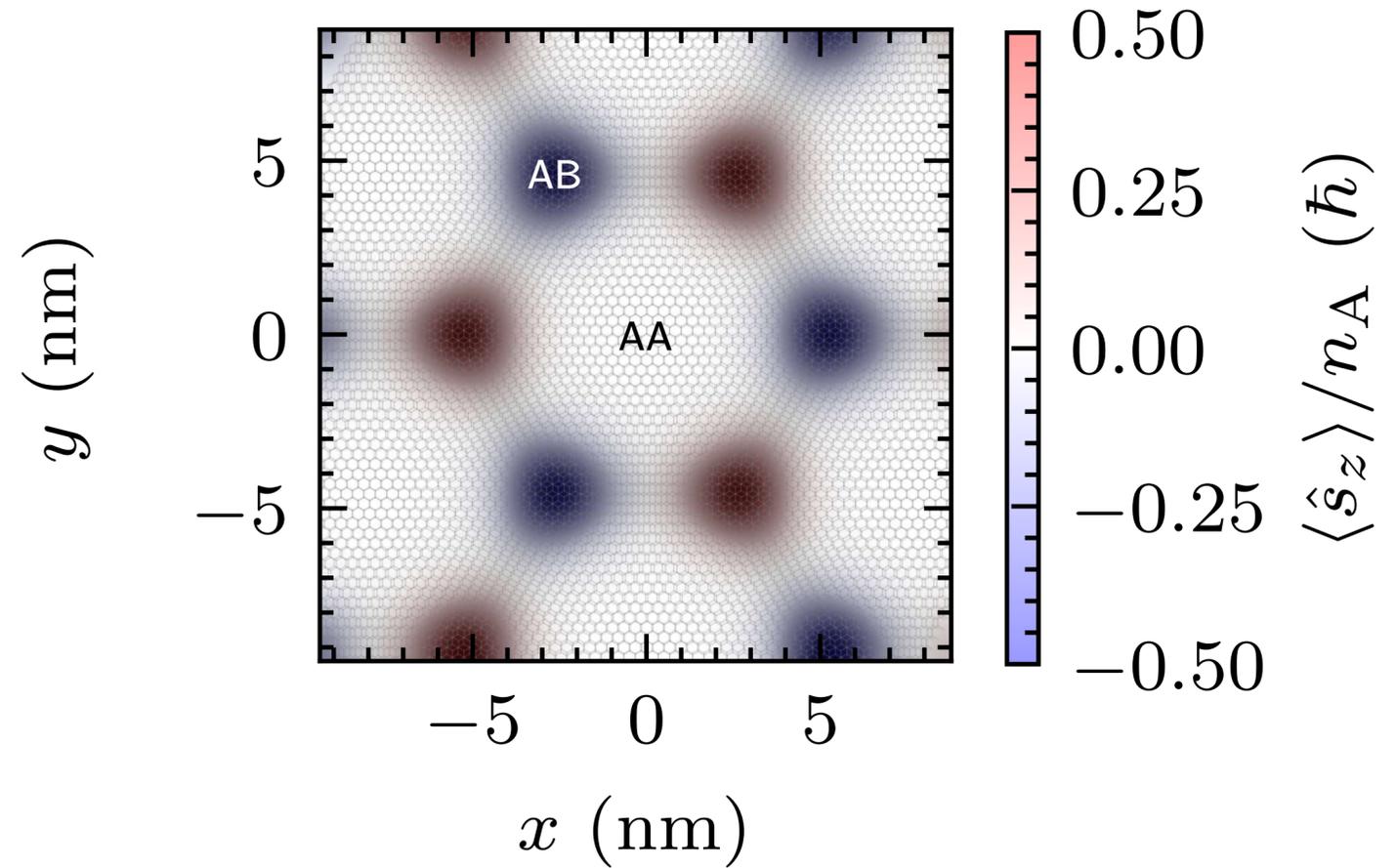
Interacting spectrum:

$$\theta = 2^\circ$$



... from selfconsistent Hartree-Fock analysis

Spin density:

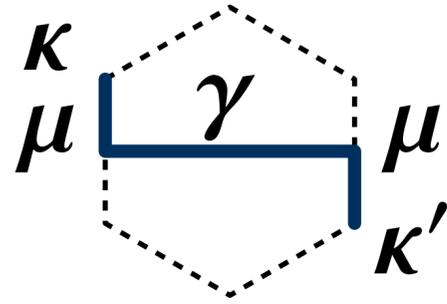
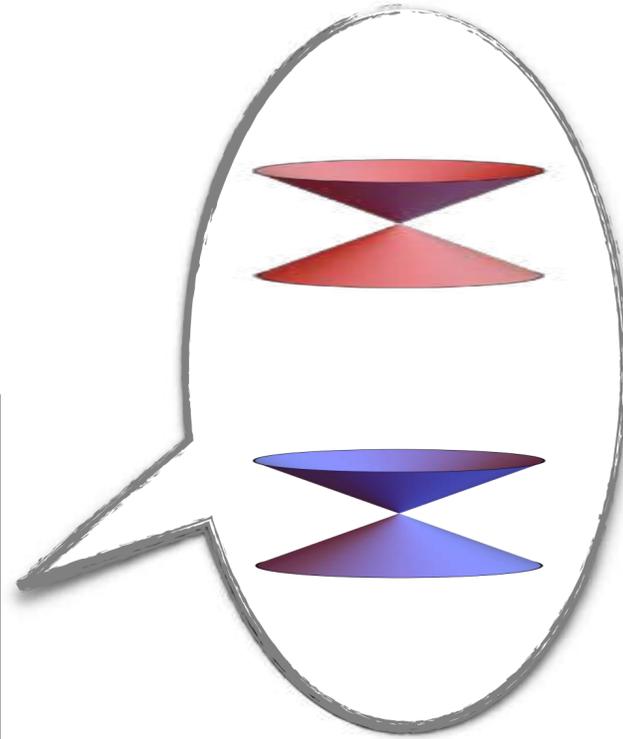
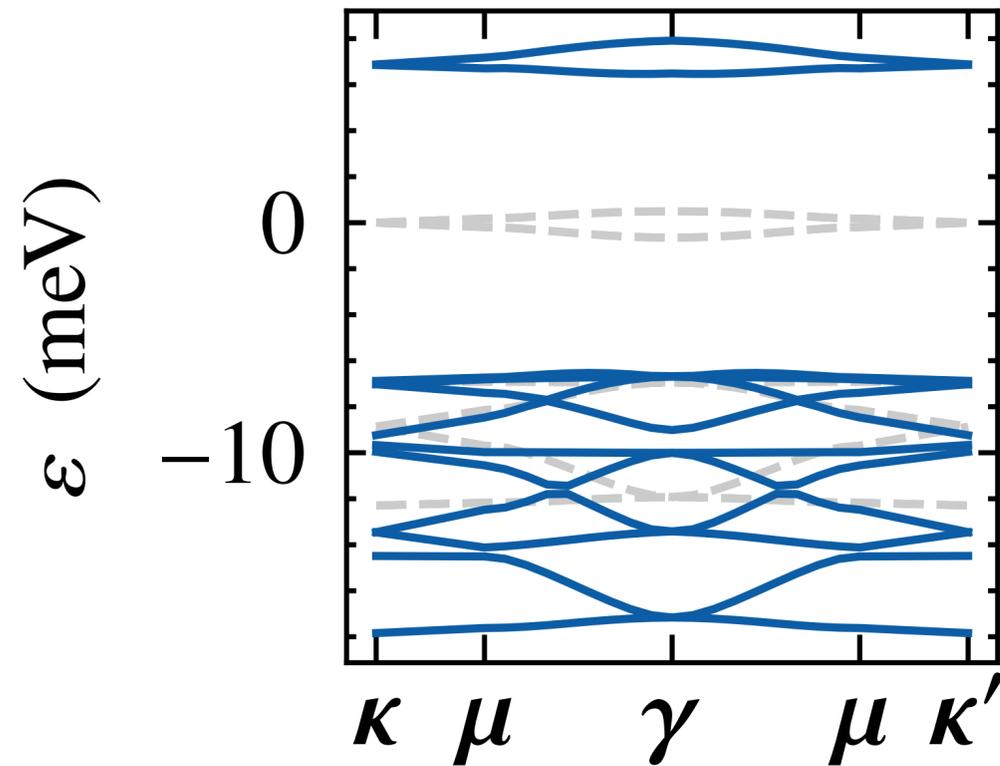


... for  $\nu = 2$  holes per unit cell

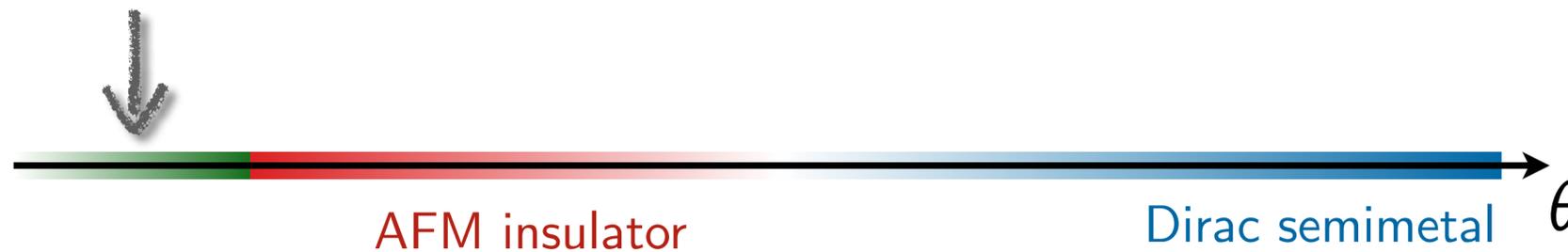
# Very small $\theta$ : Ferromagnetic insulator

