## **Ultracold Quantum Gases**

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**Problem set 5** 

to be discussed at 13:00h on 7th of July, 2011

see http://tu-dresden.de/physik/tqo/lehre

## "Black Holes" in Bose-Einstein condensates

In the hydrodynamic limit the solution of the time dependent Gross-Pitaevskii equation with external potential  $V(\mathbf{x})$  is written as  $\psi(\mathbf{x},t) = \sqrt{\rho(\mathbf{x},t)}e^{i\phi(\mathbf{x},t)}$ . Then one derives equations for  $\rho$  and  $\phi$ , neglecting gradients of  $\rho$ .

- a) Linearise these equations around a stationary solution  $\psi_s(\mathbf{x},t) = \sqrt{\rho_s(\mathbf{x})}e^{i\phi_s(\mathbf{x})-i\mu t}$  of the GP equation; i.e. write  $\rho = \rho_s + \delta\rho$  and  $\phi = \phi_s + \delta\phi$  with small  $\delta\rho$  and  $\delta\phi$ . Notice that in contrast to the lectures, we here allow for a velocity field  $\mathbf{V}(\mathbf{x}) = \frac{\hbar}{m}\nabla\phi_s(\mathbf{x})$  different from zero.
- b) Eliminate  $\delta \rho$  to derive of a wave equation for  $\delta \phi$ . Show that it can be written in the form  $\partial_{\mu}(\sqrt{-g}g^{\mu\nu}\partial_{\nu}\delta\phi)=0$  ( $\mu,\nu=0,\ldots,3$  and  $\partial_{0}=\partial_{t},\partial_{1}=\partial_{x},\ldots$ , summation convention). Here,  $g_{\mu\nu}$  is the  $(4\times 4)-$  metric tensor

$$(g_{\mu\nu}) = c \begin{pmatrix} -(c^2 - \mathbf{V}^2) & -\mathbf{V}^T \\ -\mathbf{V} & \mathbf{1} \end{pmatrix},$$

with  $c = c(\mathbf{x}) = \frac{\hbar}{m} \sqrt{4\pi a \rho_s(\mathbf{x})}$  the local speed of sound and  $\mathbf{V} = \mathbf{V}(\mathbf{x}) = \frac{\hbar}{m} \nabla \phi_s(\mathbf{x})$  the local velocity field.

- c) Interpret the above result for the dynamics of phase fluctuations  $\delta \phi$ . Draw an analogy to black hole physics.
- d) See Garay, Anglin, Cirac, Zoller: Phys. Rev. Lett. **85**, 6443 (2000) and Phys. Rev. A **63**, 023611 (2001) (find these articles on the materials web page).

Please turn over!

## **Bogoljubov transformation**

In order to discuss excitations from a condensate of an interacting homogenous gas  $(V=L^3)$ , we write  $\hat{\psi}(\vec{r})=\psi_0+\hat{\phi}(\vec{r})$ , with the condensate  $\psi_0=\sqrt{n_0}$  a real number and  $n_0$  the condensate density. The total energy is expanded in powers of  $\hat{\phi}$ . Choose momentum representation and show that the second order contribution  $H_2$  with the usual interaction parameter  $g=\frac{4\pi\hbar^2 a}{m}$  reads

$$H_2 = \sum_{\vec{k}} ' \left\{ \left( \frac{\hbar^2 k^2}{2m} + g n_0 \right) a_{\vec{k}}^+ a_{\vec{k}}^- + \frac{g n_0}{2} \left( a_{\vec{k}} a_{-\vec{k}}^- + a_{\vec{k}}^+ a_{-\vec{k}}^+ \right) \right\}.$$

Define new Bosonic quasi-particle creation and annihilation operators  $\alpha_{\vec{k}}, \alpha_{\vec{k}}^+$  through the ansatz  $a_{\vec{k}} = u_{\vec{k}} \alpha_{\vec{k}} - v_{\vec{k}} \alpha_{-\vec{k}}^+$ . Determine the coefficients  $u_{\vec{k}}, v_{\vec{k}}$  such that  $H_2$  becomes diagonal  $H_2 = \sum_{\vec{k}} {}' E_{\vec{k}} \alpha_{\vec{k}}^+ \alpha_{\vec{k}} + \Delta C$  and read off the quasi-particle excitation energies  $E_{\vec{k}}$ .

## **Attractive interaction II**

Discuss the quasiparticle excitation spectrum above for the case of attractive interaction a < 0. Consider finite values of  $N_0$  and V, and also the thermodynamic limit.