Interacting spectrum:

$$\theta = 1^\circ$$



... from selfconsistent Hartree-Fock analysis

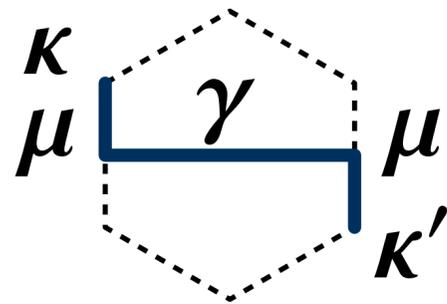
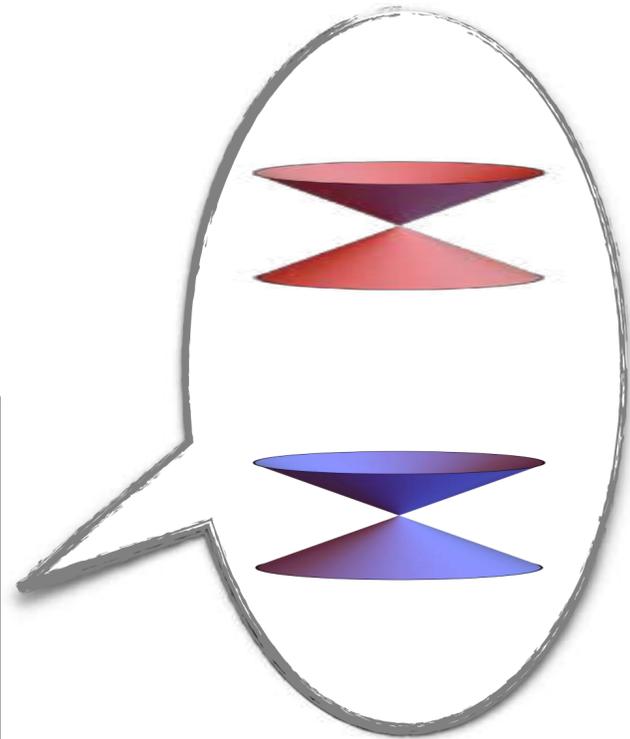
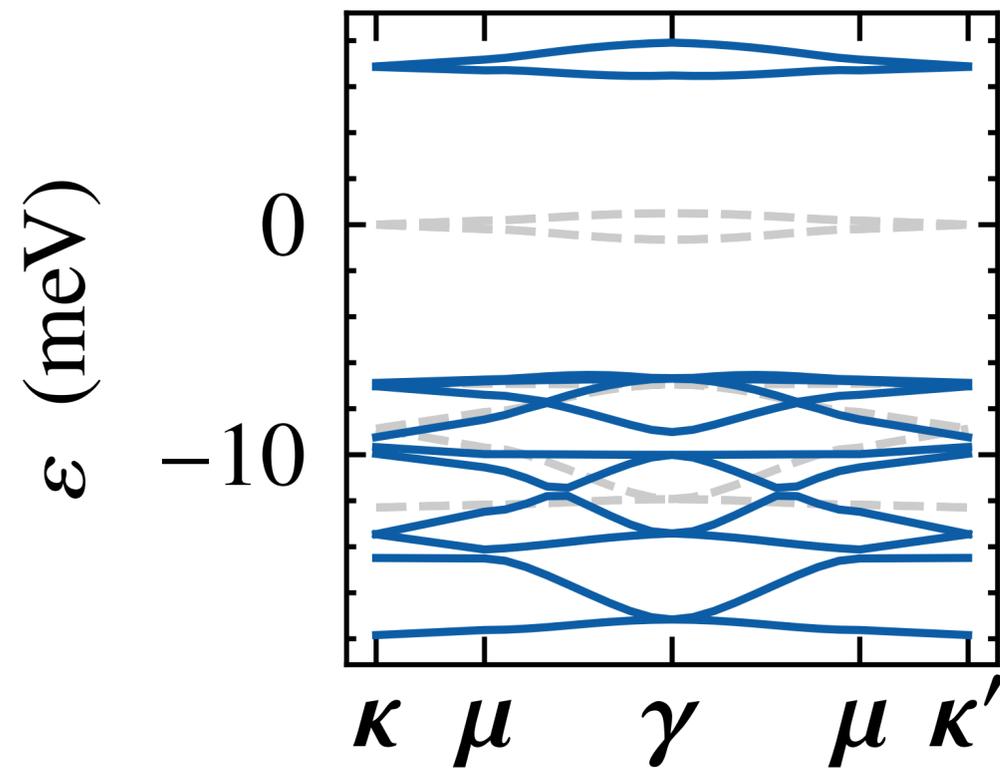


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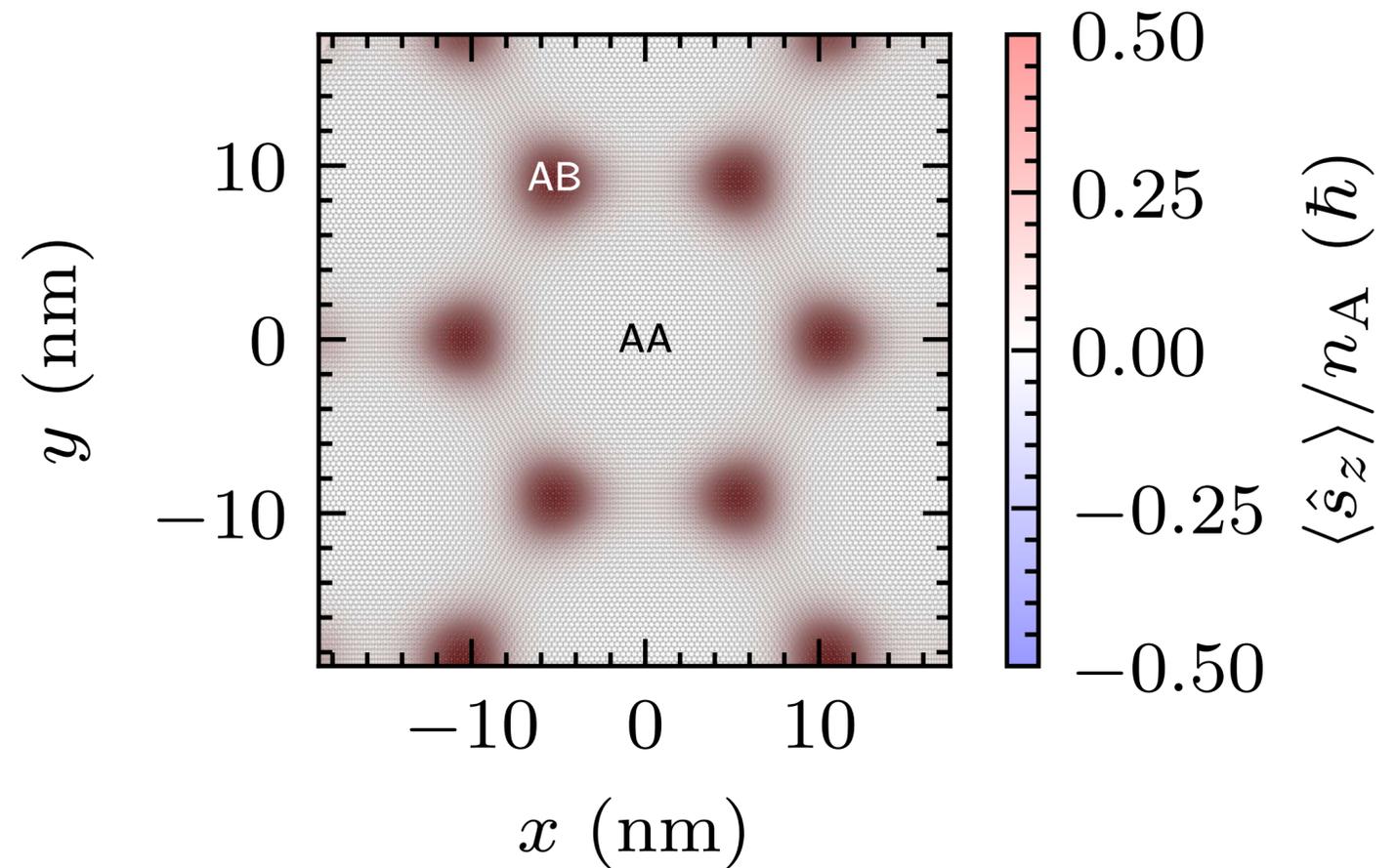
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$$\theta = 1^\circ$$



Spin density:

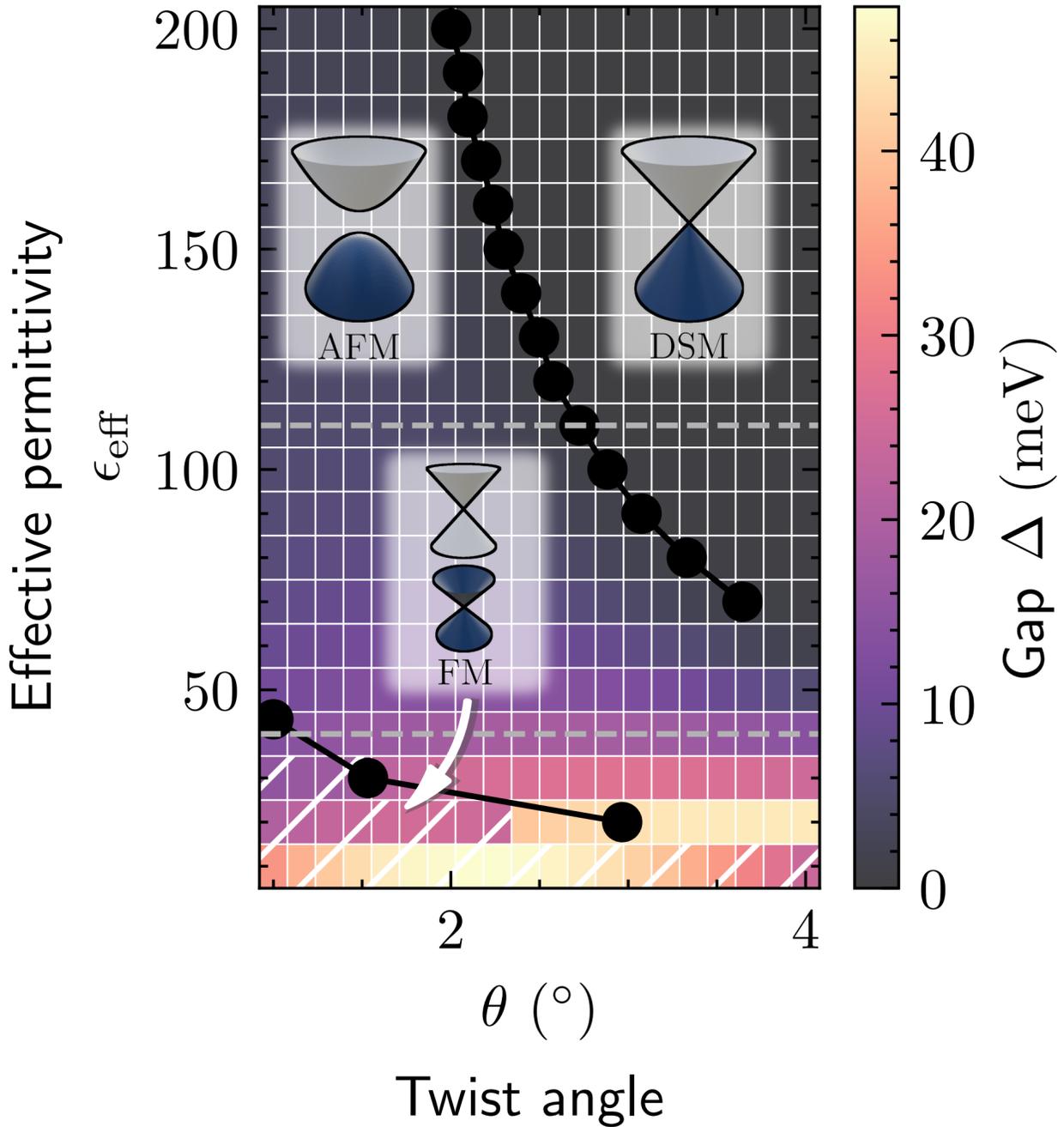


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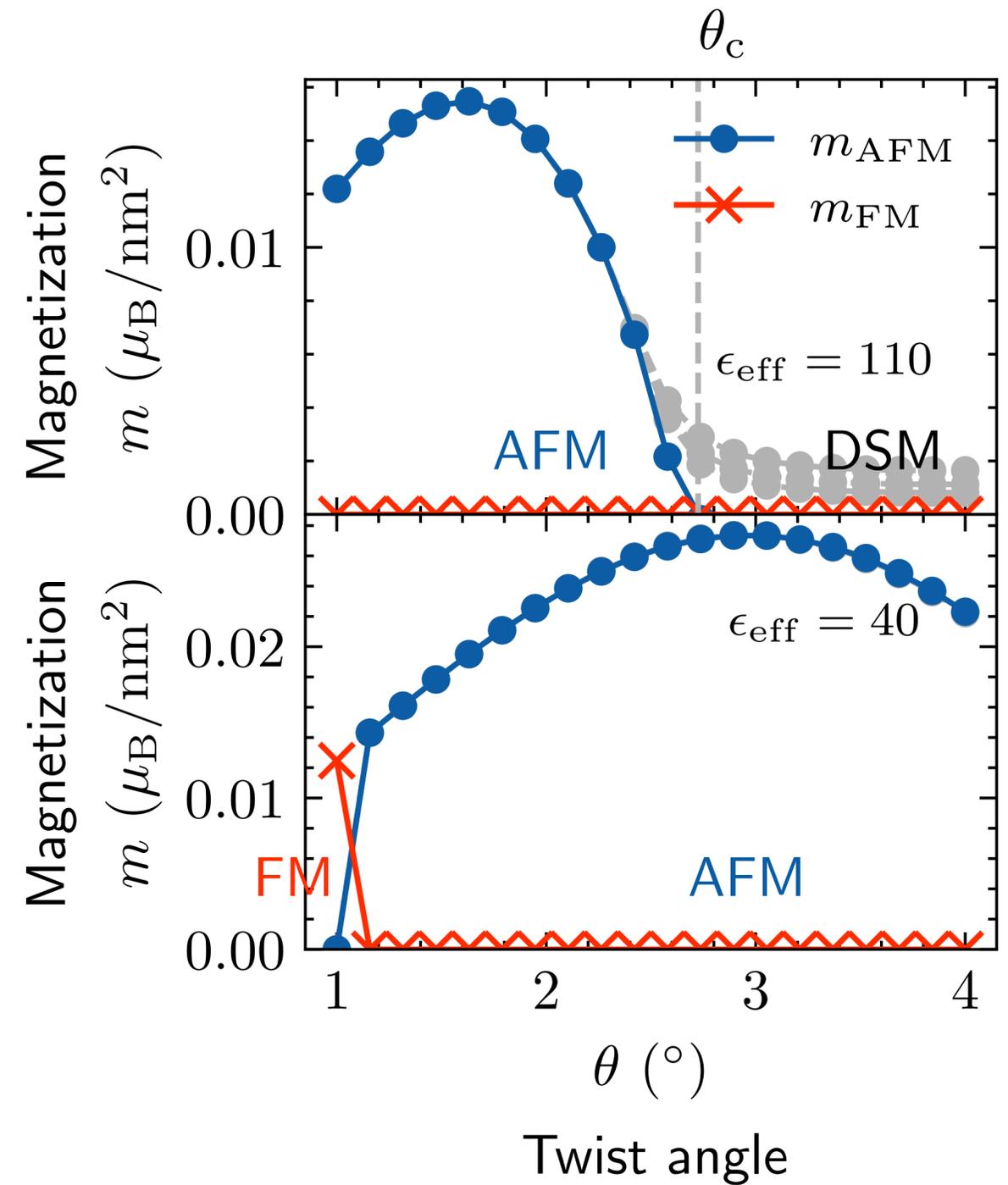
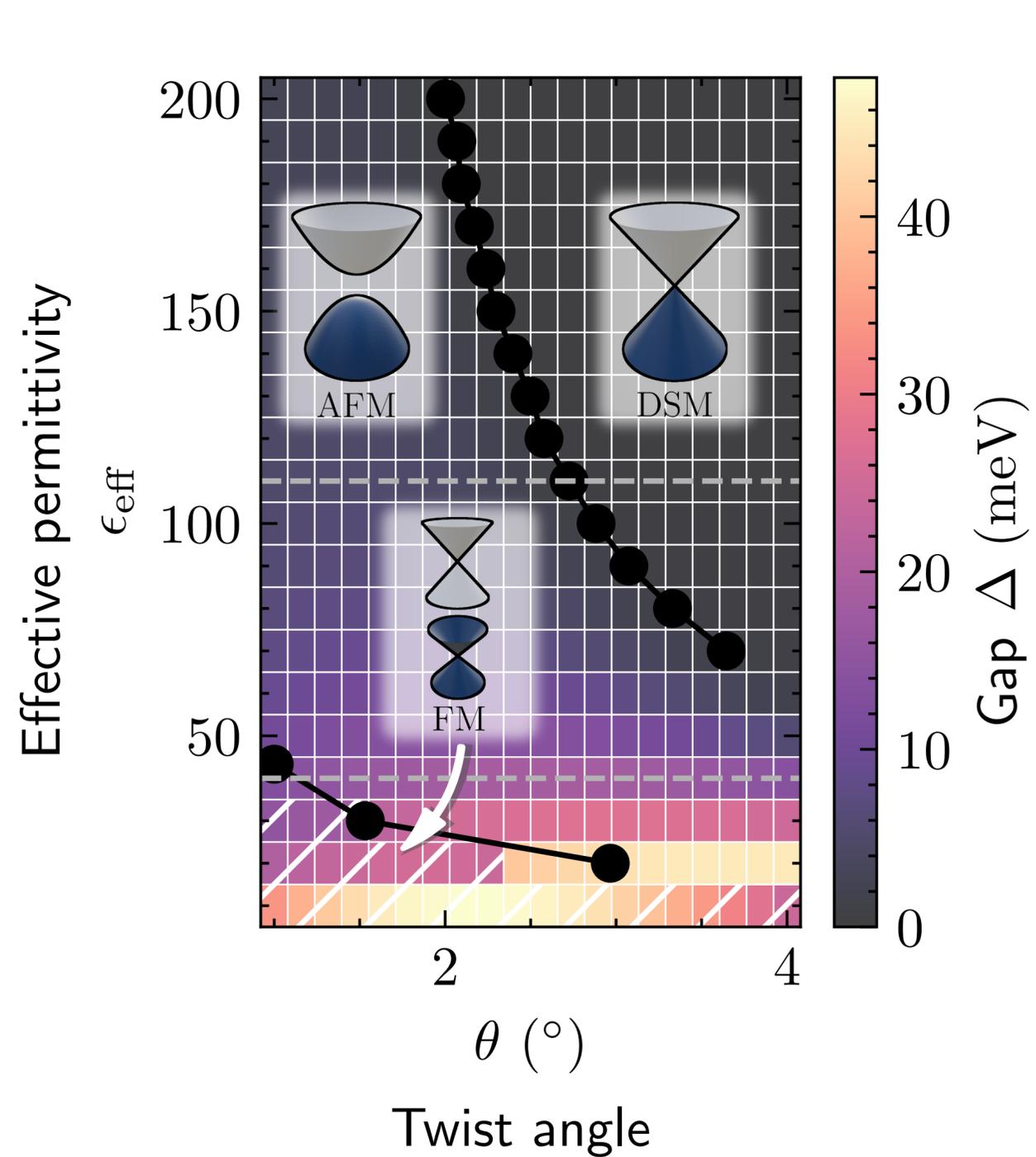
... for  $\nu = 2$  holes per unit cell

# Quantum phase diagram



... for  $\nu = 2$  holes per unit cell

# Quantum phase diagram



... for  $\nu = 2$  holes per unit cell

# DSM-to-AFM quantum criticality

Universal field theory:

$$\mathcal{L} = \frac{1}{2} \vec{\varphi} \cdot (-\partial_\tau^2 - v_B^2 \vec{\nabla}^2) \vec{\varphi} + \psi^\dagger \left( \partial_\tau + v_F \gamma_0 \vec{\gamma} \cdot \vec{\nabla} \right) \psi + g \vec{\varphi} \cdot \psi^\dagger \gamma_0 \vec{\sigma} \psi + \dots$$

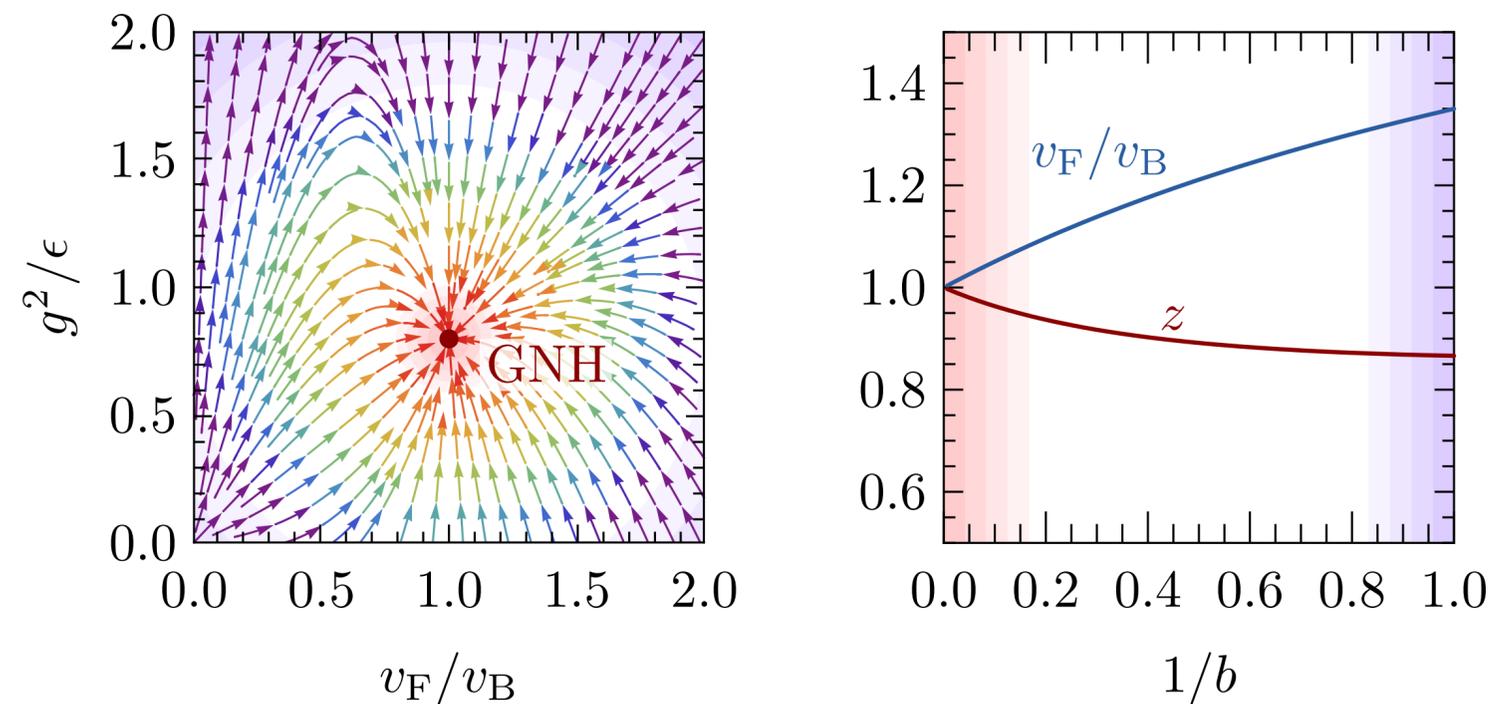
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Renormalization group flow:

... in  $d = 3 - \epsilon$  spatial dimensions



[Biedermann & LJ, arXiv:2509.04561]

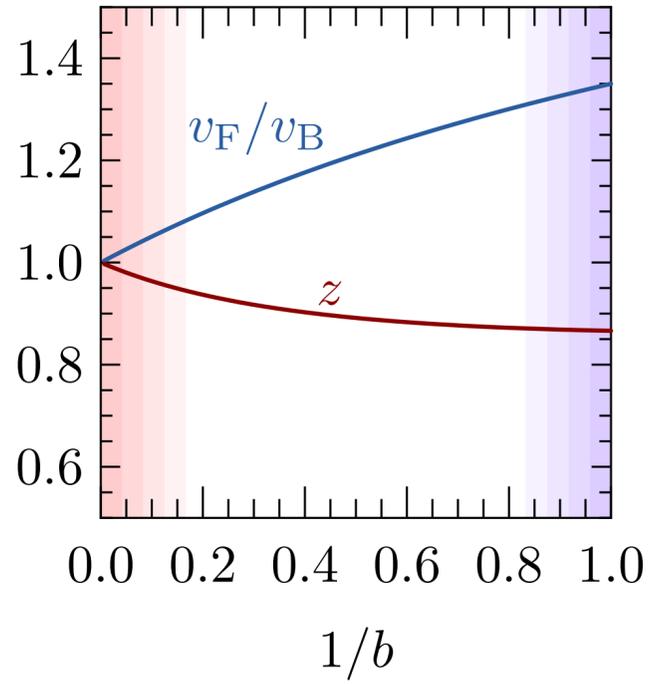
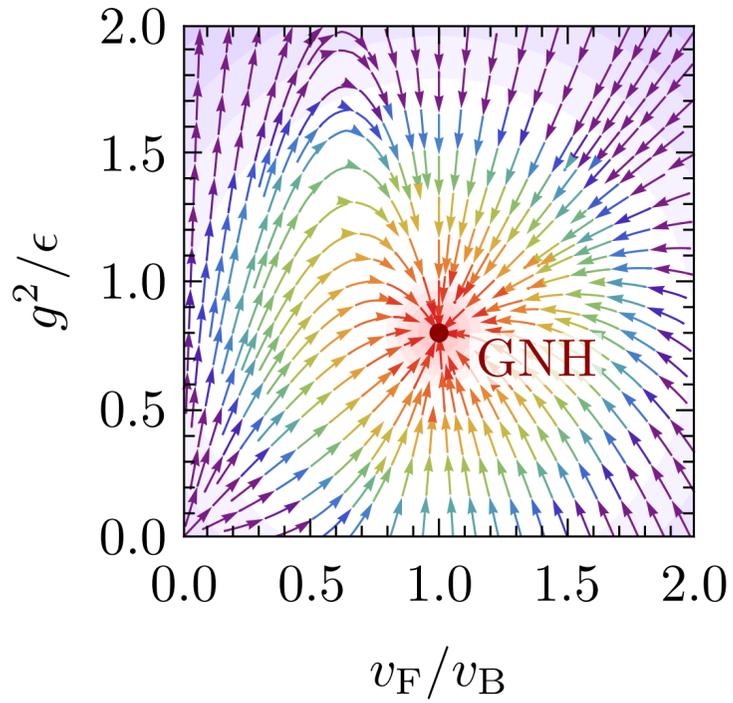
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**⇒ Emergent Lorentz symmetry**

[Biedermann & LJ, arXiv:2509.04561]

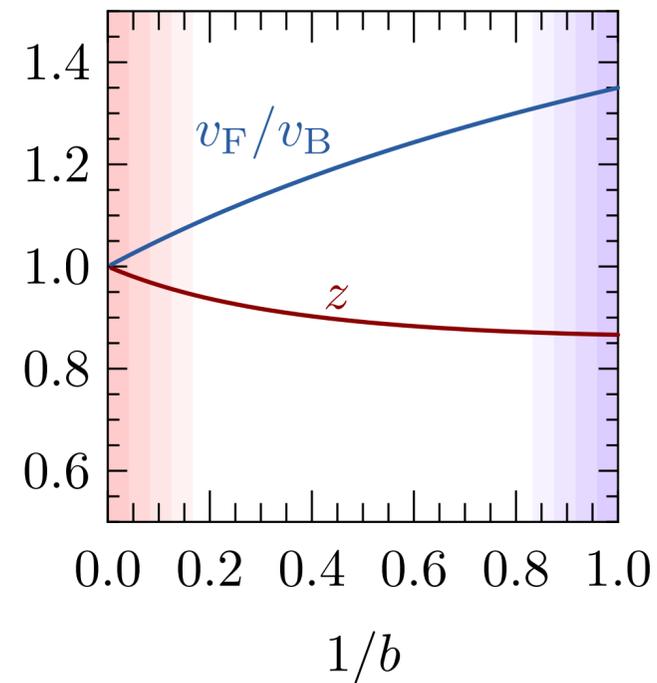
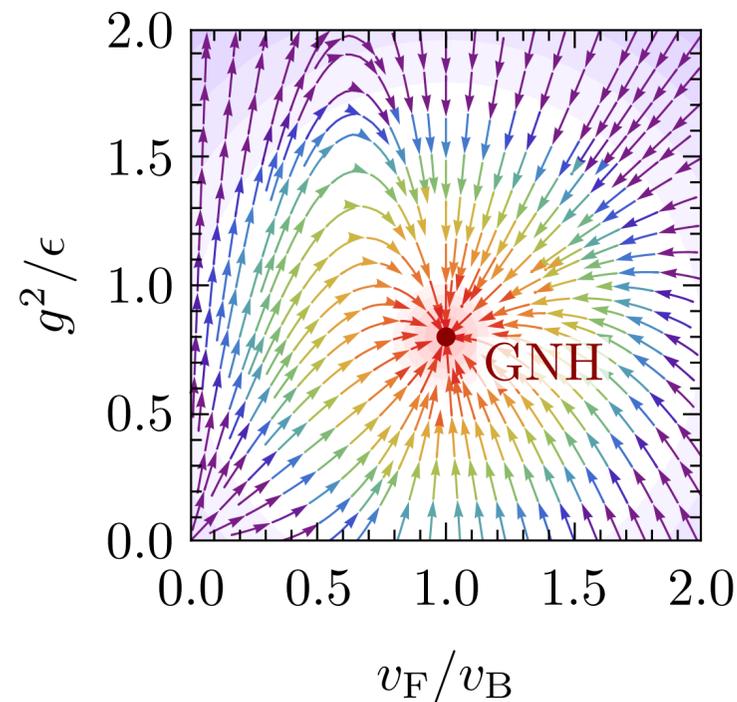
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... in  $d = 3 - \epsilon$  spatial dimensions



⇒ Emergent Lorentz symmetry

... characterized by set of universal critical exponents

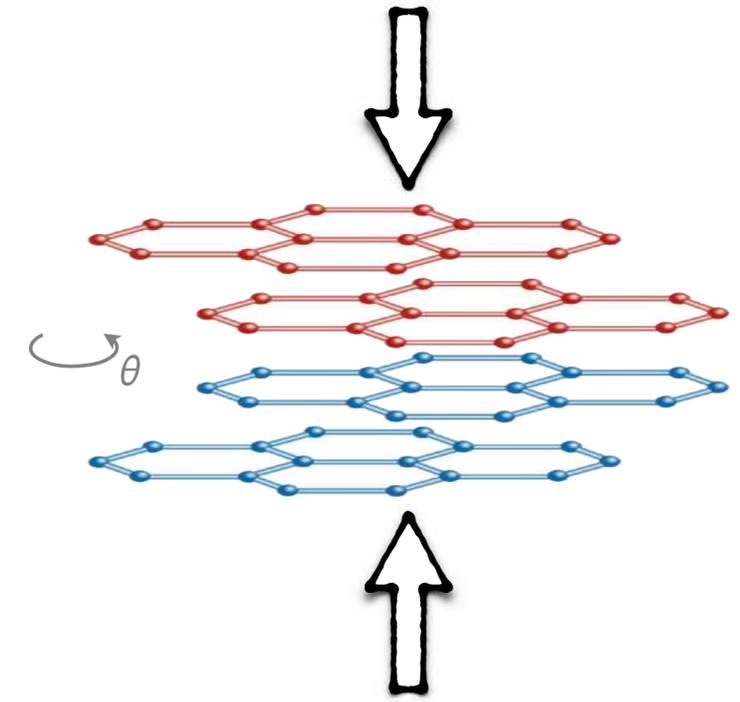
$$z = 1 \quad \nu = \frac{1}{2} + \frac{21\epsilon}{55} + \mathcal{O}(\epsilon^2)$$

$$\eta_\psi = \frac{3\epsilon}{10} + \mathcal{O}(\epsilon^2) \quad \eta_\phi = \frac{4\epsilon}{5} + \mathcal{O}(\epsilon^2)$$

⇒ Gross-Neveu-Heisenberg universality

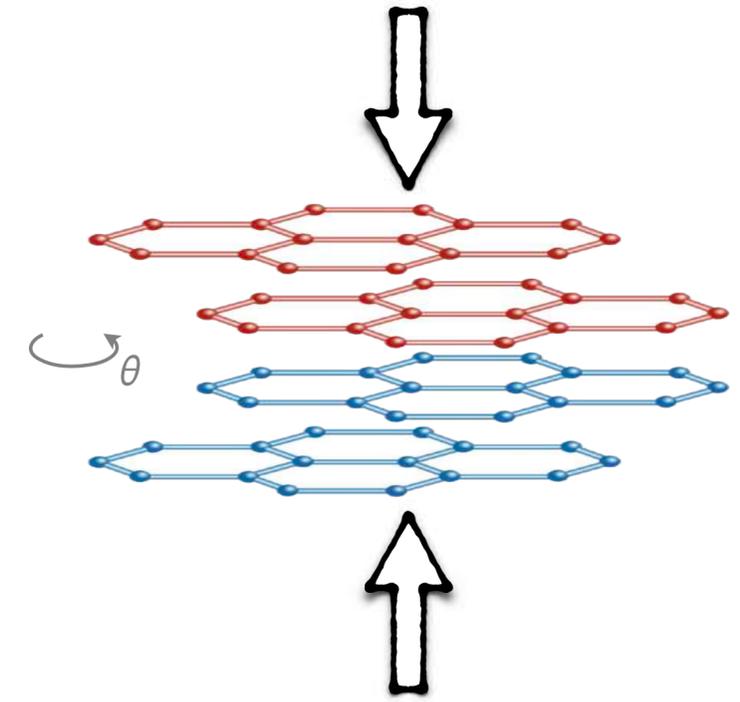
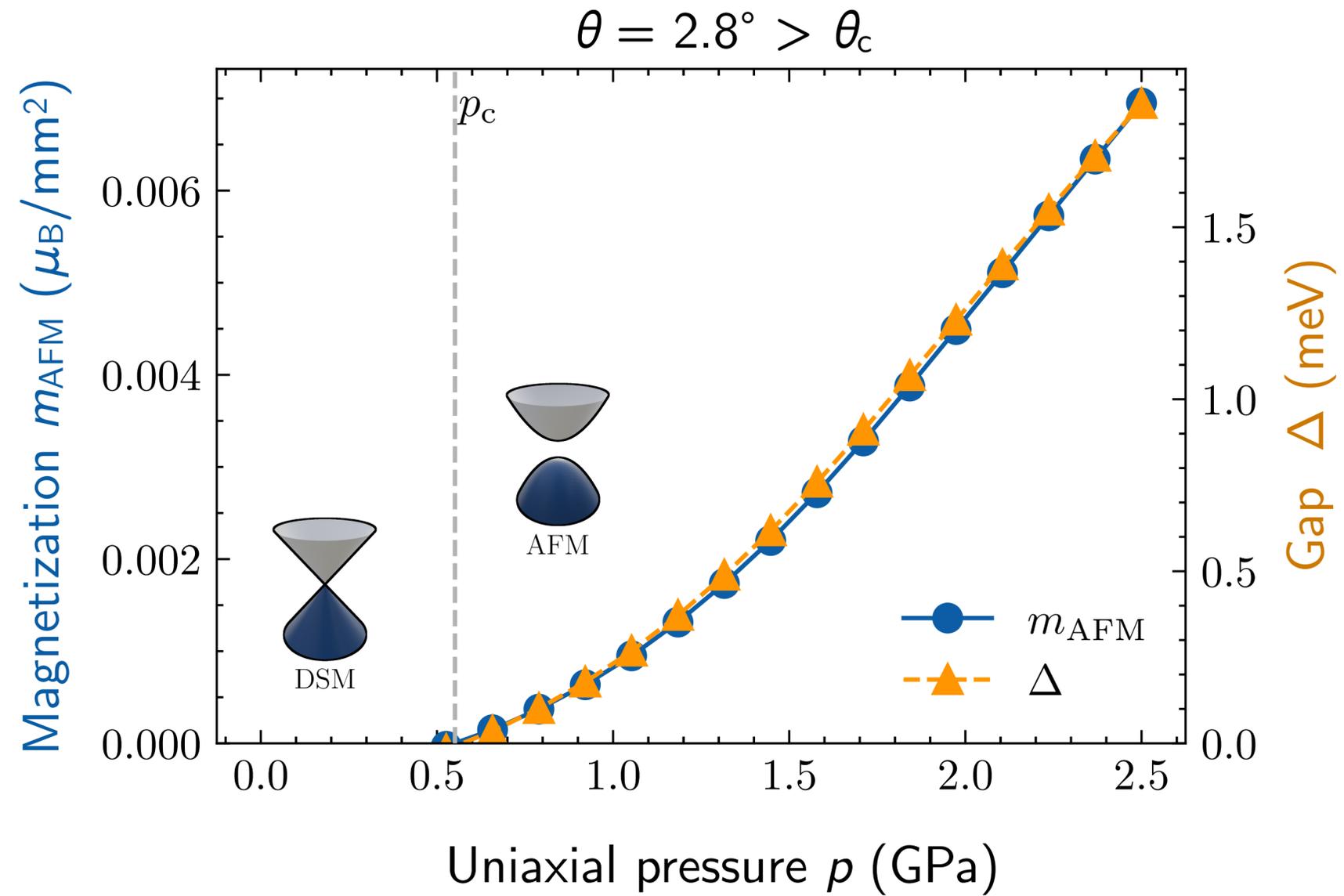
[Biedermann & LJ, arXiv:2509.04561]

# Pressure-tuned transition



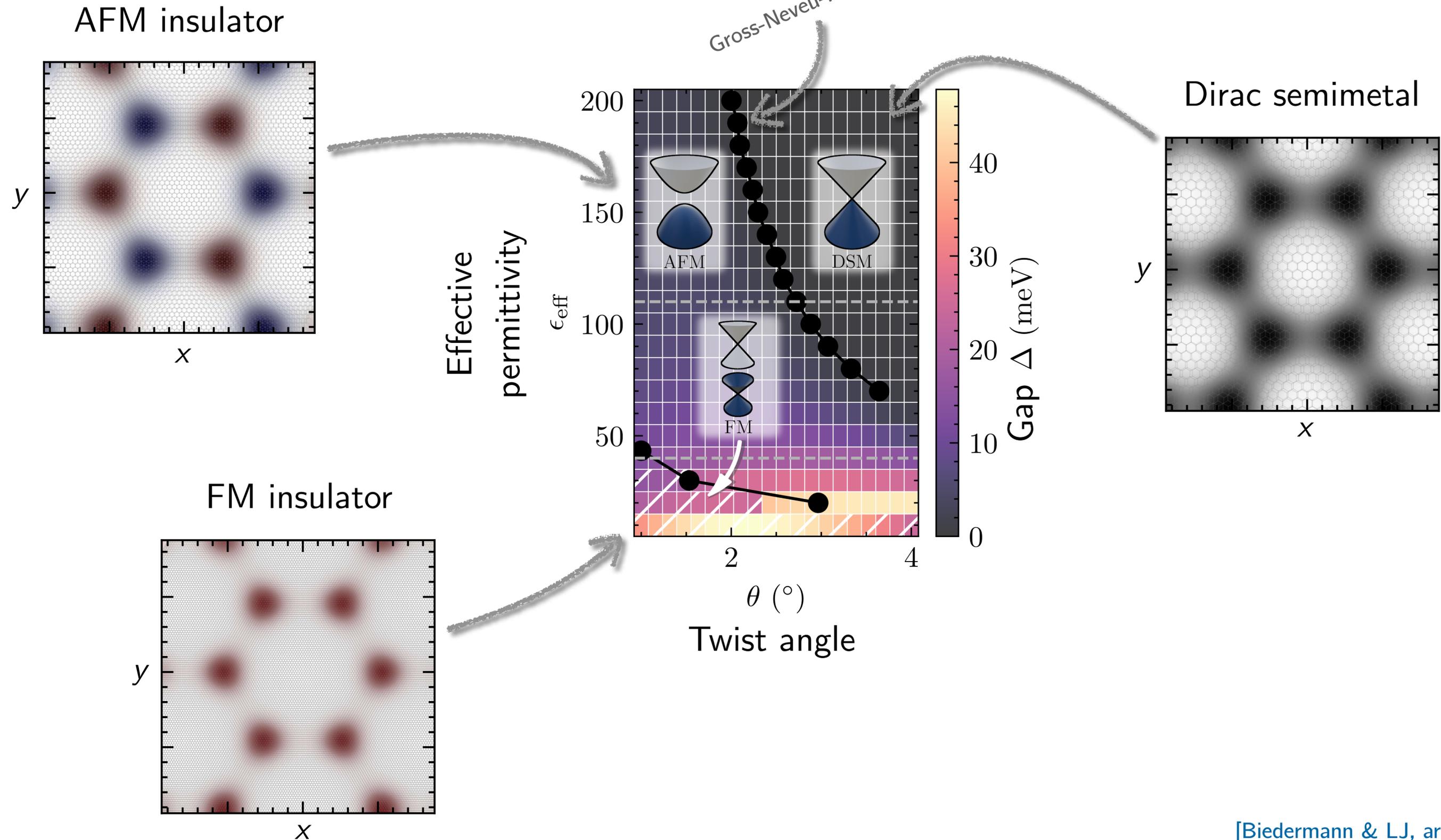
... from selfconsistent Hartree-Fock analysis

# Pressure-tuned transition



... from selfconsistent Hartree-Fock analysis

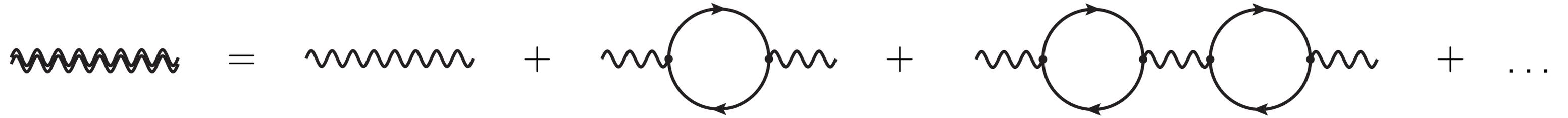
# Conclusions





# Internal screening

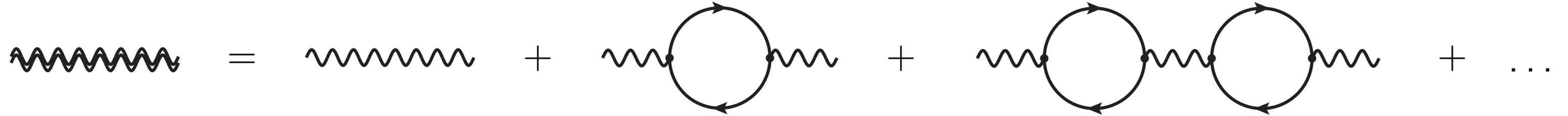
Screened Coulomb interaction:



... RPA in atomistic tight-binding model

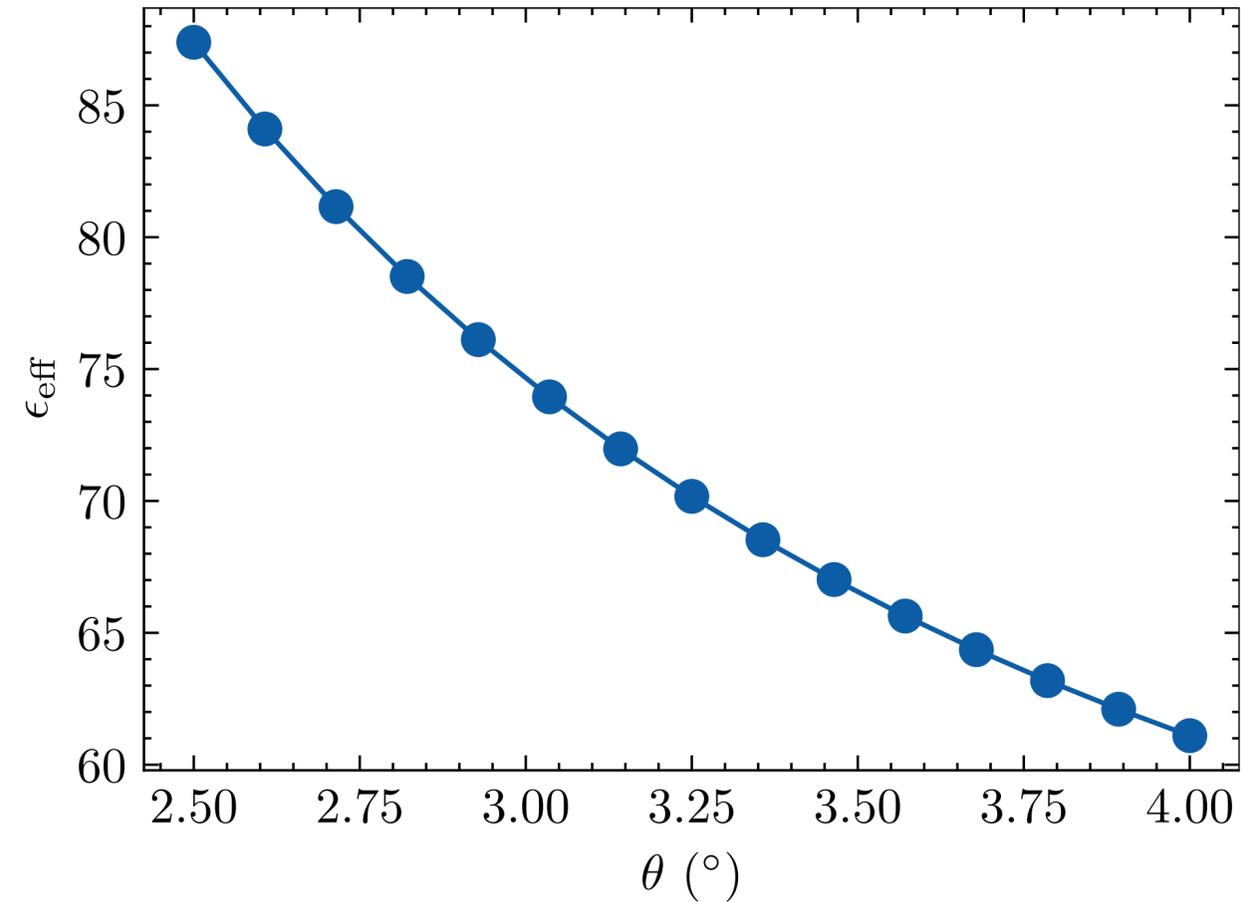
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Screened Coulomb interaction:



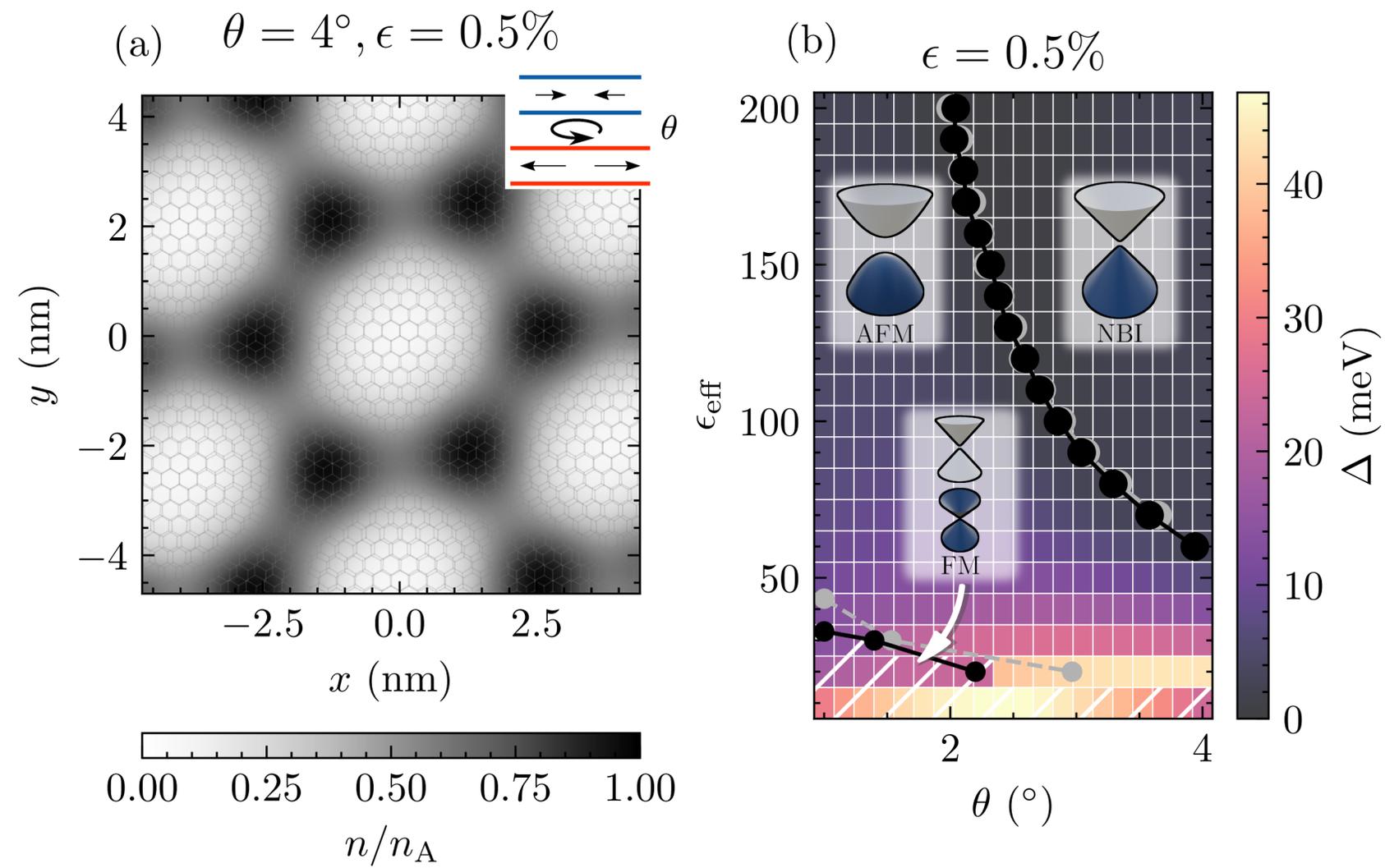
... RPA in atomistic tight-binding model

Effective dielectric permittivity:

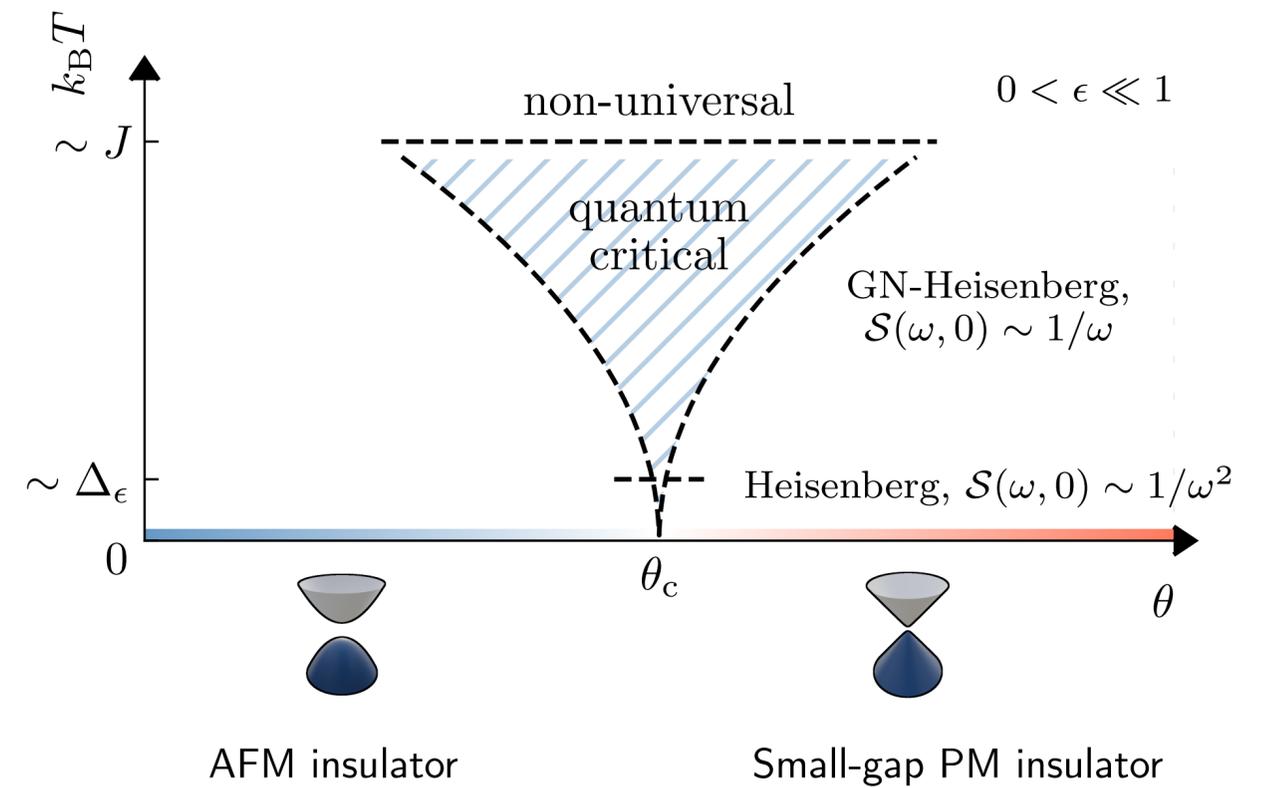
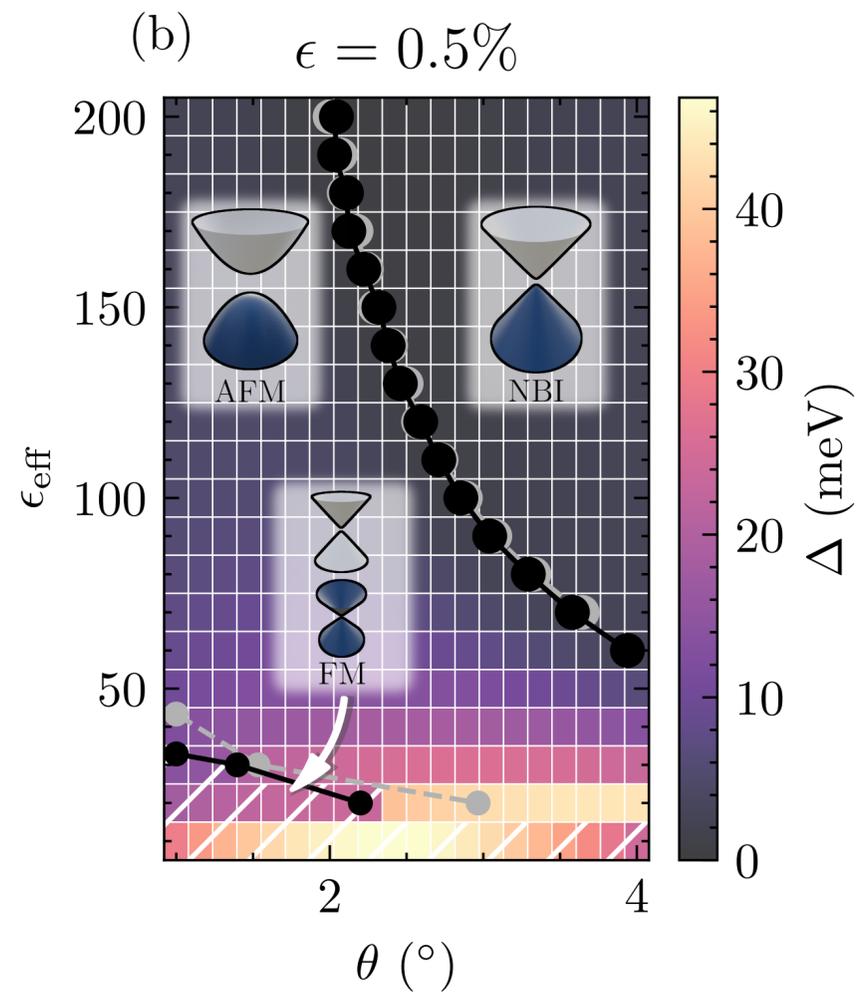
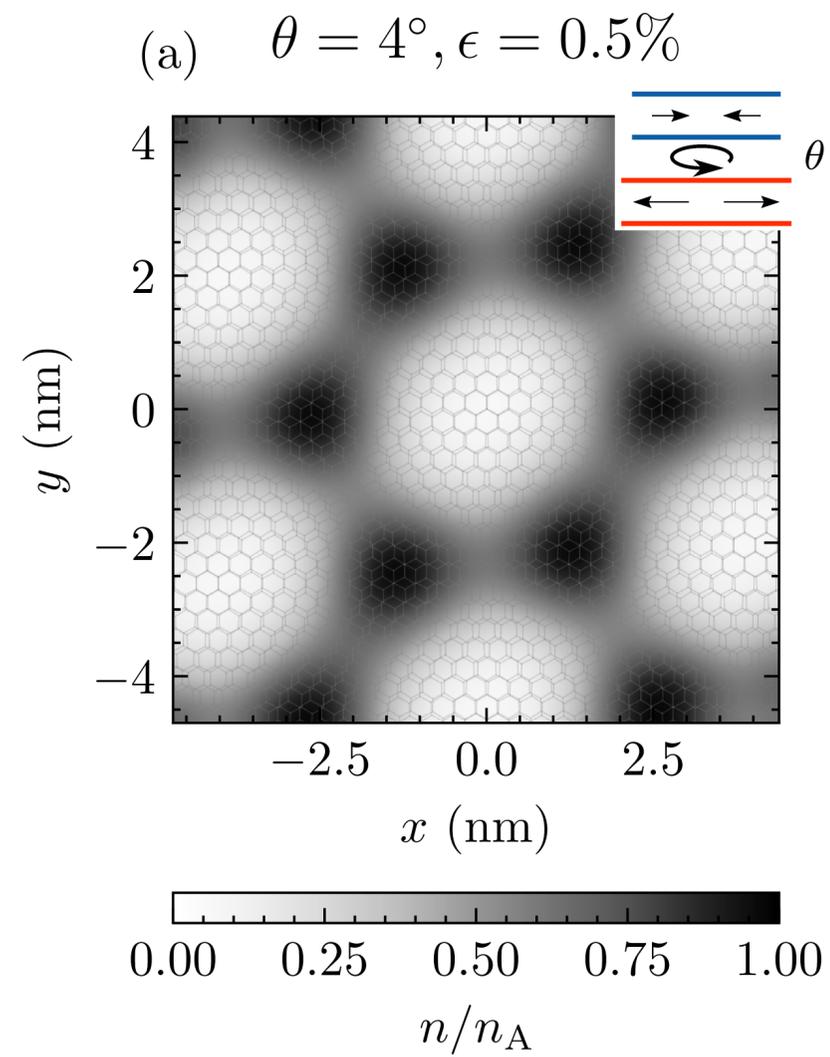


[Biedermann & LJ, arXiv:2509.04561]

# Twisted double bilayer WSe<sub>2</sub>: Heterostrain

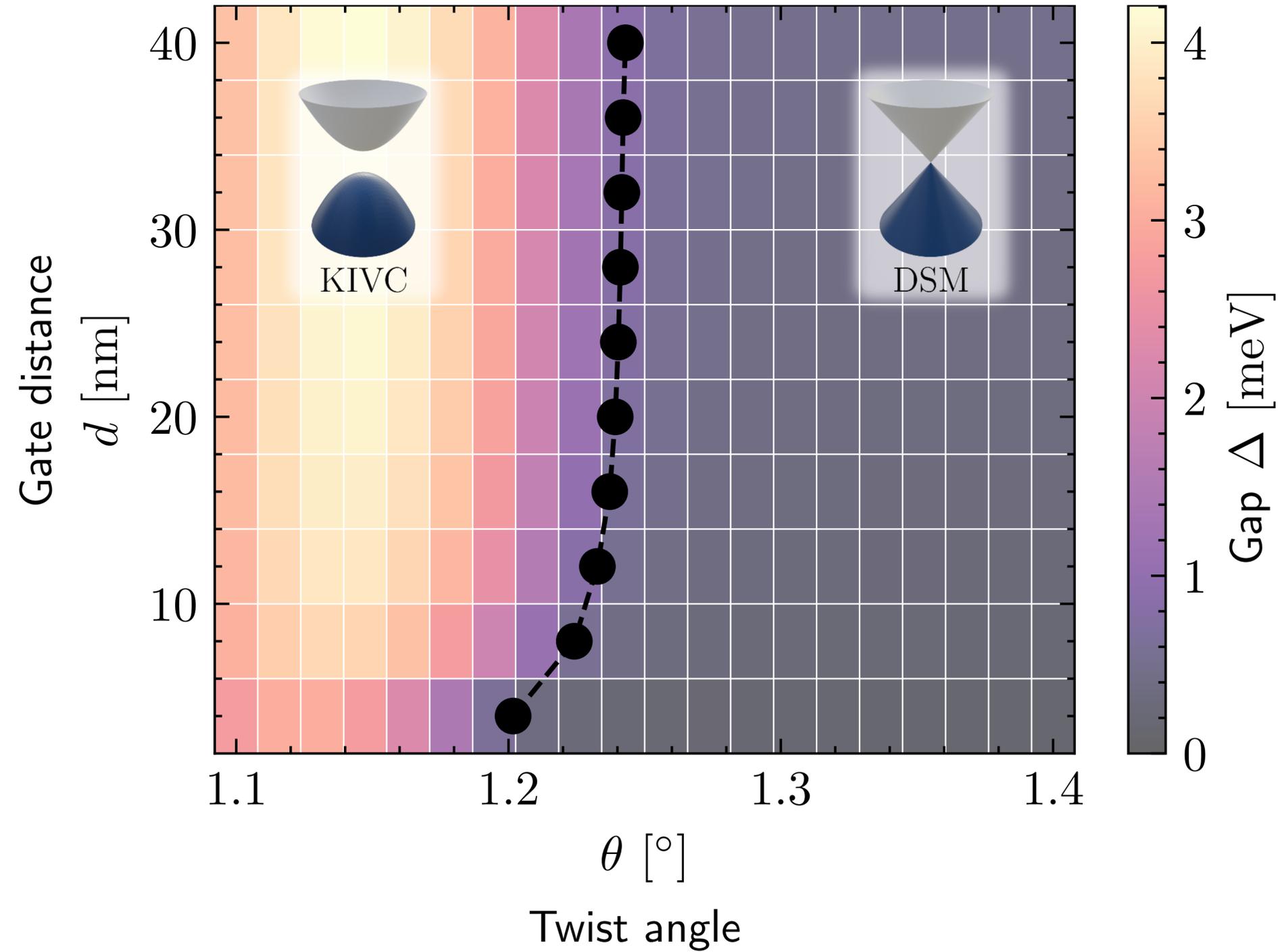
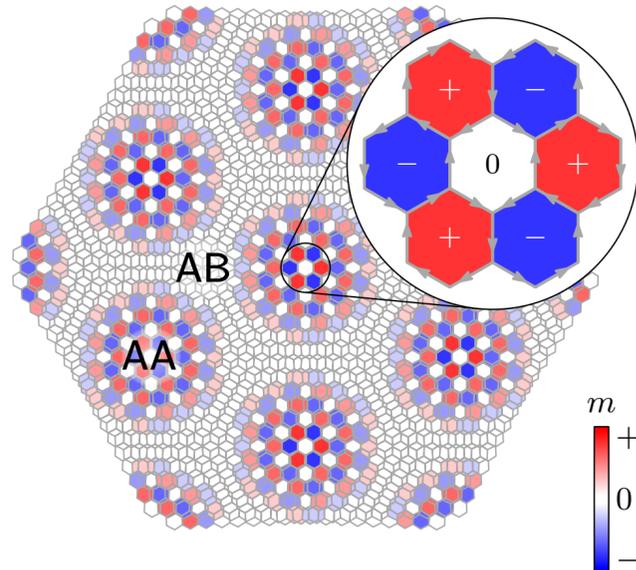


# Twisted double bilayer WSe<sub>2</sub>: Heterostrain



# Twisted bilayer graphene

Kramers intervalley-coherent insulator



[Biedermann, LJ, PRB '25]

# Gross-Neveu-Heisenberg universality

Universality field theory:

$$\mathcal{L} = \bar{\psi} \gamma^\mu \partial_\mu \psi + g (\bar{\psi} \vec{\sigma} \psi) \cdot \vec{\varphi} + \dots$$

Critical exponents:

